Earth Science Applications of Space Based Geodesy DES-7355 Tu-Th 9:40-11:05

Seminar Room in 3892 Central Ave. (Long building)

Bob Smalley Office: 3892 Central Ave, Room 103 678-4929 Office Hours – Wed 14:00-16:00 or if I'm in my office.

http://www.ceri.memphis.edu/people/smalley/ESCI7355/ESCI_7355_Applications_of_Space_Based_Geodesy.html



Determining <u>Strain</u> or strain rate from <u>Displacement</u> or velocity field

$$u_{i} = t_{i} + \frac{\partial u_{i}}{\partial X_{j}} X_{j} = t_{i} + D_{ij}X_{j} = t_{i} + (E_{ij} + W_{ij})X$$
Deformation tensor
$$E_{ij} = \frac{1}{2} (D_{ij} + D_{ji})$$

$$W_{ij} = \frac{1}{2} (D_{ij} - D_{ji})$$
Strain (symmetric) and
Rotation (anti-symmetris)
tensors

x

ū

Write it out $u_i = t_i + D_{ij}X_j$

Deformation tensor is not symmetric, have to keep d_{xy} and d_{yx} .

$$\begin{pmatrix} u_x \\ u_y \end{pmatrix} = \begin{pmatrix} t_x \\ t_y \end{pmatrix} + \begin{pmatrix} d_{xx} & d_{xy} \\ d_{yx} & d_{yy} \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

Again - this is "wrong way around"

We know <u>u</u> and <u>x</u> and want <u>t</u> and d_{ij} .

$$u_{i} = t_{i} + D_{ij}X_{j}$$
$$\begin{pmatrix} u_{x} \\ u_{y} \end{pmatrix} = \begin{pmatrix} t_{x} \\ t_{y} \end{pmatrix} + \begin{pmatrix} d_{xx} & d_{xy} \\ d_{yx} & d_{yy} \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

So rearrange it



Now we have 6 unknowns and 2 equations

So we need at least 3 data points That will give us 6 data

$$\begin{pmatrix} u_{x_1} \\ u_{y_1} \\ u_{x_2} \\ u_{y_2} \\ u_{y_2} \\ u_{x_3} \\ \vdots \\ u_{y_3} \\ \vdots \\ u_{x_n} \\ u_{y_n} \end{pmatrix} = \begin{pmatrix} 1 & 0 & x_1 & y_1 & 0 & 0 \\ 0 & 1 & 0 & 0 & x_1 & y_1 \\ 1 & 0 & x_2 & y_2 & 0 & 0 \\ 0 & 1 & 0 & 0 & x_2 & y_2 \\ 1 & 0 & x_3 & y_3 & 0 & 0 \\ 0 & 1 & 0 & 0 & x_3 & y_3 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & 0 & x_n & y_n & 0 & 0 \\ 0 & 1 & 0 & 0 & x_n & y_n \end{pmatrix} \begin{pmatrix} t_x \\ t_y \\ d_{xx} \\ d_{xy} \\ d_{yx} \\ d_{yy} \end{pmatrix}$$

And again - the more the merrier - do least squares.

For straín rate Take tíme derívatíve of all terms.

But be careful

Straín rate tensor ís NOT tíme derívatíve of straín tensor.

Spatial (Eulerian) and Material (Lagrangian) Coordinates and the Material Derivative

Spatial description picks out a particular location in space, x.

Material description picks out a particular piece of continuum material, X.

So we can write

$$x = x(A,t) \qquad \qquad x(A,0) = A$$

x is the position now (at time t) of the section that was initially (at time zero) located at A.

or

$$A = A(x,t) \qquad \qquad A(A,0) = A$$

A was the initial position of the particle now at x

$$x[A(x,t),t] = x$$
 $A[x(A,t),t] = A$

We can therefore write

$$f[x(A,t),t] = F(A,t) \qquad \qquad f(x,t) = F[A(x,t),t]$$

Next consider the derivative (use chain rule)

$$\frac{\partial}{\partial A}F(A,t) = f\left[x(A,t),t\right] = \frac{\partial f}{\partial x}\Big|_{A}\frac{\partial x}{\partial A}$$

$$\frac{\partial}{\partial t}F(A,t) = f\left[x(A,t),t\right] = \frac{\partial f}{\partial x}\Big|_{A}\frac{\partial x}{\partial t} + \frac{\partial f}{\partial t}\Big|_{A}$$

Define Material Derivative

$$\frac{\partial}{\partial t}F(A,t) = f\left[x(A,t),t\right] = \frac{\partial f}{\partial x}\Big|_{A}\frac{\partial x}{\partial t} + \frac{\partial f}{\partial t}\Big|_{A}$$

$$\frac{DF(A,t)}{Dt} = \frac{\partial F(A,t)}{\partial t} \bigg|_{A = A(x,t)}$$

$$\frac{Df(A,t)}{Dt} = \frac{\partial f(x,t)}{\partial t} + v(x,t)\frac{\partial f(x,t)}{\partial x}$$

Vector version

$$\frac{D\vec{f}}{Dt} = \frac{\partial \vec{f}(x,t)}{\partial t} + \vec{v}(x,t) \bullet \nabla \vec{f}(x,t)$$





Consider bar steadily moving through a roller that thins the bar



Examine velocity as a function of time of cross section A



The velocity will be constant until the material in A reaches the roller

At which point it will speed up (and get a little fatter/ wider, but ignore that as second order)

After passing through the roller, its velocity will again be constant



If one looks at a particular position, x, however the velocity is constant in time. So for any fixed point in space $\frac{\partial v(x,t)}{\partial t} = 0$ So the acceleration seems to be zero (which we know it is not)



The problem is that we need to compute the time rate of change of the matería which is moving through space and deforming (not rigid body) (we want/need our reference frame to be with respect to the material, not the coordinate system.



We know acceleration of material is not zero.

 $\frac{Df(A,t)}{Dt} = \frac{\partial f(x,t)}{\partial t} + v(x,t)\frac{\partial f(x,t)}{\partial x}$

Term gives acceleration as one <u>follows the material</u> through space (have to consider same material at t₁ and t₂)

Various names for this derivative

Substantive derivative Lagrangian derivative Material derivative Advective derivative Total derivative

GPS and deformation

Now we examine <u>relative</u> movement between sites



Strain-rate sensitivity thresholds (schematic) as functions of period



GPS and INSAR detection thresholds for 10-km baselines, assuming 2-mm and 2-cm displacement resolution for GPS and INSAR, respectively (horizontal only).

http://www.iris.iris.edu/USArray/EllenMaterial/assets/es_proj_plan_lo.pdf, http://www.iris.edu/news/IRISnewsletter/EE.Fall98.web/plate.html

Strain-rate sensitivity thresholds (schematic) as functions of period



Post-seismic deformation (triangles), slow earthquakes (squares), long-term aseismic deformation (diamonds), preseismic transients (circles), and volcanic strain transients (stars).

http://www.iris.iris.edu/USArray/EllenMaterial/assets/es_proj_plan_lo.pdf, http://www.iris.edu/news/IRISnewsletter/EE.Fall98.web/plate.html

Study deformation at two levels

-Kínematics – describe motions (Have to do this first)

-Dynamics relate motions (kinematics) to forces (physics) (Do through rheology/constitutive relationship/model. Phenomenological, no first principle prediction)

Simple rheological models





Apply constant stress, σ , to a viscoelastic material. Record deformation (strain, ϵ) as a function of time. ϵ increases with time.

http://hcgl.eng.ohio-state.edu/~ce552/3rdMatO6_handout.pdf



Maintain constant strain, record load stress needed. Decreases with time. Called relaxation.

viscoelastic

Kelvín rheology





Handles creep and recovery fairly well Does not account for relaxation



http://hcgl.eng.ohio-state.edu/~ce552/3rdMat06_handout.pdf



viscoelastic Maxwell rheology $\sigma = \sigma_1 = \sigma_2$ $\varepsilon = \varepsilon_1 + \varepsilon_2$ $\dot{\varepsilon} = \frac{\sigma}{\mu} + \frac{\dot{\sigma}}{k}$

Handles creep badly (unbounded) Handles recovery badly (elastic only, instantaneous) Accounts for relaxation fairly well



http://hcgl.eng.ohio-state.edu/~ce552/3rdMat06_handout.pdf



http://hcgl.eng.ohio-state.edu/~ce552/3rdMat06_handout.pdf www.mse.mtu.edu/~wangh/my4600/chapter4.ppt

víscoelastíc Standard línear/Zener



Instantaneous elastic strain when stress applied Strain creeps towards limit under constant stress Stress relaxes towards limit under constant strain Instantaneous elastic recovery when strain removed Followed by gradual recovery to zero strain

http://hcgl.eng.ohio-state.edu/~ce552/3rdMatO6_handout.pdf www.mse.mtu.edu/~wangh/my4600/chapter4.ppt

víscoelastíc Standard línear/Zener



Two tíme constants - Creep/recovery under constant stress - Relaxatíon under constant straín

http://hcgl.eng.ohio-state.edu/~ce552/3rdMatO6_handout.pdf www.mse.mtu.edu/~wangh/my1600/chapter1.ppt



Viscoelastic Response to Long-Term Loading



Can make arbitrarily complicated to match many deformation/ strain/time relationships

Three types faults and plate boundaries

- Faults -

Strike-slip Thrust Norma

- Plate Boundary -

Strike-slip Convergent Divergent

How to model

Elastic Viscoelastic

Half space Layers Inhomogeneous



2-D model for strain across strike-slip fault in elastic half space.

Fault is locked from surface to depth D, then free to infinity.

Far-field displacement, V, applied.



w(x) is the equilibrium displacement parallel to y at position x.

|w| is 50% max at x/D=.93; 63% at x/D=1.47 & 90% at x/ D=6.3



Effect of fault díp.

The fault is locked from the surface to a depth D (not a down dip length of D).

The fault is free from this depth to infinity.



Surface deformation pattern is <u>SAME</u> as for vertical fault, but centered over <u>down</u> <u>dip end</u> of dipping fault.

Dip estimation from center of deformation pattern to surface trace and locking depth.











Interseismic velocities in southern California from GPS



Fault parallel velocities for northern and southern "swaths".

Total change in velocity ~42mm/yr on both.

Meade and Hager, 2005



Residual (observed-model) velocíties for block fault model (faults in grey)

41

Modeling velocities in California

$$\vec{V}(\vec{r}) = \Omega(\vec{r}) \times \vec{r} + \sum_{f=1}^{F} G \bullet s_f$$

 Ω is the angular velocity vector

effect of interseismic strain accumulation is given by an elastic Green's function G response to backslip distribution, s, on each of, f, faults.

Modeling Broadscale Deformation From Plate Motions and Elastic Strain Accumulation, Murray and Segall, USGS NEHRP report.

$$\vec{V}(\vec{r}) = \Omega(\vec{r}) \times \vec{r} + \sum_{f=1}^{F} G \bullet s_f$$

In general, the model can accommodate zones of distributed horizontal deformation if Ω varies within the zones

latter terms can account both for the Earth's sphericity and viscoelastic response of the lower crust and upper mantle.

Modeling Broadscale Deformation From Plate Motions and Elastic Strain Accumulation, Murray and Segall, USGS NEHRP report.

$$\sum_{f=1}^{F} G \bullet s_{f} \to -\frac{a}{\pi} \sum_{f=1}^{F} \Delta \omega_{f} \sin \phi_{f} \tan^{-1} \left(\frac{d_{f}}{a(\phi - \phi_{f})} \right)$$

Where a is the Earth radius distance from each fault located at ϕ_f is a $(\phi - \phi_f)$. Each fault has deep-slip rate $a\Delta\omega_f \sin\phi_f$, where $\Delta \omega_f$ is the difference in angular velocity rates on either side of the fault.