

Effects on Seismic Absorption due to Changed Pore Surface

Properties Resulting from Exposure to Propanol:

A Study Utilizing Artificial Glass Cracks

by

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ABSTRACT

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Effects on Seismic Absorption due to Changed Pore Surface Properties Resulting From Exposure to Propanol: A Study Utilizing Partially Saturated Artificial Glass Cracks

Thesis directed by Professor Hartmut Spetzler

The study of multi-phase flow has been expanding in recent times due to a deepening appreciation of, and a greater ability to predict, the system's complex behavior. This thesis shows how altering the solid phase in the presence of gas and liquid phases produces definite changes in the measured quantities, attenuation and stiffness, for a partially saturated, sinusoidally stressed, artificial glass crack.

The two key means of dissipation for this system are energy lost due to viscous effects and energy lost due to restricted movement of the contact line, the interface between the gas, liquid, and solid phases. The viscous effects are relatively well known and simple, but contact line movement, involving the interaction of the three phases, requires a broader understanding. The chemistry between the solid phase and liquid phase and the intermolecular forces in the liquid phase become exceedingly important.

Model predictions are tested against experimental data obtained on an attenuation spectrometer. The spectrometer produces oscillatory deformation of a partially saturated artificial crack made of glass from .001 to 100 Hz with amplitudes up to .5 μm . The phase of the displacement relative to an elastic standard, a measure of dissipated energy, is deduced after digitizing and fitting the measured signal.

Viscous effects are found to explain the measured attenuation above 10 Hz and restricted contact line movement does its job below 1 Hz. Their combination allows for nearly perfect agreement over the entire frequency range. However, large discrepancies exist between predicted and measured stiffness.

ACKNOWLEDGMENTS

Many people are responsible for the completion of this thesis. I would like to thank Hartmut Spetzler for all of his ideas and constant encouragement. Ivan Getting has been an invaluable resource in the experimental design and implementation of the attenuation meter. I would also like to thank Pamela Burnley and Bill Waite for attempting to answer most of my questions. A special thanks go to Rainer Moerig for all that I learned during our endless conversations and my father, Jeff Boyd, whose programming experience and keen mind helped reduce my programming bugs to theoretical contradiction. I am also grateful to my mom and sister who also helped keep me together.

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Chapter I

INTRODUCTION

Characterizing upper crustal response to imposing pressure gradients provides valuable information to all who use pressure waves to image subsurface phenomena. Much effort has been expended in making observations and developing theories about how energy is propagated in the subsurface. This thesis is primarily concerned with the role of partial fluid saturation in the dissipation of seismic energy.

The partially saturated system has been studied extensively yet little is known about the role of the contact line, the interface between the three phases, gas, liquid, and solid. Through the application of experimental and theoretical techniques, this thesis will show how the properties of the solid phase at its surface affect the mobility of the contact line and thereby substantially affect the magnitude and frequency dependence of contact line and viscosity related dissipation. The measurements are performed on a cylindrical crack geometry with Propanol as the surface affecting agent.

Background

Many precious resources lie below our feet. There are reservoirs full of water, oil and gas. There are faults and folds and the remains of our ancestors. Approximately 100% of the Earth's mass is below our feet. The scientific, cultural, and economic evolution of our species depends on an understanding of the subsurface. There are several ways to gather information about the material within the Earth depending on the depth and attribute of inquiry and the resources available to pursue such inquiry. Direct methods such as bore holes provide a wealth of direct observations that allow for characterization on a sub-meter scale. However, they are relatively expensive and suffer an extremely low resolution on larger scales. Direct methods also become unfeasible with depths greater than several kilometers. Indirect methods allow a much greater resolution on a much larger scale. Measurements can be made on phenomena from the crust to the core to distant stars and their planets. The indirect methods require a significant amount of theory to describe the observed phenomena and frequently result in non-unique interpretations. At the forefront of acquisition techniques for the earth's interior are the methods associated with seismic wave propagation.

Seismic waves are pressure waves that propagate through the subsurface. Their velocity and phase is dependent on the complex modulus of the medium the waves are passing through. Their frequency content is dependent on the depth of the source and the dispersion of the intervening material. Common velocities are on the order of several km/s and frequencies on the order of 1 Hz. These conditions make

measurement difficult in the lab because the samples are small compared to the wavelength of the passing waves. The sample's influence on the dissipation of seismic energy is virtually undetectable and so empirical development has relied on measurements made in the sonic and ultrasonic frequency ranges.

Important steps toward an understanding of wave propagation in porous rocks were made by Born and Owen in 1935. Measuring the width of a resonant peak in the sonic to sub-sonic frequency range, they determined the attenuation and discovered that the behavior of their rock samples possess a strong dependence on pore content. To explain this Biot published several papers outlining the theory behind the propagation of seismic waves in saturated porous solids (1956a; 1956b; 1962). Much of his theory is based on inertial effects of the pore content. But since its inception, the theory has not adequately modeled experimental results for velocity dispersion and attenuation (O'Connell and Budiansky, 1977; Mavko and Jizba, 1991). Bulau et al. (1984), Jones and Nur (1983), and Dvorkin et al. (1993) note that experimental observations show a frequency dependence that is in contradiction to Biot theory. With increasing oil viscosity in a Berea sandstone, measured peaks in attenuation and dispersions in velocity moved to lower frequencies (Winkler, 1983). Biot predicted their movement to higher frequencies. Nur and Simmons (1969) and O'Connell and Budiansky (1977) suggested that another mechanism more strongly dependent on the viscosity of the pore content was involved. This mechanism was referred to and is called local fluid flow.

During and after the development of local fluid flow theory (Murphy et al., 1986; Norris 1993; Gurevich and Lopotnikov, 1995) and because of improvements in

measurement techniques, additional patterns began to emerge. Scattering in dry rocks (Burridge and Chang, 1989; Shapiro et al., 1994) and partially saturated rocks, simulated (Gurevich et al., 1997), squirt flow between pores (Palmer and Traviolia, 1980; Murphy et al., 1986; Akbar et al., 1994), and contact line movement (Miksis, 1988; Waite et al., 1997) have been proposed.

Gurevich et al. (1997) combined Biot theory, local fluid flow, and scattering in a simulated stratified rock and found that in the seismic to sonic frequency range, local fluid flow and scattering dominate over Biot theory. They also found that the attenuation peak due to interlayer fluid flow occurred at lower frequencies than scattering. Akbar et al. (1994) modeled squirt flow between pores and local fluid flow within them and found good agreement to data measured by Paffenholz and Burkhardt (1989). Akbar et al. note that the resulting squirt flow attenuation occurs at lower frequencies than do losses due to local fluid flow. Ongoing measurements by Moerig et al. (1996) in the seismic frequency range begin to address the role of restricted contact line movement and subsequent energy loss. In their measurements the solid surfaces of an artificial crack made of microscope slides separated by thin wires are brought together at varying velocities causing pressure to be built up within the fluid and energy to be stored in the elastic spring mechanism of the sample. The fluid is displaced producing some combination of shear, contact line movement and change in shape of the meniscus depending on the chemical properties of the solid surface (Waite et al., 1997, Moerig et al., 1997). Moerig et al. (1996) made the observation that on surfaces that have been exposed to carbon based molecules, an increase in low frequency attenuation and an increase in stiffness occurs (figure I.1).

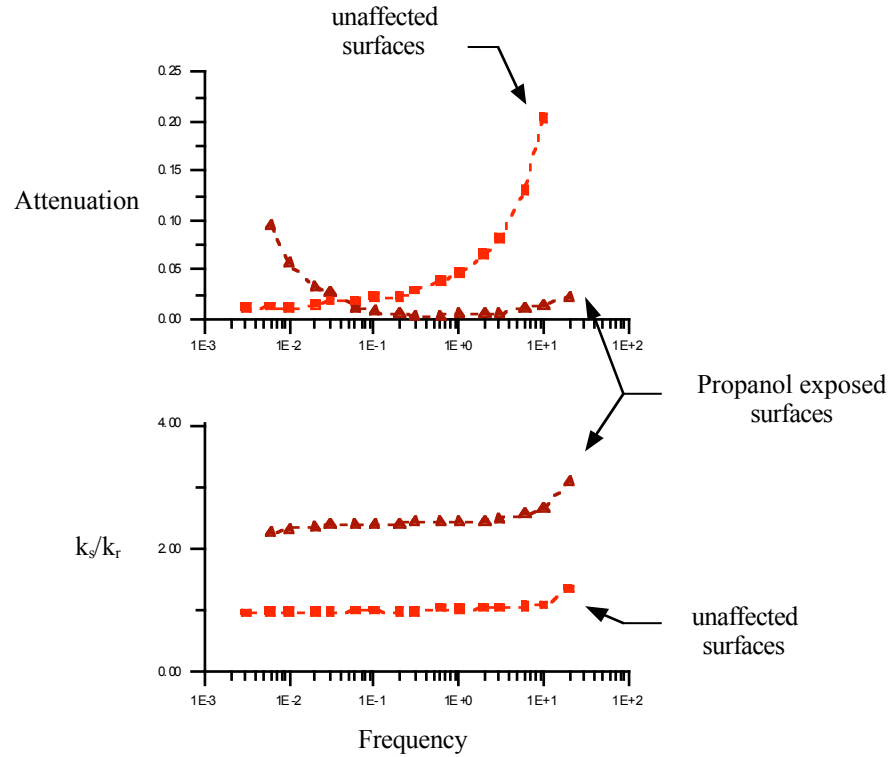


Figure I.1 Attenuation and relative stiffness ($k_r = 50 \text{ kN/m}$) versus frequency for artificial cracks partially saturated with water. Squares represent data taken on surfaces that have not been exposed to carbon based molecules. Triangles represent data taken on surfaces exposed to Propanol.

The types of mechanisms addressed in this thesis were originally investigated using real rock samples. The measured attenuation and stiffness of these samples were very difficult to interpret (Chelidze et al., 1996). There was no quantitative means for predicting the measured values. Since that time, efforts have been made to reduce the complexity of the system by studying single cracks. Observations by Moerig et al. (1996) utilized a rectangular crack where one side was fixed, the opposite side allowed to move and deform, and the orthogonal sides allowed to deform. The geometry for this sample was still fairly complex but pointed towards viscous relaxation and a definite dissipation mechanism related to a restricted contact

line. With this information, the single crack was simplified further such that the fluid film was nearly cylindrical. This geometry allows for quantitative interpretations and are the focus of this thesis. Future work will begin with measurements on cindered glass bead samples and then real rocks.

Chemistry

A silica surface that is exposed to a humid environment will adsorb water to produce hydroxyls (Iler, 1979). On a flat quartz surface, the surface density of hydroxyls are on average 4.6 OH nm^{-2} (Peri and Hensley, 1968). Amorphous silica has a slightly smaller surface density of hydroxyls. The hydroxyls allow hydrogen bonding to other molecules containing OH groups such as Water, H_2O (Iler, 1979), Methanol, CH_3OH (George, 1994), Acetone, $\text{C}_3\text{H}_5\text{OH}$, and Propanol, $\text{C}_3\text{H}_7\text{OH}$ (Appendix A). The most likely molecular configurations appear in figure I.2. Others have been noted (Iler 1979).

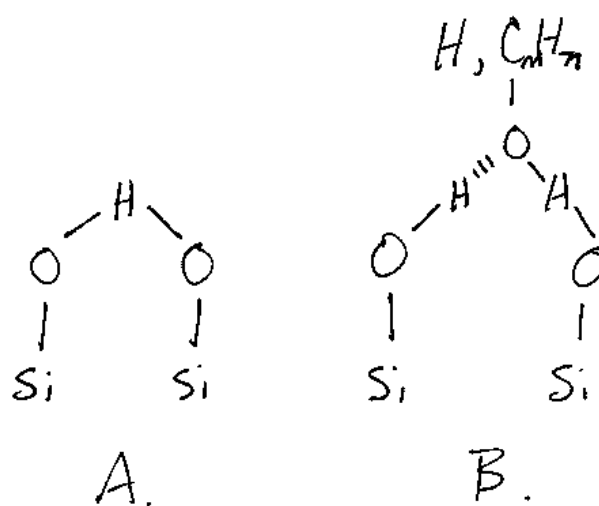


Figure I.2 A. Hydroxylated silica surface. B. Hydrogen bound water and carbon based molecules through the OH group.

Water reacts quite differently to a surface that contains carbon based groups as opposed to a simple hydroxylated surface. A simple hydroxylated surface allows polar interactions between water molecules and surface hydroxyls, a process that is considered hydrophilic. Carbon groups containing hydrogen are nonpolar and have no affinity for water producing a hydrophobic effect (Iler, 1979; Ben-Naim, 1980). In terms of a moving contact line this means that on a hydrophilic surface the water molecules have a relatively high probability of bonding with a surface site and moving some distance in some interval of time. As the surface sites become hydrophobic, the probability for water molecules to bond to the surface decreases resulting in a lower contact line velocity. The water molecules must either find a surface hydroxyl or replace a nonpolar molecule with a smaller probability.

Attenuation Spectrometer

Definition of $1/Q$

Attenuation, the inverse of the quality factor, Q , is defined as the fractional energy lost during one cycle of deformation of a material,

$$\frac{1}{Q} = \frac{\Delta W}{2\Delta W}. \quad (I.1)$$

The work in terms of stress and strain is

$$W = \int_0^t \sigma \partial \epsilon. \quad (I.2)$$

The phase with which the strain lags the stress is a measure of attenuation. To show this, let's begin by writing the stress as

$$\sigma = M^* \epsilon \quad (I.3)$$

where M^* is the complex modulus and ϵ is the strain. The strain can be expressed as

$$\epsilon = \epsilon_0 \sin(\omega t + \phi) \quad (I.4)$$

$$\text{and} \quad \partial \epsilon = \epsilon_0 \cos(\omega t + \phi) \partial \omega t \quad (I.5)$$

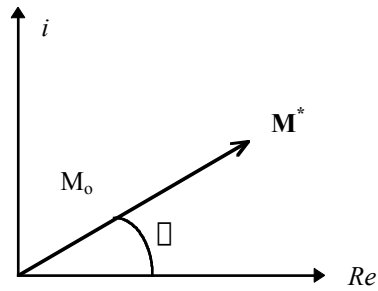


Figure I.3 Complex vector M^* in the real-imaginary plane. ϕ is the angle between M^* and the real axis and corresponds to the phase lag between stress and strain. The anelastic component of M^* is $M_0 \sin \phi$ and the elastic, $M_0 \cos \phi$

where ϵ_0 is the strain amplitude, ω is the circular frequency, and δ is the phase lag between stress and strain and the angle between the real axis and the complex vector M^* on the real-imaginary plane (figure I.3). The energy lost is the solution to eq. I.2 where the limits of integration are one full cycle,

$$\Delta W = \oint \sigma d\epsilon = \epsilon_0^2 M_0 \sin \delta. \quad (I.6)$$

The maximum energy stored is the solution to eq. I.2 where the limits of integration are over a quarter cycle,

$$W = \int_0^{\pi/2} \sigma d\epsilon = \frac{\epsilon_0^2 M_0 \cos \delta}{2}. \quad (I.7)$$

The fractional energy lost is then eq. I.6 divided by eq. I.7,

$$\frac{\Delta W}{W} = \frac{2 \sin \delta}{\cos \delta} = 2 \tan \delta. \quad (I.8)$$

For small δ , $\tan \delta$ is approximately equal to δ . The result is a direct relationship between the phase lag and the attenuation,

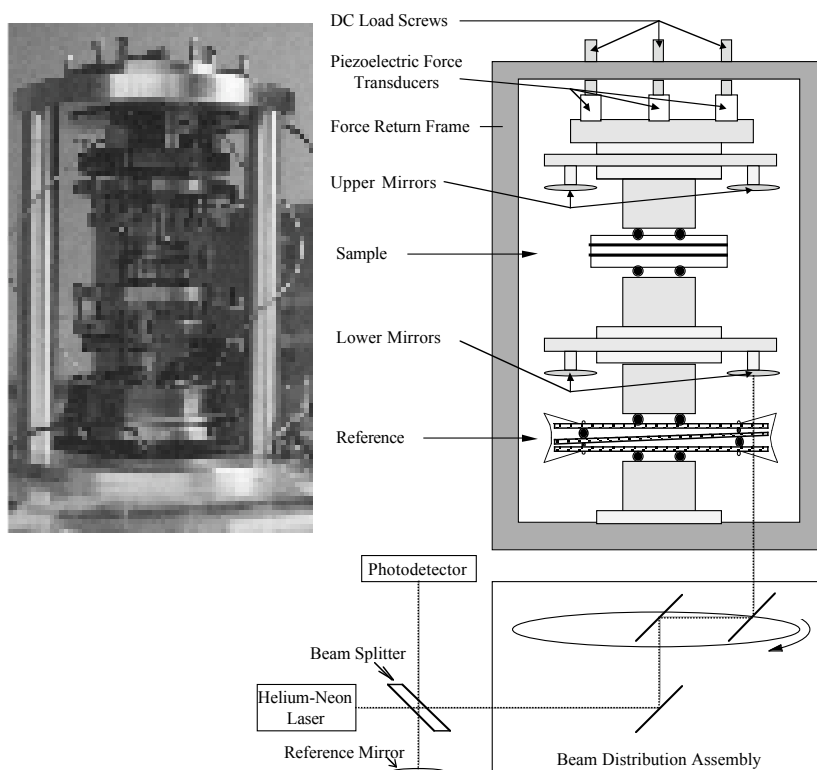
$$\frac{\Delta W}{W} = 2 \delta. \quad (I.9)$$

Apparatus

Work by Moerig et al. (1996) and continuing measurements made by members of this group utilizing an attenuation spectrometer (Cherry et al., 1996; figure I.4) will be the central focus of the experimental concepts presented in this thesis. An extended discussion of the sample is presented on page 31. The instrument employs optical interferometry techniques to measure the phase and magnitude of the displacement of a sample and elastic standard. With this information, the attenuation and stiffness of the sample are deduced.

Besides the sample/standard column, the apparatus consists of a waveform generator, amplifier, force transducers, He-Ne laser, photodetector, A-D converter, a computer, and other miscellaneous devices. The waveform generator produces a sinusoidal voltage that varies in frequency between .001 Hz and 100 Hz. The waveform passes through three independently adjustable amplifier channels. From there each amplified signal is sent to a piezoelectric transducer where up to 1 μm of displacement is produced. This displacement compresses the column causing deformation in the sample and standard. A Michelson interferometer detects the movement through changes in the lengths of 6 object beams relative to a reference beam. Three of the object mirrors constitute a plane above the sample and three constitute a plane between the sample and standard. The object and reference beams are interfered at the photodetector which continuously feeds the signal to the A-D converter. The A-D converter stores 16384 data points comprising several cycles of

Instrument



Sample

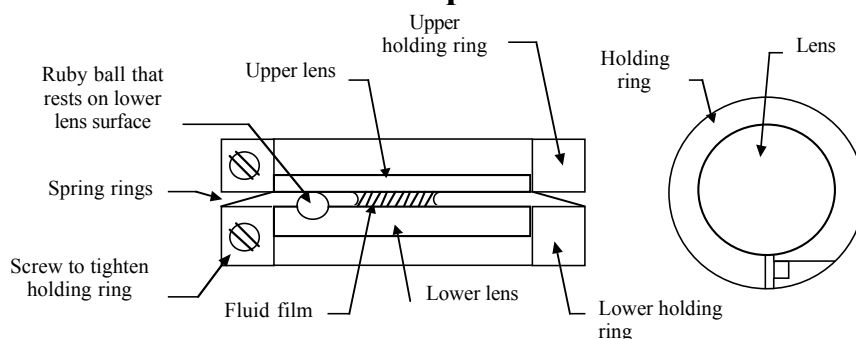


Figure I.4 Photograph and schematic of the attenuation spectrometer and sample. In the photograph, the two mirror plates can be seen along with the cylindrical sample and rectangular standard. The sample/standard column is surrounded by a rigid exoskeleton made of aluminum. Leads from a 500 volt amplifier wind up the outer structure to the top of the column to meet the transducers. The schematic shows the interferometer technique including laser, reference mirror, beam distribution assembly, and photodetector. A basic schematic of the sample includes upper and lower holding rings, upper and lower lenses, screws to tighten holding rings, spring rings, ruby ball, and fluid film.

the seismic waveform and relays it on to the computer at the end of the measurement (Appendix D). After several signals have been acquired, they are reduced to a seven parameter equation (eq. I.10, figure I.5) using a Levenberg-Marquardt nonlinear least-squares algorithm (Appendix E).

$$z(t) = A_3 \sin[A_1 \sin(2\pi ft + A_2) + A_5 + A_6 t + A_7 t^2] + A_4 \quad (\text{I.10})$$

A_1 is the sinusoidal displacement of the objective mirror. A_2 is the phase of the objective mirror relative to the waveform generator. $A_3 + A_4$ and A_4 represent the

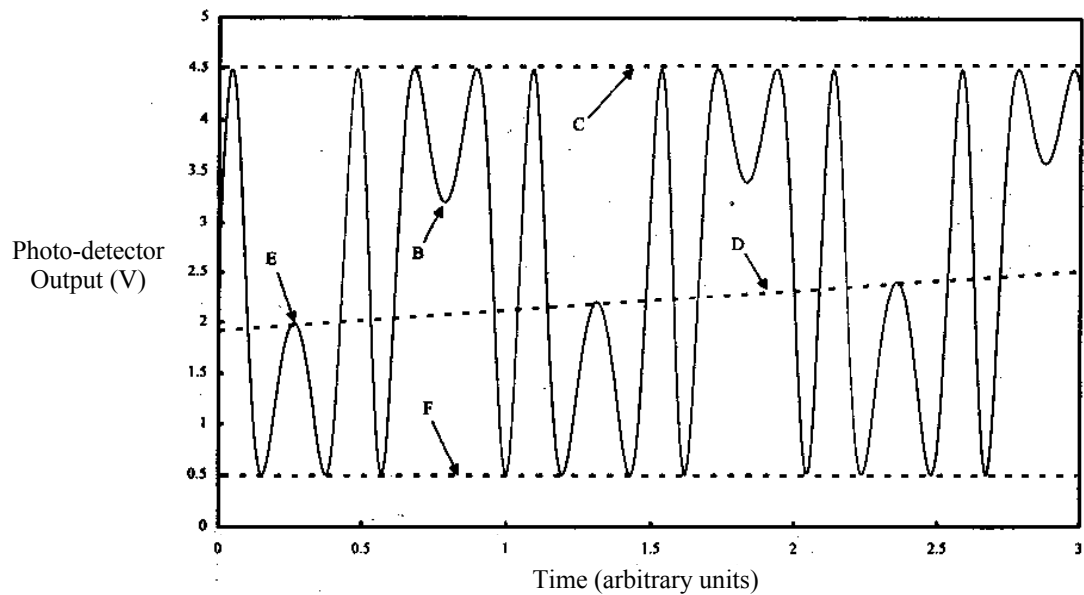


Figure I.5 Figure from Cherry et al., 1996. Interference pattern produced by combining the moving object beam and stationary reference beam. A_1 is the sinusoidal displacement of the objective mirror represented by the number of fringes between locations E and B. A_2 is the phase of the displacement. A_4 is the mean amplitude. $A_3 + A_4$ is the peak amplitude of the photodetector output corresponding to the line labeled C. $A_4 - A_3$ is the minimum output represented by the line labeled F. A_5 is the starting position of the objective mirror relative to the reference mirror measured along the beam path, labeled E. A_6 and A_7 are measurements of the linear and second order drift of the objective mirror, the line labeled D.

maximum and mean amplitude respectively of the photodetector output. A_5 is the starting position of the objective mirror relative to the reference mirror as measured along the beam path, and A_6 and A_7 are the values of the linear and second order drift of the objective mirror.

To deduce the attenuation and stiffness of the sample, the parameters A_2 and A_5 must be known. With these coefficients, the complex displacement of the sample is calculated. The vectorial difference between the upper and lower objective beams gives the displacement and phase of the sample. The phase of the sample minus the phase of the standard is a direct measure of the attenuation in the sample. The displacement of the standard relative to the displacement of the sample is the stiffness ratio between the sample and standard.

Chapter II

THEORETICAL OVERVIEW

The sample which these models describe consists of a cylindrical film of fluid between parallel plates (figure II.1). The plates are separated by an elastic spring of known stiffness. The distance between the plates, referred to as the gap height h , is approximately $100\text{ }\mu\text{m}$ and is varied sinusoidally with frequencies ω , from $.001\text{ Hz}$ to 100 Hz . The amplitude of oscillation is approximately $.02\%$ of the initial gap height and the aspect ratio of the fluid film, h/r , is approximately $.01$.

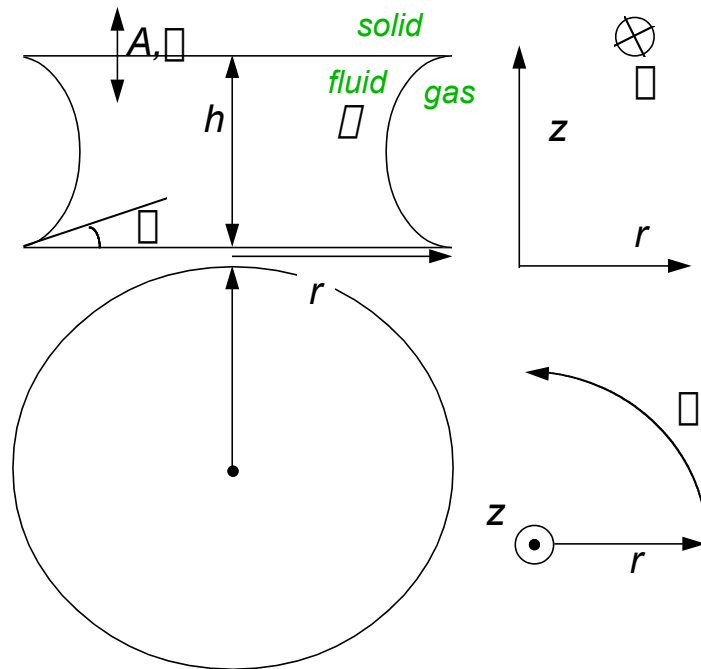


Figure II.1. Schematic illustrating the relation of key parameters in the sample. The fluid has viscosity η and sample has stiffness K . h is the distance between the plates, r is the radius of the fluid film, A is the amplitude of vibration of the upper plate and ω is the frequency of vibration of the upper plate.

Attenuation Due to Viscous Dissipation

Viscosity is a property of materials that describes how quickly a material will strain in response to an applied stress. The energy given to the material through viscous dissipation is unrecoverable and by definition leads to attenuation.

In order to calculate the energy absorbed by viscous dissipation, the strain rate or, equivalently, the velocity field within the material and its viscosity must be known. The velocity field can be described by the Navier-Stokes equations and the viscosity has been empirically measured for many materials.

Energy and Force

One path for the transfer of energy is to impart a force on an object over a distance. Quantitatively this is given by the expression

$$E = \int \hat{F} \cdot d\hat{l} . \quad (\text{II.1})$$

The total force, F , acting externally on a body is the normal component of stress, $\hat{\sigma} \cdot \hat{n}$, acting on that body integrated over the body's entire surface area, A . Gauss's theorem states that the integral over the surface of a body is related to the integral over its volume through the divergence of the integrals argument. The force field can then be expressed as the total change in stress as a function of space within the volume, V ,

$$\hat{F} = \int \tau \cdot \hat{\gamma} dA = \int \tau \cdot \hat{\gamma} dV. \quad (\text{II.2})$$

The stress relates to the strain rate as

$$\tau = \eta \hat{\gamma} \quad (\text{II.3})$$

where η is the materials Newtonian viscosity and $\hat{\gamma}$ is the strain rate . The strain rate is expressed as the gradient of the velocity field, $\nabla \hat{U}$. In terms of the velocity and a constant viscosity, the viscous force becomes

$$\hat{F} = \eta \nabla^2 \hat{U} dV. \quad (\text{II.4})$$

Velocity Field

The Navier-Stokes equation (Gerhart and Gross, 1985), a statement of cause and effect, is the governing principle that allows for the solution of the velocity field within continuous bodies. It simply states that the time rate of change in linear momentum of a body is equal to the sum of forces acting on that body. This relationship is customarily expressed in terms of the density ρ , velocity U , pressure gradients ∇p , and viscous forces $\eta \nabla^2 \hat{U}$,

$$\frac{1}{V} \frac{d(m\hat{U})}{dt} = \rho \frac{d\hat{U}}{dt} = \rho \rho p + \rho \rho^2 \hat{U}. \quad (\text{II.5})$$

The general solution to this equation does not exist due to the equations complexity. In order to extract the velocity field, the system must be known well enough to understand where simplifications can be made.

Some of the terms in the Navier-Stokes equation can be neglected due to their relatively small magnitude. The expression on the left hand side of the Navier-Stokes equation can be broken up into the force per unit volume needed to accelerate the body and a force per unit volume associated with the inertia of the body,

$$\rho \frac{\partial \hat{U}}{\partial t} + \rho \hat{U} \cdot \nabla \hat{U}, \quad (\text{II.6})$$

respectively. The velocity of the upper plate is on the order of $2\rho A\rho$ and dt on the order of the period, $1/\rho$. The Spatial derivatives, dr and dz , are on the order of the gap dimensions, r and h . The expressions giving rise to the order of magnitude of the accelerating, inertial, and viscous terms are

$$\rho 2\rho A\rho^2, \rho \frac{4\rho^2 A^2 \rho^2}{h}, \rho \frac{2\rho A\rho}{h^2}. \quad (\text{II.7})$$

For the values of the parameters used in this model, $\rho = 1000 \text{ kg/m}^3$, $\rho = .001 \text{ kg/ms}$, $A = .2 \text{ m}$, $\rho = 10 \text{ Hz}$, and $h = 100 \text{ m}$, the size of these equations are .0126, .0016,

and $.126 \text{ N/m}^3$, respectively. At 10 Hz, an error of less than 1% is incurred in the Navier-Stokes equation by neglecting the inertial effects of the fluid and an error of 10% develops by neglecting the acceleration of the fluid. These errors decrease quickly with decreasing frequency. In terms of the velocity profile, which is the parameter of interest, there may be no error for all frequencies by neglecting equation II.6, the acceleration and inertia of the fluid film. The simplified Navier-Stokes equation is the Stokes equation,

$$\nabla^2 p = \mu \nabla^2 \hat{U}. \quad (\text{II.8})$$

The current form of this force balance is still cumbersome and can be reduced further by evaluating the Laplacian.

The Laplacian of a vector, \hat{U} , in cylindrical coordinates is

$$\begin{aligned} \nabla^2 \hat{U} = & \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) (\hat{U}) + \frac{1}{r} \frac{\partial}{\partial z} \left(r \frac{\partial}{\partial z} \right) (\hat{U}) - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r \hat{U}_r) \right) - \frac{\partial}{\partial z} \left(\frac{1}{r} \frac{\partial}{\partial z} (r \hat{U}_z) \right) \\ & + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) (\hat{U}) + \frac{1}{r} \frac{\partial}{\partial z} \left(r \frac{\partial}{\partial z} \right) (\hat{U}) - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r \hat{U}_r) \right) - \frac{\partial}{\partial z} \left(\frac{1}{r} \frac{\partial}{\partial z} (r \hat{U}_z) \right) \\ & + \frac{\partial}{\partial z} \left(\hat{U} \right) + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) (\hat{U}) - \frac{\partial}{\partial z} \left(\frac{1}{r} \frac{\partial}{\partial z} (r \hat{U}_z) \right) - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r \hat{U}_r) \right) \end{aligned}$$

Due to the high modulus of water and the lack of confining pressure, the water volume can be considered incompressible, thus $\nabla \cdot \hat{U} = 0$. Because of the symmetry of the fluid volume and of the forces acting on it, angular dependence is absent. This condition requires all terms containing derivatives with respect to the angle ϕ to be zero. Further simplification arises when the assumption is made that the velocity in

the z direction does not depend on the radius, i.e. $\frac{\partial U_z}{\partial r} = 0$. This assumption is made without proof but is believed to be true for parallel plates. The resulting Stokes equations are

$$\frac{\partial p}{\partial r} = \mu \frac{\partial^2 U_r}{\partial z^2} \quad (\text{II.9})$$

and

$$\frac{\partial p}{\partial z} = \frac{\mu}{r} \frac{\partial}{\partial r} \left(r \frac{\partial U_r}{\partial z} \right). \quad (\text{II.10})$$

After taking the derivative of the first with respect to z and the second with respect to r , the two equations can be set equal to each other.

$$\mu \frac{\partial^3 U_r}{\partial z^3} = \mu \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial U_r}{\partial z} \right) \right) \quad (\text{II.11})$$

The assumption of the water volume being incompressible allows U_r to be written in terms of U_z . Starting with the divergence of the velocity vector,

$$\nabla \cdot \hat{U} = \frac{1}{r} \frac{\partial}{\partial r} (r U_r) + \frac{\partial U_z}{\partial z} = 0, \quad (\text{II.12})$$

U_r is found to be

$$U_r = -\frac{r}{2} \frac{\partial U_z}{\partial z}. \quad (\text{II.13})$$

Substituting U_r in terms of U_z into equation II.11 remembering that the z component of velocity does not vary in the radial direction, namely $\frac{\partial U_z}{\partial r} = 0$, equation II.11 becomes

$$\frac{\partial^4 U_z}{\partial z^4} = 0 \quad (\text{II.14})$$

where the general solution is

$$U_z = a \frac{z^3}{6} + b \frac{z^2}{2} + cz + d. \quad (\text{II.15})$$

It is now time to invoke the boundary conditions. At the lower solid boundary where $z = 0$, $U_z = 0$ and $U_r = 0$, a no slip boundary condition. At the upper solid boundary where $z = h$, $U_z = U_h = A$ and $U_r = 0$, also a no slip boundary condition. Using equation II.13 and applying these boundary conditions reveals U_r to be

$$U_r = 6A \frac{r}{h} \frac{z^2}{h^2} - \frac{z}{h}. \quad (\text{II.16})$$

The velocity profile for various r is shown in figure II.2. At $z = 0$ and $z = h$, the velocity is zero. At $h/2$, the velocity is at a maximum and as r increases, so does the velocity.

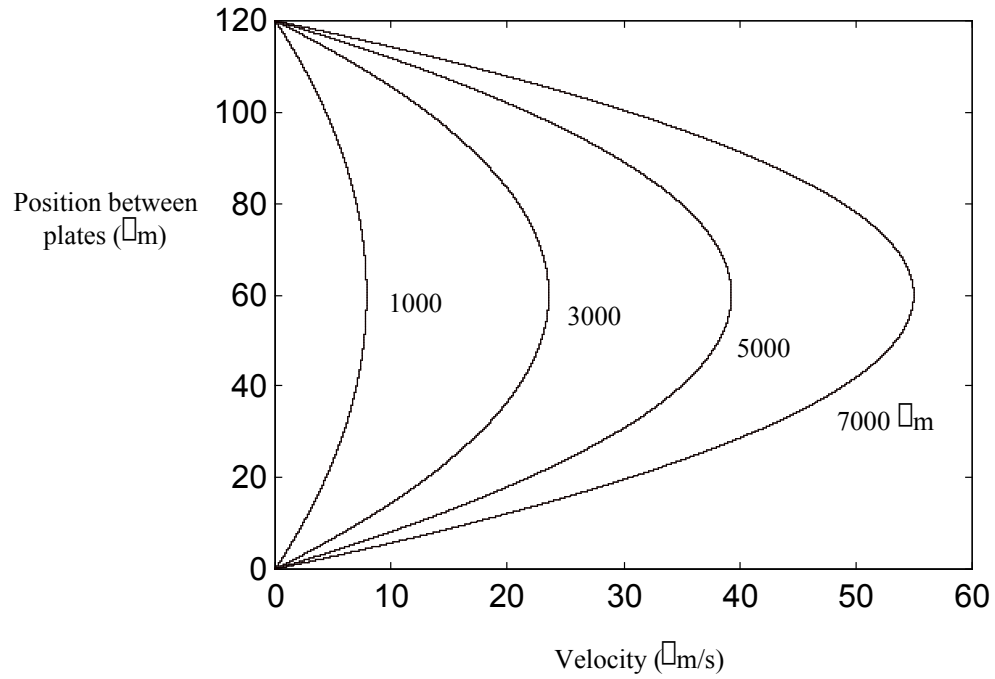


Figure II.2 Velocity field dependence on position between plates for various distances from center of fluid film. Curves are calculated using $A = .2 \mu\text{m}$, $h = 120 \mu\text{m}$, $\omega = 1 \text{ Hz}$, and $r = 1000, 3000, 5000, 7000 \mu\text{m}$.

Dissipation

All of the ingredients for determining the energy lost through viscous dissipation are present. Applying the Laplacian to the radial component of the velocity (eq. II.16), integrating over the volume and multiplying by the viscosity yields the restrictive force acting in the radial direction (eq. II.4). Integrating this force over the radius of the fluid volume gives the energy lost due to viscous dissipation (eq. II.1, II.17)

$$E = 2\eta^2 A \frac{r^4}{h^2} \quad (\text{II.17})$$

The key to this effect lies within two fundamental concepts, viscosity and the Laplacian of the velocity field. As discussed earlier, viscosity is a property of materials that governs the rate at which the material will strain in response to an applied stress. The Laplacian of a vector is a measure of how the individual components of that vector vary with respect to each of the coordinate directions. For the geometry discussed in this paper, the magnitude of the resistive force depends on

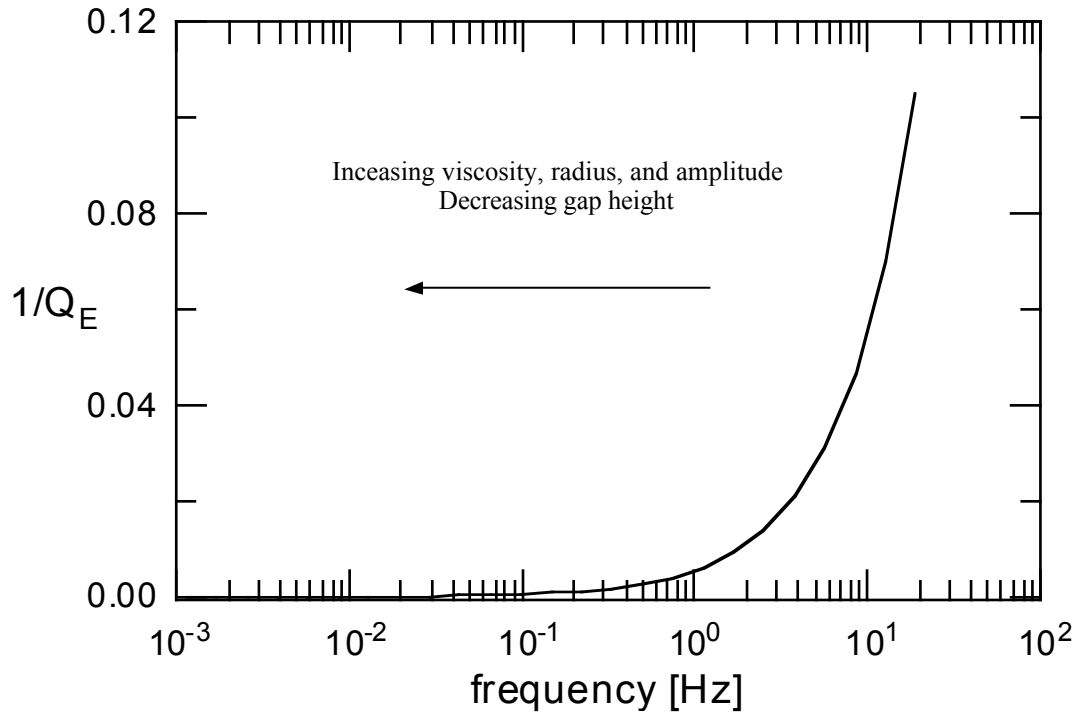


Figure II.3 - Attenuation versus Frequency. Increasing viscosity, radius, and amplitude, or decreased gap height result in increased shear and would move the curve to the left. This causes greater energy lost and higher attenuation for a given frequency. Input parameters are $h = 120 \mu\text{m}$, $r = 6\text{mm}$, $A = .2 \mu\text{m}$ and $\eta = .001 \text{ kg/ms}$. The curve does not account for inertial effects or the force needed to accelerate the fluid.

how quickly the velocity in the radial direction changes with respect to z . The dissipated energy is a result of the shearing of the fluid volume. A faster change in velocity produces greater shear and increased attenuation (figure II.3).

Attenuation due to Contact Line Movement

As was inferred in the introduction, the mobility of the contact line can be represented as the probability of water molecules at the contact line to bond to the solid surface. The probability for a new bond to be created depends on the departure from equilibrium of the forces on the contact line. As pressure is built up within the fluid due to a decrease in volume of the pore space, the meniscus changes shape bringing molecules on its surface closer to the solid surface. This is represented by an increase in contact angle (figure II.1). Given enough time, new bonds will form and the contact line will move effortlessly. But if the pressure within the fluid builds too quickly, the molecules at the gas fluid interface will not have enough opportunity to bond to the solid surface and the sample will stiffen. Let this regime of contact line motion be considered a sticking regime and described with a probability of the form

$$P_i = \frac{(\theta_i - \theta_e)^b}{\theta_h^b} \quad (\text{II.18})$$

where θ_i is the contact angle at the i th interval of time, θ_e is the equilibrium contact angle, θ_h is the contact angle hysteresis and b is a constant that controls how quickly

the probability changes during sticking. Slipping occurs above θ_h where the contact line motion is governed by other processes.

Energy is lost when some change in fluid pressure balances a related change in fluid radius. To determine the energy lost due to restricted contact line motion, the resistive force on the contact line and the distance over which the contact line moves must be known. Because of the expression for the volume of the fluid film, the force on the contact line is not known analytically as a function of distance. The integral in equation II.1 must be made a sum and the force and distance calculated for each interval of time.

$$E_{lost} = \sum_i F_i dr_i \quad (II.19)$$

The force on the contact line is due to a change in capillary pressure (Corey, 1994)

$$\Delta P_c = \frac{4\sigma}{h_i} (\cos\theta_i - \cos\theta_e) \quad (II.20)$$

where σ is the surface tension of the fluid. The force is given by

$$F_i = \sigma r_i h_i \Delta P_c = 4\sigma r_i (\cos\theta_i - \cos\theta_e) \quad (II.21)$$

The calculation of the change in radius is broken up into several stages. At each time step, dt , the velocity of the upper plate, $\partial h / \partial t$, is calculated from the

frequency, ω , and amplitude, A , of the passing waveform to determine the new height.

For infinitesimal dt , h is

$$\begin{aligned} h_i &= h_0 + \frac{\partial h}{\partial t_i} dt \\ \frac{\partial h}{\partial t_i} &= 2\omega A \cos(2\omega t_i) \end{aligned} \quad (\text{II.22})$$

The velocity is used to determine the change in separation between the solid surfaces. For an incompressible fluid and stationary contact line, the contact angle must change. The volume for a cylindrical film of fluid with a semicircular meniscus between parallel surfaces is approximately the volume of a cylinder, $\pi r^2 h$, minus the fractional volume of the surrounding annulus (Jeff Boyd, personal communication),

$$V = \pi r_i^2 h_i - \frac{\pi h_i^2 r_i}{2} \frac{1}{\cos^2 \theta_i} + \tan \theta_i. \quad (\text{II.23})$$

Because the contact angle can not be solved for analytically, an iterative procedure is used to converge on the correct value. The resulting change in contact angle produces some probability for a bond to be created at the contact line. If the probability is less than 1 and greater than a randomly generated number between 0 and 1, the contact line advances an appropriate distance, dr_m , the distance considered

to be the maximum change in fluid radius when sticking. Averaged over the entire contact line, the change in fluid radius is

$$dr_i = p_i dr_m \quad (\theta < \theta_h) \quad (\text{II.24})$$

If the probability is greater than 1, the contact line motion is determined by alternate processes. The behavior has been studied by Moerig et. al. (1996). They found that the velocity of the contact line has an approximately linear dependence on contact angle. The fluid radius may then be given by

$$dr_i = m \frac{dr_m}{\theta_h} [\theta_i - (\theta_e + \theta_h)] + dr_m \quad (\theta > \theta_h) \quad (\text{II.25})$$

where m scales the slope of the line, dr_m/θ_h . The resulting velocity versus contact angle relationship is seen in figure II.4. The lost energy (eq II.19) is solved using either equation II.25 or II.24 and II.21.

The time steps progress until a minimum in separation exists between the solid surfaces, 1/4 cycle. After the last time step, the change in stiffness of the sample and energy stored in the fluid film are calculated. The added stiffness is simply the change in capillary pressure (eq. II.20) multiplied by the area of the fluid-solid interface, θr^2 , divided by the amplitude of motion of the upper surface,

$$\partial k = \frac{\theta r_i^2 \Delta P_c}{A} = \frac{4 \theta r_i^2 \Delta}{A h_i} (\cos \theta_i - \cos \theta_e). \quad (\text{II.25})$$

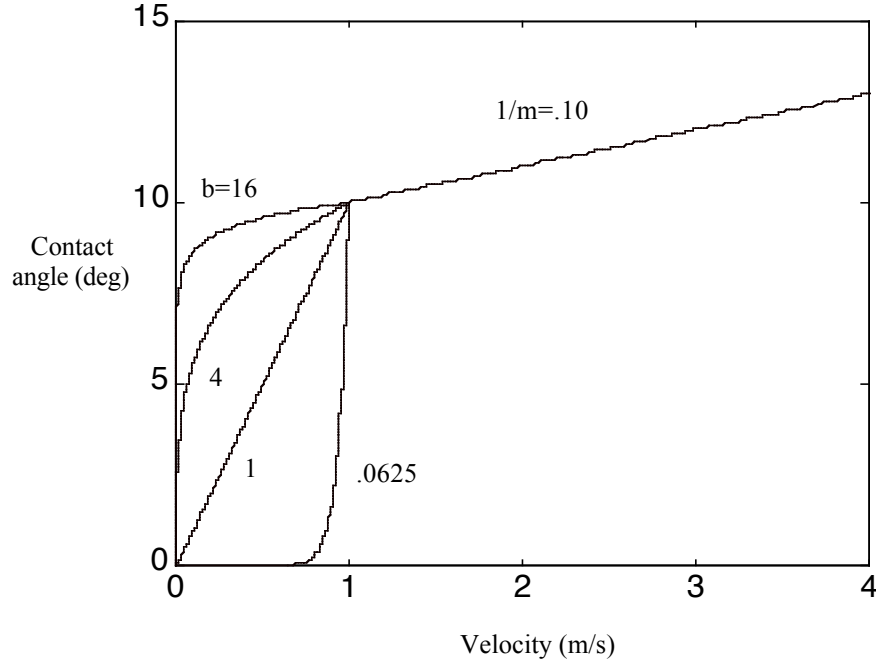


Figure II.4 Contact angle versus velocity; $\partial r_m / \partial t = 1$ and $\theta_h = 10^\circ$. Values used to calculate fits are given in table V.1.

The energy stored is the energy that would be given back by the system if all external stresses were relaxed. The energy is approximately equal to the change in capillary pressure times the area of the fluid solid interface times the distance over which the capillary force acts $A/2$, plus the energy stored in the spring rings of the sample,

$$E_{stor} = A r_i^2 P_c + \frac{1}{2} k A^2. \quad (\text{II.26})$$

By having multiple regimes of contact line movement, multiple shapes in attenuation and stiffness result. In figure II.5, all parameters are held fixed while the constant b is varied. Recall that b is the power to which the ratio θ/θ_h is raised. As b increases, the peak in attenuation and dispersion in stiffness below 1 Hz. move to lower frequencies. If the slope of the contact line velocity versus contact angle

relationship decreases then the peak in attenuation and dispersion in stiffness shift to lower frequencies. With increasing equilibrium contact angle the magnitude of both

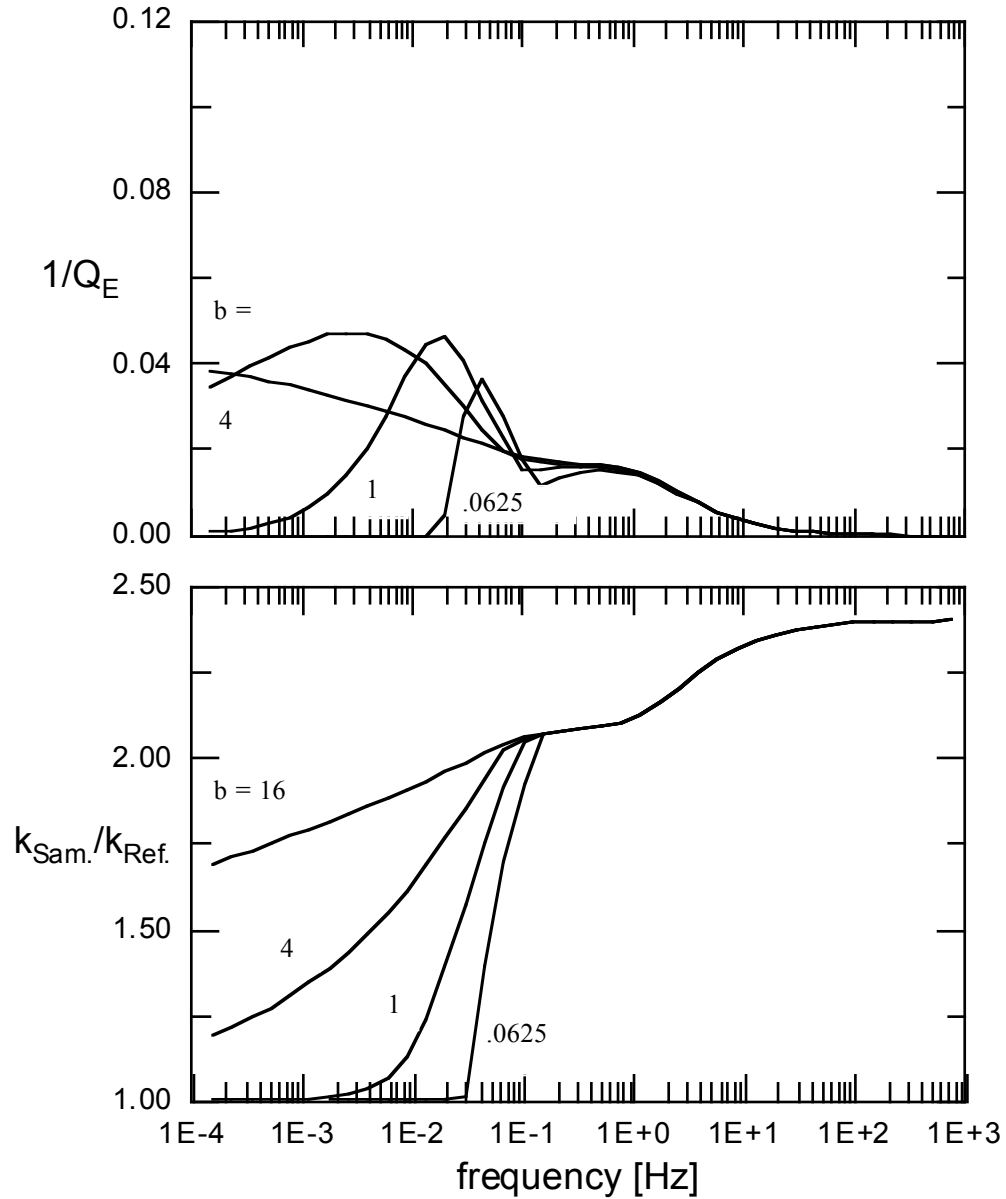


Figure II.5 Attenuation and stiffness versus frequency in dependence on the constant b . The constant b controls the behavior of sticking during low velocities or low frequencies. As b increases from .0625 to 16, the attenuation peak and stiffness dispersion move to lower frequency. This is due to the fact that the slope of the contact line velocity versus contact angle relationship is decreasing. Values for b used in the modeling of attenuation and stiffness versus frequency span this range.

the attenuation and stiffness increase. This is due to the fact that to displace the same volume of fluid during deformation, a higher equilibrium contact angle needs a greater contact angle change. With increasing hysteresis, the attenuation due to slipping decreases. This is simply because more fluid can be taken up in deformation of the meniscus rather than moving the contact line. This also means that the stiffness increases due to the increased contact angle. Above .1 Hz the attenuation and stiffness are independent of b . This means that at any given frequency or contact line velocity, the measured attenuation and stiffness do not depend on the shape of the contact line velocity versus contact angle relationship below that point.

Chapter III

EXPERIMENTAL PROCEDURE

Sample

The sample consists of nearly parallel glass surfaces separated by a thin cylindrical film of fluid (figure III.1) and elastic spring.

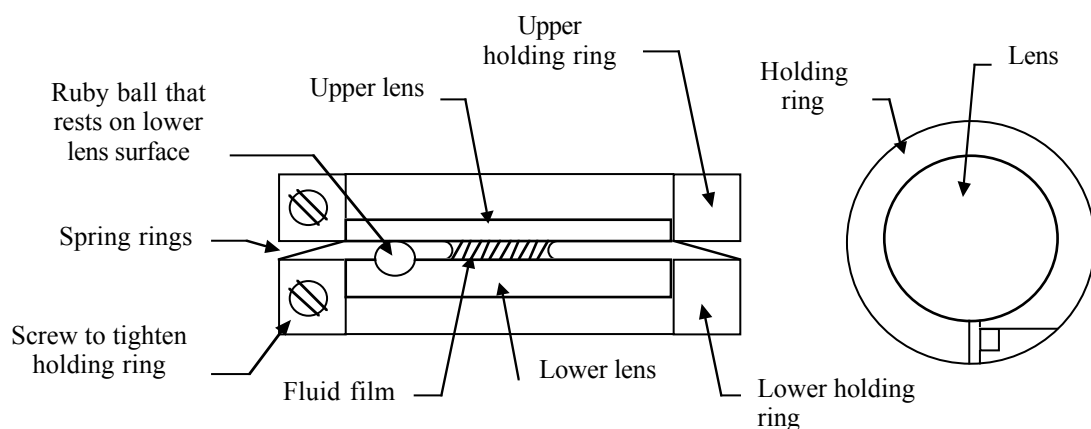


Figure III.1 Basic schematic of the sample design. Included are upper and lower holding rings, upper and lower lenses, screws to tighten holding rings, spring rings, ruby ball, and fluid film.

A small ruby ball rests on the lower surface acting both as a spacer between the solid surfaces and as an indicator for the gap height, a crucial parameter in the models considered for this sample. The glass surfaces are held rigidly by holding rings. These rings are separated by an elastic spring mechanism consisting of donut like metal disks. Although elastic, the mode of deformation of the metal disks is such that

a stress dependent stiffness emerges (Appendix B). This produces problems in reproducing the same dry stiffness between measurements.

Preparation

The basic procedure for preparing the sample is the following: 1. The glass lenses are cleaned. 2. The lenses are oriented in the holding rings. 3. The ruby ball is placed on the bottom lens. 4. The sample is partially fluid saturated and placed in a plastic bag to reduce evaporation.

To clean the lenses, they are placed in an oven at 420 degrees C for two hours. This cleaning procedure produces a condition of the surfaces which is relatively reproducible and referred to as unaltered. The condition is reflected in a relatively mobile wetting front. The sample surfaces may then be contaminated. Before any data are taken, all of the instruments are turned on and allowed to warm up. This reduces the chance of fluctuations in the movement of the upper mirror plate and wavelength of the laser light. The block upon which the sample and column sit is checked to see that the only differential forces acting on it are due to the piezoelectric transducers.

During sample assembly, the critical task is orienting the upper lens such that during all measurements, the sample is in the same prestressed state for a given surface separation. This will insure the same dry sample stiffness between measurements. The upper surface of the lower lens is made flush with the top of the bottom holding ring. On this the spring rings are placed. The top holding ring with its lens loosely inside is then added. The upper lens slides down and comes in contact

with the lower lens. A Plexiglas container holding up to 2 liters of water is placed on the upper holding ring. In the center of the container is a dry cylinder 8 cm in diameter with a 3 cm diameter hole at its center through which the upper surface of the upper lens is exposed. With the desired mass of water in the container, ~ 2 kg, the exposed surface of the upper lens is tapped by a small Allen wrench. Newton rings are observed at the contact between the two inner glass surfaces. When the Newton rings are centered, the upper holding ring is tightened, thereby orienting the upper lens. The actual point of contact is up to 1 Newton of reduced force. The sample may then be saturated and placed in a plastic bag.

When the instruments are ready, the sample is incorporated into the column. The reference, lower mirror plate and upper mirror plate are aligned with a mechanism fixed to the side of the column. The sample is aligned manually. Each component of the column is positioned such that its mass is distributed symmetrically about a single axis, \hat{a} , parallel to gravity. The upper mirror plate is allowed to rest freely on the sample and reference producing an initial prestress. Further stress is added by the transducer screws. The column is compressed until the upper surface within the sample comes into contact with the ruby ball on the lower surface. For partially saturated measurements, the fluid film must be centered on the \hat{a} axis in order to determine the separation between upper and lower surfaces. The transducer screws are then backed off slightly so that there is no contact between the ruby ball and upper lens.

All six mirrors are aligned and the gains of the amplifier channels adjusted until the amplitudes of the upper mirrors are between 2.8 and 3.0 fringes of laser

light, approximately 1 μm . One fringe corresponds to one wavelength of laser light, which for a He-Ne laser is $\sim 600\text{ nm}$. Three fringes is 1800 nm. The displacement of the upper mirror plate is 1/2 the change in distance of the sample beam, $\sim 1\text{ }\mu\text{m}$. Two Styrofoam boxes are then placed around the column to reduce air flow.

Dry Sample

Before reliable measurements can be made on a partially saturated sample, the dry system must be understood. The two elements of the dry sample are the spring rings and the plastic bag. First, the sample is constructed such that it lacks only fluid. This is done to insure that the basic sample design will not effect the results obtained later. Frequency independent and low magnitude attenuation along with a constant stiffness should be observed. According to the process described above, measurements are made on this sample multiple times with slight variations of the sample orientation. Initially nothing changes, only time. Then the column is disassembled and reassembled and the measurements are repeated.

The next step involves the removal of the upper holding ring, the process involved when saturating the sample. As before, the sample is placed in a plastic bag and measurements are taken. The sample is removed from the plastic bag and the upper holding ring set aside. The sample is reassembled, placed back in a plastic bag, and measurements are again performed. Like the basic sample design, the removal and reassembly of the sample should not give rise to attenuation or changing stiffness that would interfere with the principle investigation.

Partially Saturated Unaltered Sample

The sample surfaces are cleaned (see preparation p. 32) before the sample is partially saturated with water. At this point, the saturant is seen to easily flow along the sample surfaces. Because of the unaltered condition of the sample surfaces, great care must be taken when placing the sample within the column so that the fluid does not flow into the spring rings. The sample and column are assembled as described above. A measurement is taken and repeated until all of the fluid has evaporated. A dry measurement follows to determine the dry sample stiffness. Before and after the measurements, the saturation state of the sample is observed to give an estimate of the fluid radius. This process is repeated for several clean samples.

Partially Saturated Surface Affected Sample

After a good understanding of the unaltered surface measurements has been obtained, the sample surfaces are contaminated with Propanol. The contamination is incurred by wetting the sample surfaces and allowing the contaminant to fully evaporate. This is done twice to insure the conversion of the surface from hydrophilic to hydrophobic. The sample and column are assembled and measurements are taken and repeated until the fluid has evaporated. This process is repeated once more for comparison.

Chapter IV

RESULTS

The following results were obtained using the sample discussed in the previous chapter. All samples were measured between .002 and 20 Hz. For those results that do not encompass this frequency range, the data has been discarded because the phase and amplitude values could not be extrapolated or the A/D converter stopped functioning during the measurement. The amplitude for all measurements can be found by solving the equation

$$A = \frac{x_A}{\frac{2k_2}{k_1} + 1} \quad (\text{IV.1})$$

where x_A is the peak to peak displacement at the upper mirror plate ($\sim .9$ μm), k_1 is the sample stiffness, and k_2 is the reference stiffness (~ 60 kN/m).

Dry

The dry system was measured on three different samples. For each sample, a new bag was used and measurements were made multiple times. Initially the sample was left untouched and the measurements were repeated. Without taking the sample

out of the plastic bag, the sample column was disassembled and reassembled. All measurements show roughly the same behavior (figure IV.1). The attenuation averages around zero, has a standard deviation of $\sim .004$ rad, and increases slightly with increasing frequency. The stiffness has a virtually constant value up to 10 Hz. Above this frequency, the data exhibit a repeatable rise.

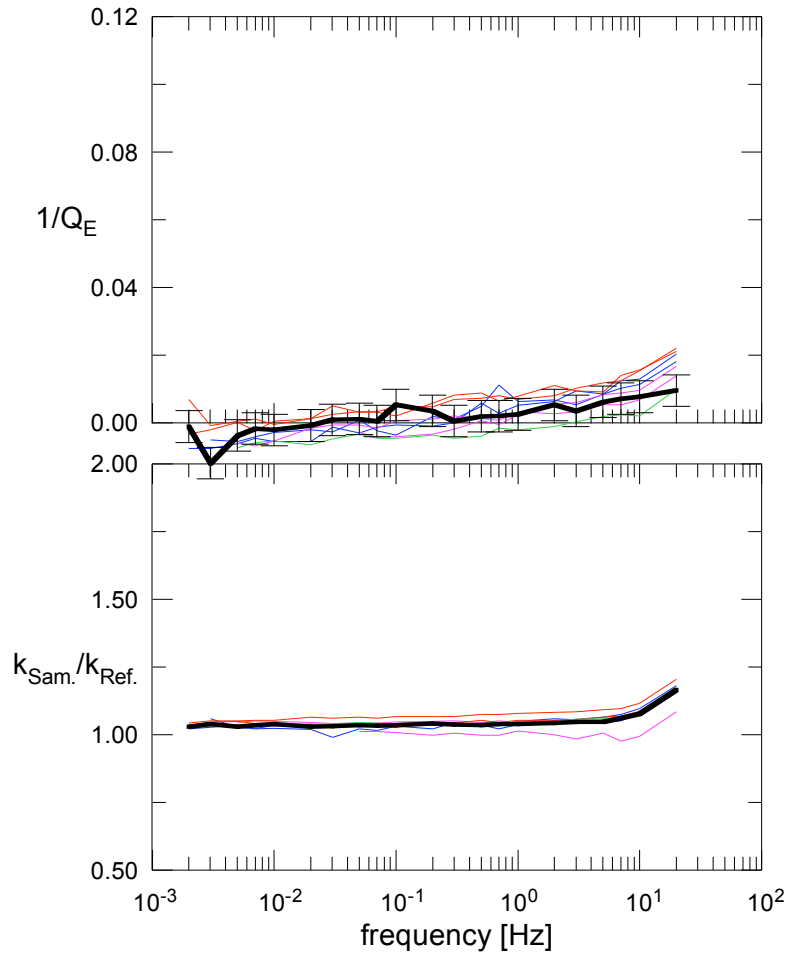


Figure IV.1 Measured data on dry cylindrical samples. The upper graph shows attenuation vs. frequency. The error bars on the thick curve are an averaged standard deviation over three of the measurements. The lower graph contains measurements of the sample/reference stiffness ratio.

Saturated

For all saturated cases presented, the saturant was water. Water has a viscosity of $.001 \text{ kg/m s}$ and a surface tension of $.073 \text{ N/m}$ at 20°C . The fluid volume injected was $\sim 2 \times 10^{-8} \text{ m}^3$. The fluid radius ranged from ~ 4 to 6 mm depending on the amount of evaporation between measurement and injection. The amplitude of deformation is given by equation IV.1. The following values are unknown: equilibrium contact angle, contact angle hysteresis, velocity of the contact line at the stick-slip transition, dr_m/dt , slope of the contact line velocity versus contact angle above θ_h , and dry sample stiffness. The range of the contact line-contact angle parameters is currently being measured. Unfortunately the samples dry stiffness was not obtained due to a precipitate that had formed after evaporation of the fluid (figure IV.2).

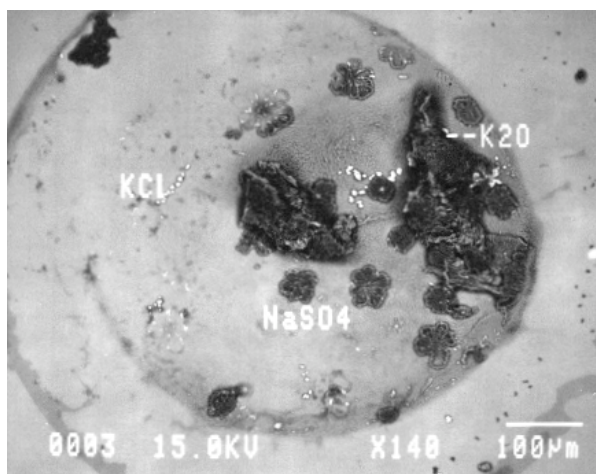


Figure IV.2 Precipitates that had formed at the end of all fluid saturated measurements. The small white speckles are potassium chloride. The flower like disk structures are sodium sulfate, and the larger structures are potassium oxide. This last precipitate was responsible for preventing the acquisition of a dry sample stiffness. There was enough of the compound to span the gap separation.. Photo and analysis courtesy of John Drexler, University of Colorado at Boulder.

Clean

During injection of the fluid, the contact line moves smoothly. There is little resistance. A slight tilt of the sample is enough to cause the fluid to flow quickly. As stated before, great care must be taken to keep the fluid from flowing into the spring rings. Looking down the axis of axial symmetry of the sample, the fluid film is nearly a perfect circle. This was the basic theme behind all three sets of partially saturated unaltered samples.

The first sample shows measurable attenuation in the entire frequency range, $\sim .04$ rad (figure IV.3). Generally, the attenuation increases slightly with increasing frequency. Above 1 Hz, the attenuation seems to increase more rapidly. At the mid to low frequencies the attenuation shows little frequency dependence. The corresponding stiffness exhibits a gradual increase from $1.25 k_{\text{ref}}$ at .002 Hz to $1.5 k_{\text{ref}}$ at 10 Hz.

The second sample looks much like the first (figure IV.4). Because more of the high frequency data was measured, a rapid increase in attenuation with increasing frequency is very apparent. The sample displays similar frequency independent attenuation between .01 and 1 Hz. What has not been seen before is the sudden decrease in attenuation below .01 Hz.

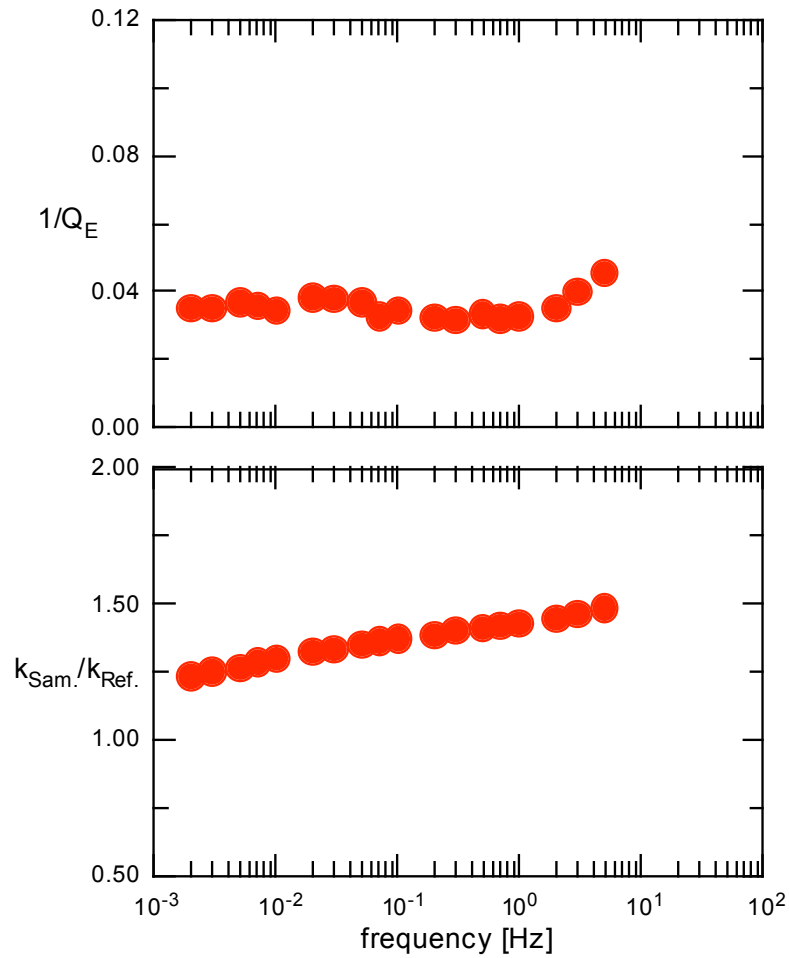


Figure IV.3 First clean sample; attenuation and stiffness versus frequency. The attenuation and stiffness show a slight increase with increasing frequency. A more rapid increase occurs in the attenuation above 1 Hz. Each solid circle represents a data point. For the attenuation curve, the vertical size of the circle is on the order of the standard deviation as derived from the dry experiments.

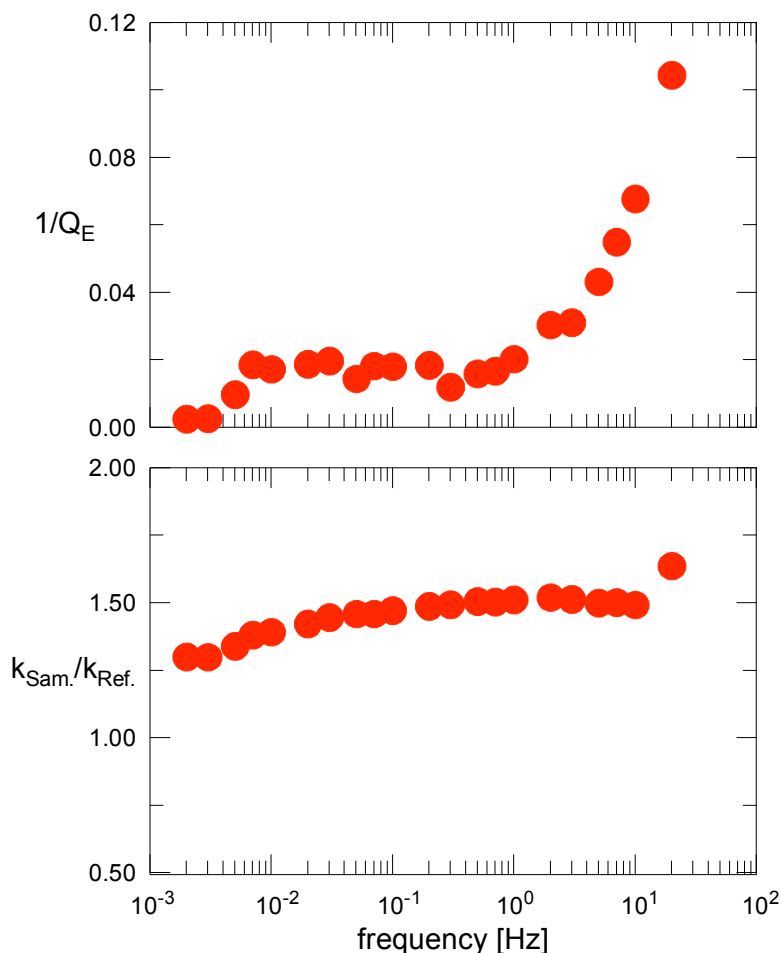


Figure IV.4 Second clean sample; attenuation and stiffness versus frequency. Attenuation and stiffness continue to increase with increasing frequency. As compared to the previous sample, a more dramatic rise in attenuation occurs above 1 Hz. Below .01 Hz, a sudden decrease is noted.

The third sample was not as fortunate as the first two due to a failure of the A/D converter to finish much of the first measurement. A second measurement was completed but at a lower saturation (figure IV.5). The fluid radius was approximately 4 mm with all other parameters as before. The resulting measurement closely resembles the data taken on the second sample. The major difference is lower attenuation and stiffness at all frequencies.

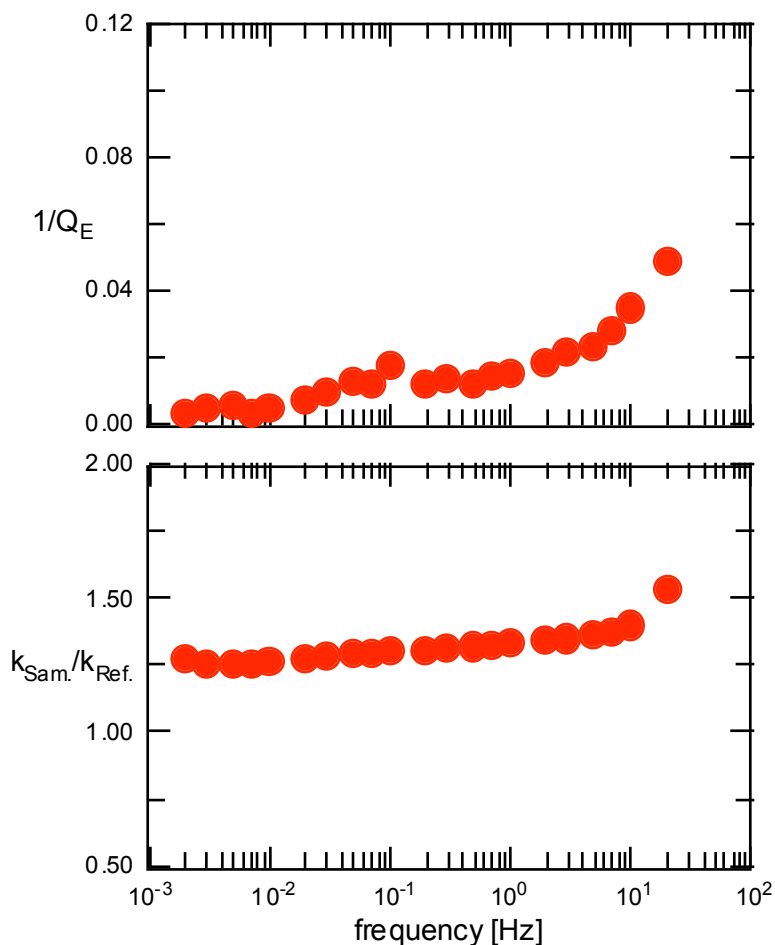


Figure IV.5 Third clean sample; attenuation and stiffness versus frequency for a fluid radius of 4 mm. The frequency dependence resembles that of the second sample with the exception of a nearly constant decrease at all frequencies.

Contaminated

After exposing the sample surfaces to either Propanol or acetone, new sets of attenuation and stiffness measurements were made. The surface alteration produces an observable effect on fluid flow. When saturating the sample, the contact line moves irregularly. A slight tilt of the sample now has visually no observable effect.

Looking down the axis of axial symmetry of the sample, the fluid film takes on a very irregular, oblong shape. The sample is stressed multiple times to force the fluid into a more circular form and more closely match the model constraints, i.e. cylindrical symmetry.

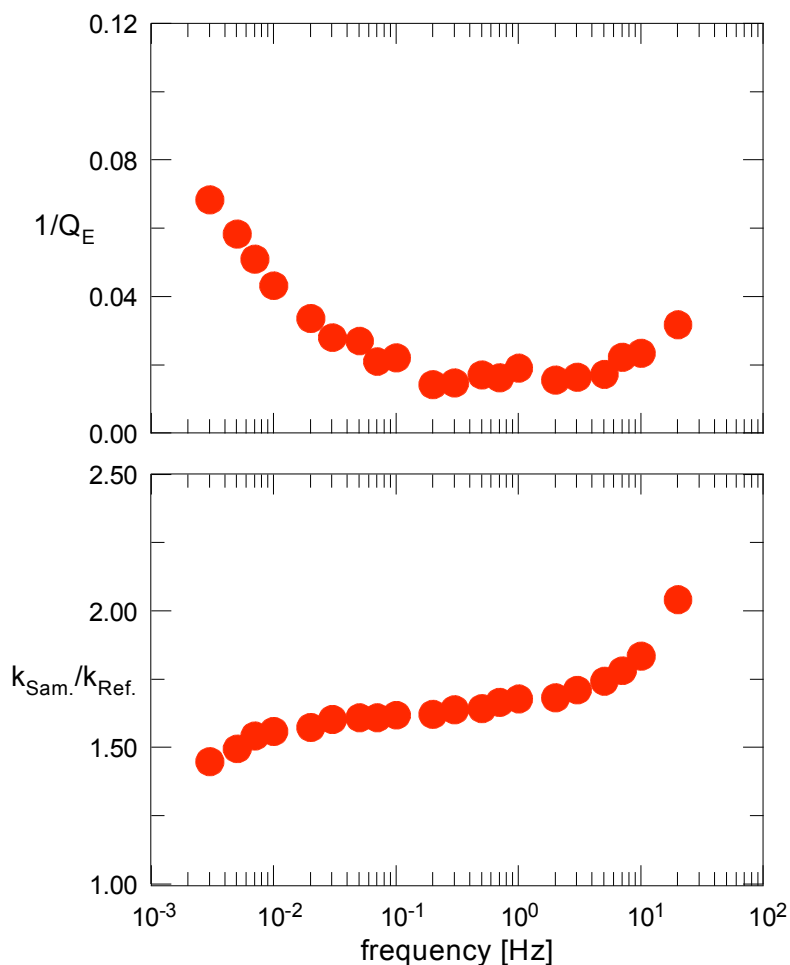


Figure IV.6 First Propanol contaminated sample; attenuation and stiffness versus frequency. The sample shows a very different frequency dependence in comparison to the clean measurements. Note the change in stiffness axis from .5 to 2.0 in the clean figures to 1 to 2.5 in the contaminated figures.

Two contaminated measurements were made, both on surfaces exposed to Propanol (figure IV.6 and figure IV.7). The Propanol measurements look very

similar to each other but remarkably different from the clean samples. In addition to the high frequency attenuation, they both show a large increase in attenuation at low frequencies. Accompanying this new attenuation behavior is an increased stiffness.

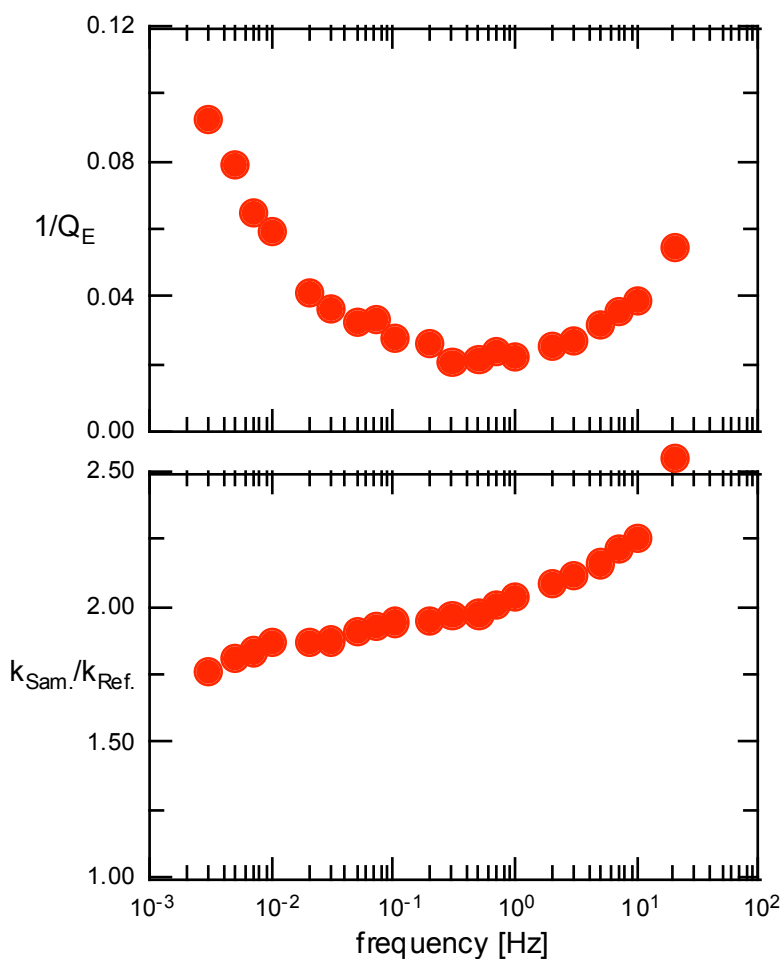


Figure IV.7 Second Propanol contaminated sample; attenuation and stiffness versus frequency. The sample looks almost identical to the other Propanol contaminated sample. The only difference is a slight increase in stiffness and attenuation over the entire frequency range.

Chapter V

INTERPRETATION

In the modeling of the attenuation and stiffness for the three clean samples and first Propanol contaminated sample, there are several characteristics that need to be kept in mind. All samples show some amount of the following trends except where noted: 1. Attenuation increases with increasing frequency above 1 Hz. 2. Attenuation is nearly frequency independent between .01 and 1 Hz. 3. For the clean samples, the attenuation decreases at the low frequencies. 4. For the contaminated sample, the attenuation increases below .1 Hz. 5. Stiffness increases with increasing frequency throughout the entire frequency range. 6. The contaminated sample is markedly stiffer than clean samples.

All calculations assume a gap height of 123 μm , displacement amplitude of .19 μm , and dry sample stiffness 1.05 times that of the reference. A water volume of .02 cm^3 was injected into the samples. For a cylindrical fluid film, this volume and gap height, a radius of 6 mm results but due to the curvature of the upper lens, a radius of 5 mm is used. The true value is unknown. A radius of 4 mm is applied to the third clean measurement because approximately 30% of its volume had evaporated. The values for the equilibrium contact angle, contact angle hysteresis, velocity of the contact line at the contact angle hysteresis, and constants b and m that control the shape of the contact angle-contact line velocity relationship are all chosen

to provide the best fit to the data. The ranges for the equilibrium contact angle and contact angle hysteresis are currently being measured. The ranges for the velocity of the contact line at the contact angle hysteresis and constants b and m are unknown.

Table V.1 lists the values used for these five free parameters.

Table V.1 Free parameters used in the modeling of contact line movement, eqs. II.18, II.24, & II.25.

	$\theta_e(\text{deg})$	$\theta_h(\text{deg})$	$v_m(\text{m/s})$	b	m
first clean	28.0	5.0	.03	8	194
second clean	19.5	4.0	.00007	.5	12467
third clean	24.0	3.5	.0007	.0625	582
first prop.	35.0	6.0	.0005	4	209

Attenuation is described using both viscous and contact line effects. Stiffness is described only by contact line effects. There is a contribution to stiffness from viscosity, but it is not considered. Moerig et al. (1996) showed that the contribution to stiffness from viscosity is negligible.

Viscous Dissipation

All measurements exhibit increasing attenuation with increasing frequency. This phenomena can be easily explained by viscous dissipation. However, the fits in figure V.1 clearly demonstrate that this mechanism for attenuation is inadequate to explain the data below 10 Hz. The measured attenuation is always higher than that predicted by viscous effects. For the first contaminated measurement, a serious

discrepancy is observed. Not only is the measured attenuation much higher below .01 Hz, it is lower above 10 Hz, indicating an additional source of stored energy.

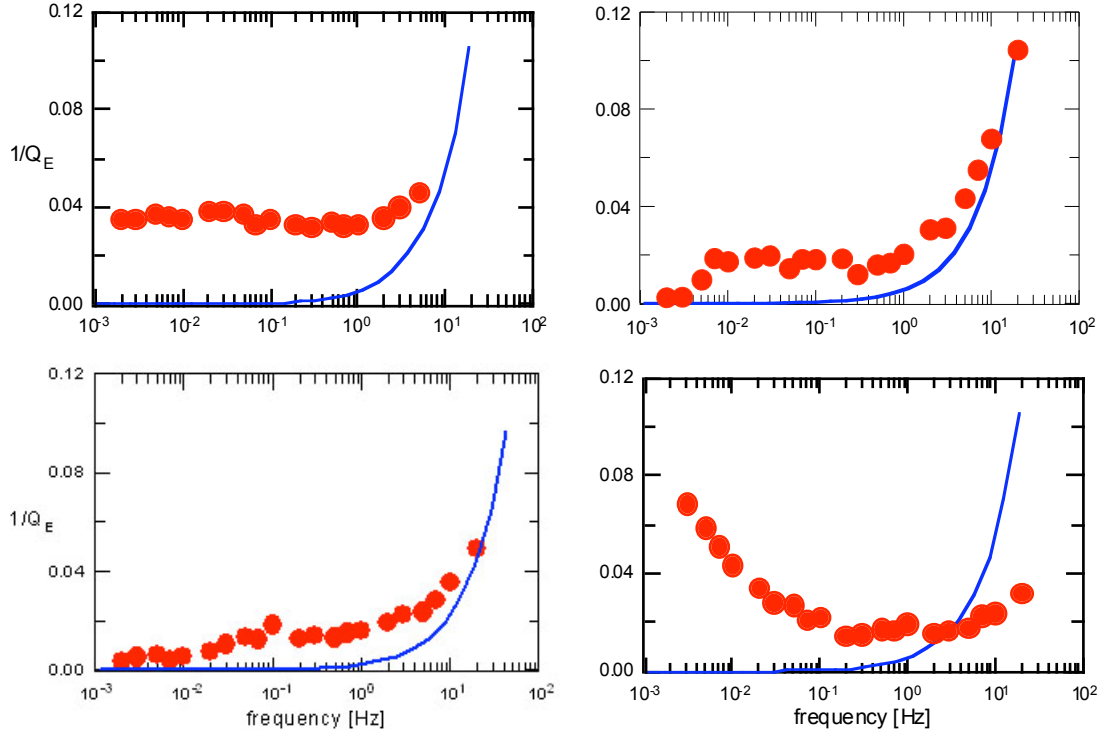


Figure V.1 Attenuation versus frequency; measured data and calculated viscous dissipation. From left to right, top to bottom, first second and third clean measurements, and first contaminated measurement. The clean measurements are modeled well by viscous dissipation at the higher frequencies. The contaminated measurement falls below theoretical predictions indicating either a lack of lost energy or a source of stored energy.

Restricted Contact Line Motion

Below 10 Hz, a definite trend in attenuation is observed on all samples. The clean measurements have a plateau down to $\sim .01$ Hz and then tend to decrease. The contaminated measurements have a similar plateau but begin to rise below .1 Hz. The stiffness increases with increasing frequency and is much higher for the contaminated samples.

The attenuation measurements for all samples can be modeled fairly well by restricted contact line motion between $\sim .01$ and 1 Hz (figures V.2 V.3 V.4 V.5). The modeled attenuation diverges more strongly for the contaminated sample at very low frequencies. The stiffness is much more sensitive. Above $\sim .1$ Hz, the second and third clean measurements have stiffness' that coincide with the model. Below this frequency, the modeled stiffness decreases quickly to the samples dry stiffness. The measured data continues with little change. The first clean sample has a measured stiffness with a slope closely resembling the model. The magnitude, however, is less than calculated. The first contaminated sample is even more anomalous. Measured values are far too low above .01 Hz.

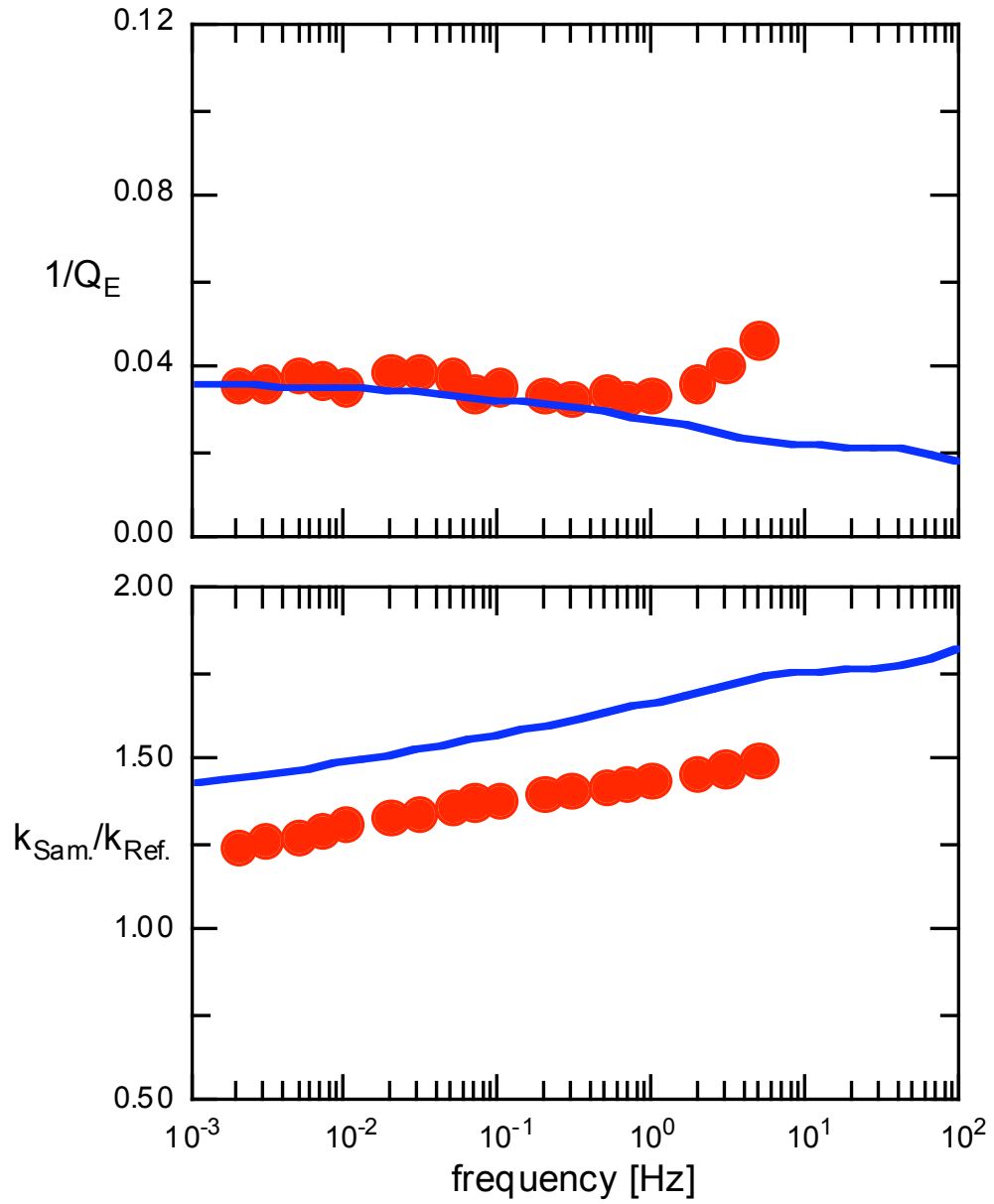


Figure V.2 First clean sample; attenuation and stiffness versus frequency. Solid line indicates modeling from restricted contact line motion only. Note the discrepancy in the stiffness between model and experiment. Model and experimental attenuation diverge above 1 Hz.

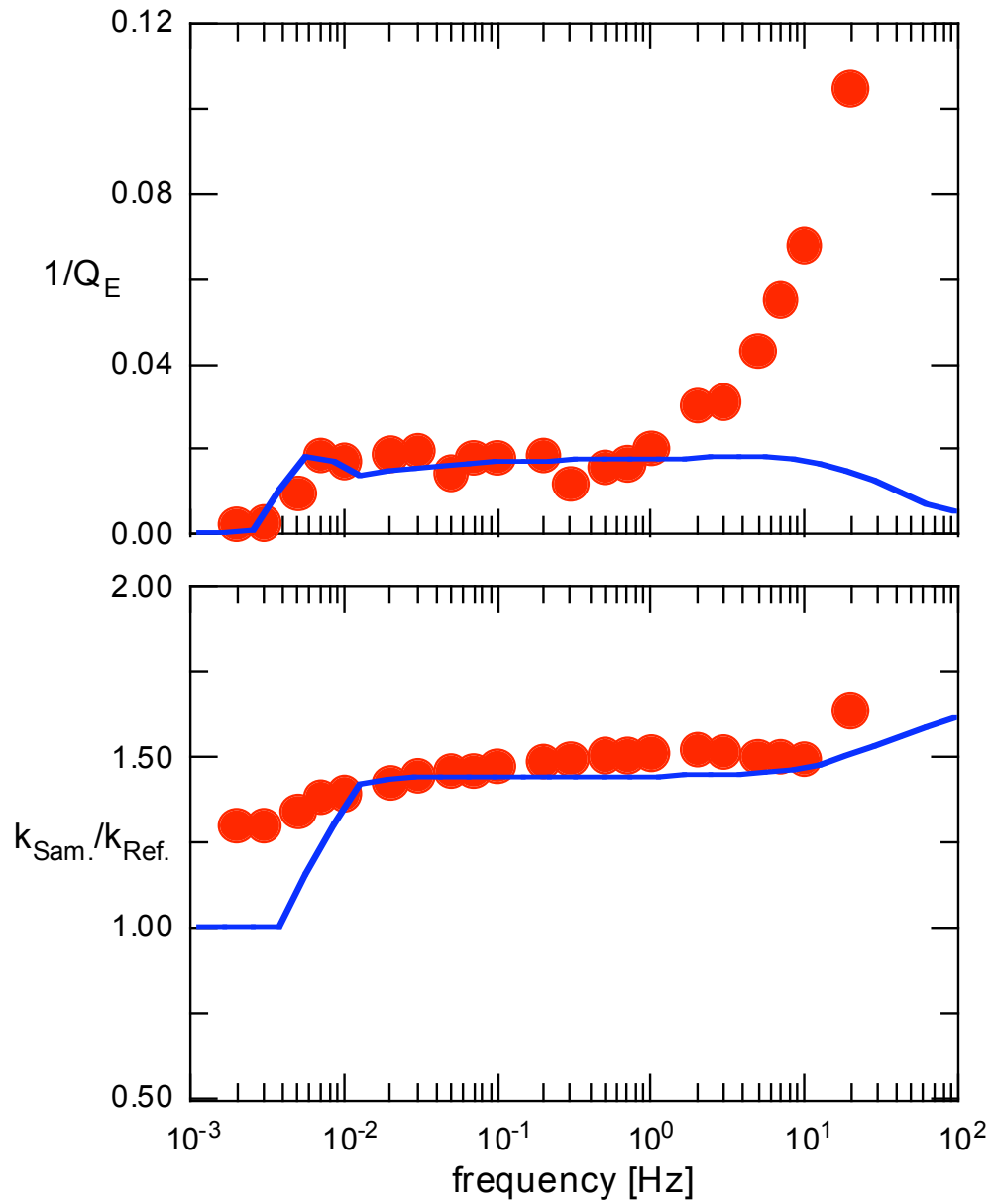


Figure V.3 Second clean sample; attenuation and stiffness versus frequency. Solid line indicates modeling from restricted contact line motion only. Model and experimental attenuation diverge above 1 Hz. Model and experimental stiffness diverge below .01 Hz.

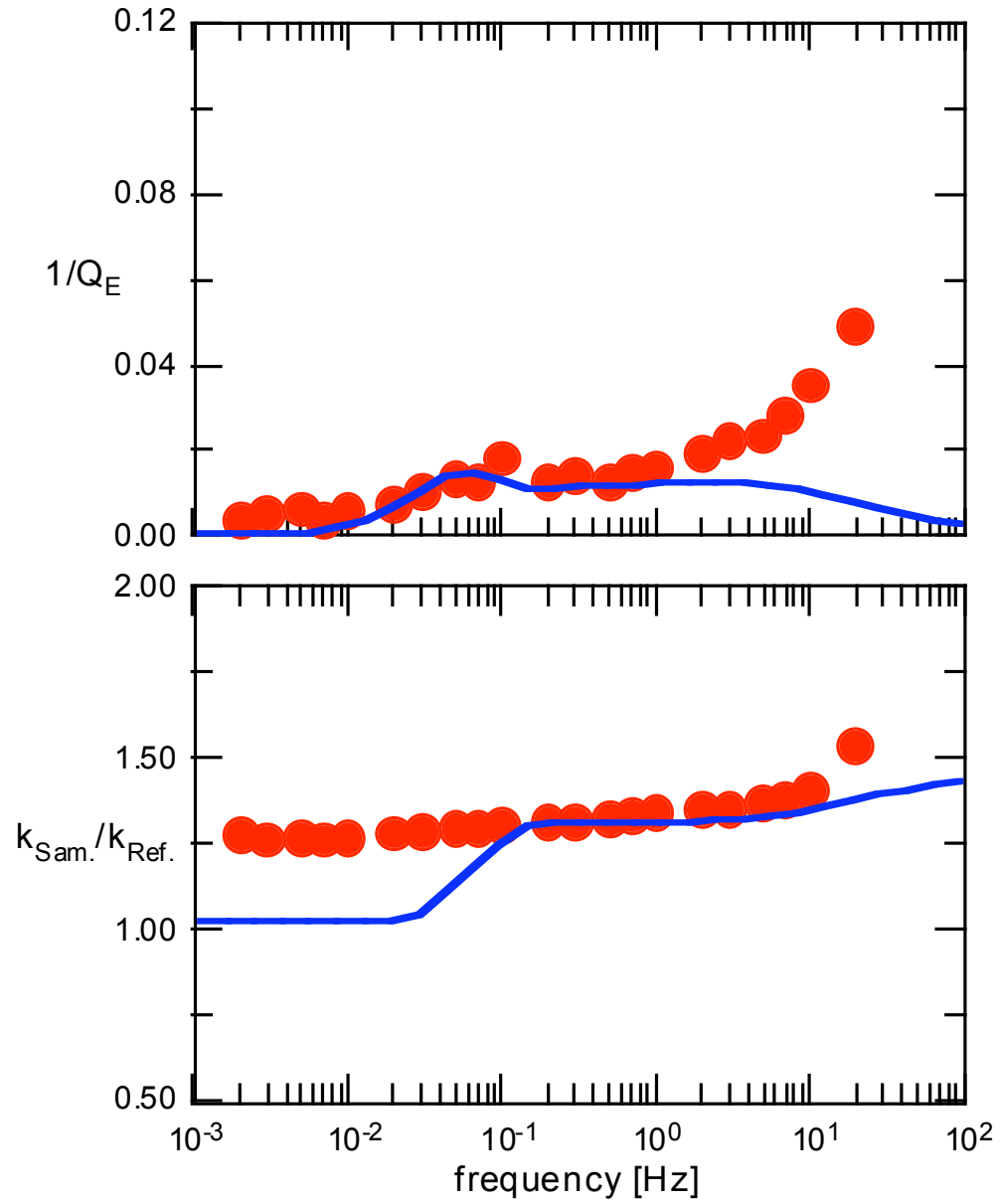


Figure V.4 Third clean sample; attenuation and stiffness versus frequency. Solid line indicates modeling from restricted contact line motion only. Model and experimental attenuation diverge above 1 Hz. Model and experimental stiffness diverge below .1 Hz.

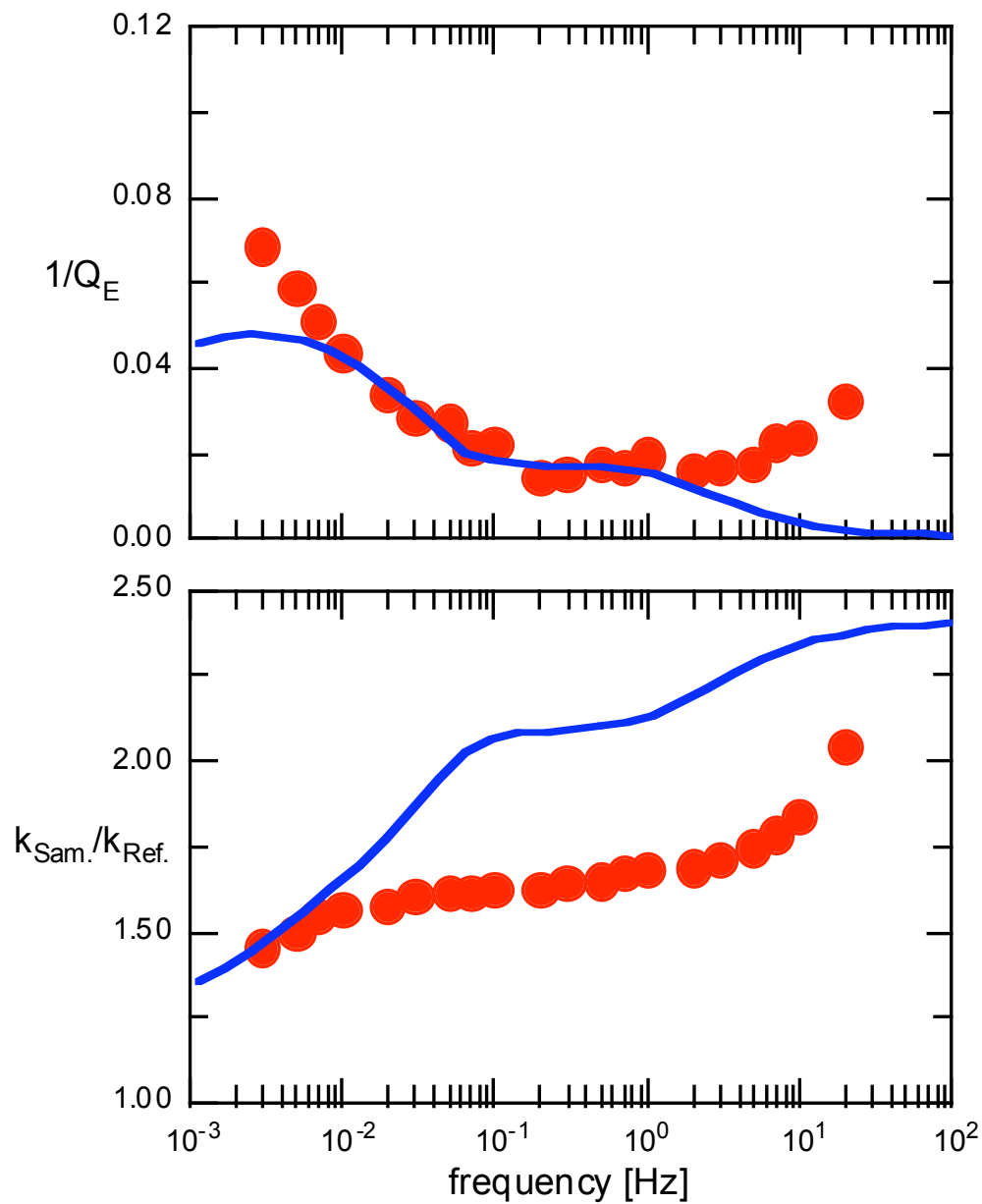


Figure V.5 First Propanol contaminated sample; attenuation and stiffness versus frequency. Solid line indicates modeling from restricted contact line motion only. Note the large discrepancy between model and experimental stiffness. Model and experimental attenuation diverge above 1 Hz and below .01 Hz.

Viscous Dissipation and Restricted Contact Line Motion

The combination of the two mechanisms, viscous dissipation and restricted contact line motion, accounts for nearly all of the attenuation (figure V.6). The restricted contact line model breaks down for contaminated samples below .01 Hz. The contact line resistance in the contaminated sample has provided a source for stored energy and reduced the attenuation at high frequencies.

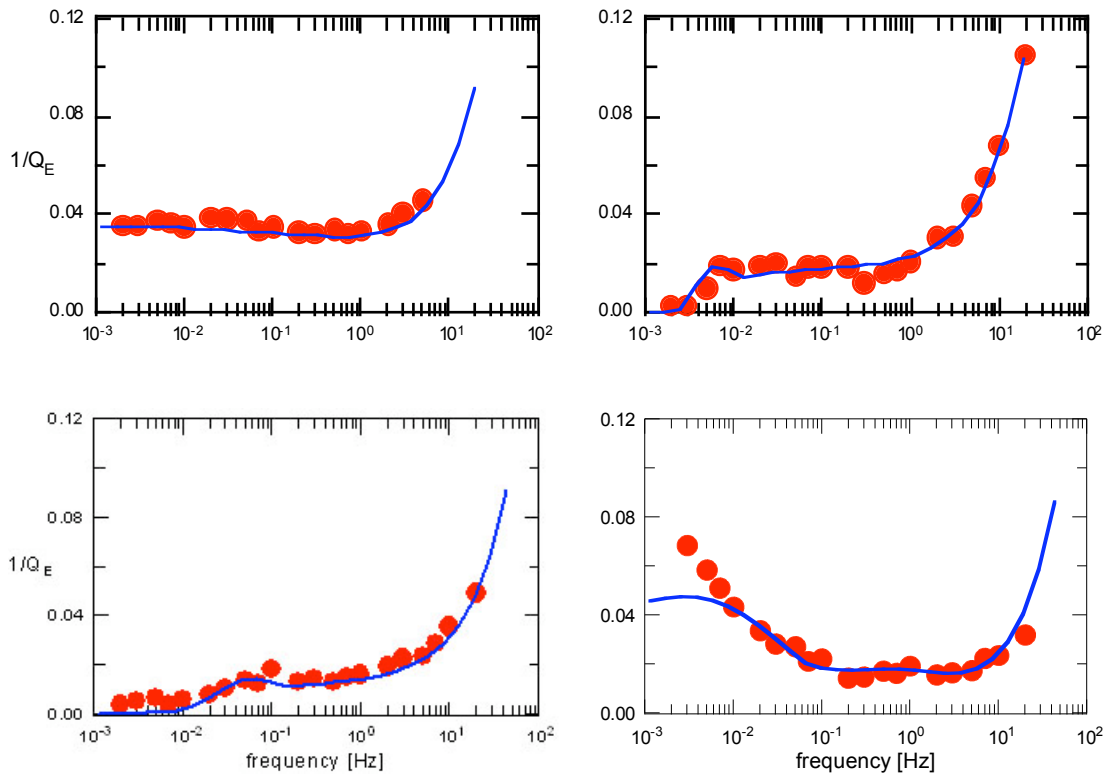


Figure V.6 Attenuation versus frequency; measured data and calculated viscous and contact line dissipation. The combination of the two mechanisms account for nearly all of the experimentally measured attenuation. The source of stored energy in the contaminated measurement needed to reduce the measured attenuation noted earlier can be attributed to contact line resistance.

Chapter VI

DISCUSSION

The description of contact line movement coupled with viscous dissipation as being a source for attenuation shows quantitative agreement with the measured data. The differences between model and experiment in both attenuation and stiffness can probably be attributed to a slightly different shape of the contact line velocity versus contact angle relationship both while sticking and slipping.

The first step would be to make the transition between sticking and slipping continuous. For the sticking regime, a probability that considers the activation energy and other realistic parameters may give the needed relation between contact line velocity and contact angle. It may be possible to derive the contact line velocity versus contact angle relationship empirically from the measured data. Another point to keep in mind is that the dry sample attenuation increased slightly with increasing frequency. This attenuation was not accounted for in the calculations of viscous and contact line related dissipation.

In terms of the values used to model restricted contact line motion, preliminary work suggest that the contact angles are in the proper range. Measurements were made on the slope of the contact line velocity versus contact angle but to date the results are unreliable. Further work should be done to determine the contact line velocity versus contact angle relationship, especially in the sticking regime.

Further work will include measurements on cindered glass bead samples and eventually real rocks. With these new measurements, a higher sample stiffness will increase stored energy and may obscure contact line dissipation.

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Appendix A

MEASUREMENT OF THE ACTIVATION ENERGY FOR THE BONDING OF PROPANOL TO A SODA-LIME GLASS SURFACE

Introduction

The wetting of artificial crack surfaces with water is severely inhibited when the surfaces have been contaminated by 1-propanol, methanol, or acetone. This is observed as a highly irregular solid-liquid-gas contact line and a greatly reduced capillary pressure.

Restoration of precontaminated wetting behavior has been achieved through the application of sustained heat. It is proposed that the restoration of Propanol contaminated soda-lime slides is a thermally activated process.

Method

Soda-lime glass slide surfaces are cleaned by heating the slides in an oven at 420 C for two hours. They are then cooled, fully wetted with 1-propanol, allowed to dry, and placed in the oven at a specific temperature. After a known period of time, two slides are removed and allowed to come to room temperature. The slides are separated by wires (figure 1) and the capillary rise of water is measured.

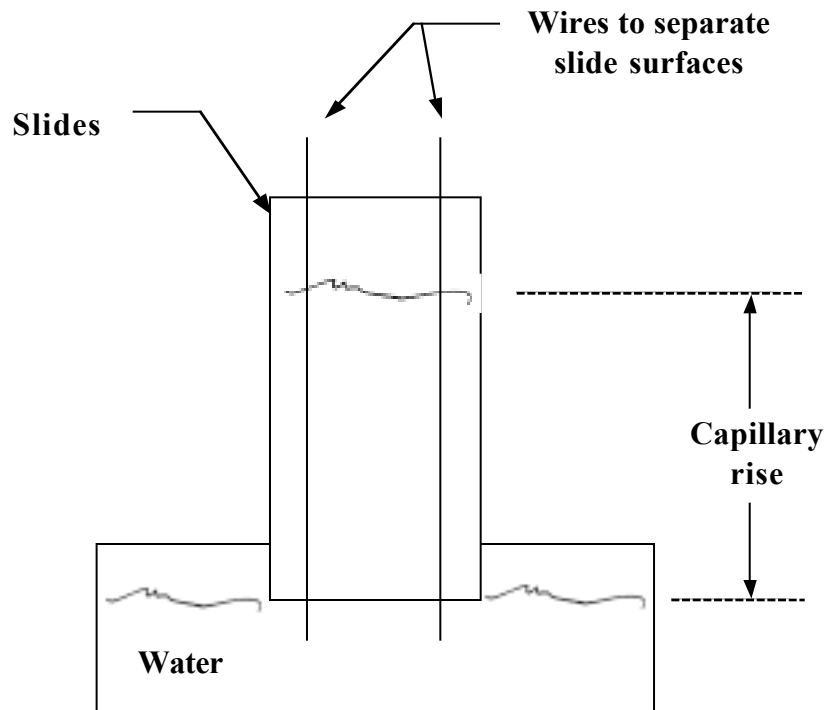


Figure A.1 Measurement of the capillary rise between parallel plates. Water is drawn into the sample. The capillary rise is dependent on the condition of the solid surfaces.

The capillary rise is an indication of the surface contamination. For this experiment, the assumption is made that the greater the capillary rise of water, the fewer the number of affected sites that remain. A site on the surface is one where a bond may be created between the site and an externally imposed molecule (i.e. water).

Assume that the number of affected sites, N , that remain on the surface after exposure to a constant temperature for some period of time is directly proportional to the initial number of affected sites,

$$N = N_0 e^{-Kt}, \quad (\text{A.1})$$

where N_0 is the original number of affected sites and K is a rate constant. If this process of reversion is thermally activated, then

$$K = K_0 e^{-\frac{E_A}{k_B T}}, \quad (\text{A.2})$$

where K_0 is the rate constant at infinite temperature, E_A is the activation energy for this process, T is the temperature and k_B is Boltzman's constant.

K_0 can not be measured directly, but remains constant for the same thermally activated process. By taking the ratio between a reaction rate at one time and temperature and a reaction rate at a different time and temperature, K_0 will cancel. If we choose a constant ratio between remaining affected sites and initial affected sites, N/N_0 , this ratio will also cancel. After taking the logarithm of both sides and some simple algebra, we are left with

$$\ln \frac{t_1}{t_2} = \frac{E_A}{k_B} \left(\frac{1}{T_1} - \frac{1}{T_2} \right). \quad (\text{A.3})$$

This is a straight line for $\ln(t)$ vs. $1/T$ with slope E_A/k_B , characteristic of a thermally activated process.

In order to make use of equation A.3, we must measure the temperatures and find the times for a constant ratio between remaining affected sites and initial affected sites. A constant ratio, N/N_0 , is assumed equivalent to a constant capillary rise. It is

important to choose a capillary rise which is substantially less than that for clean surfaces and substantially more than for fully contaminated surfaces.

It is highly unlikely that several sets of slides removed from the oven at different temperatures would provide the same capillary rise. The time for a constant capillary rise must be extrapolated from multiple sets of slides exposed to the same temperature for differing amounts of time. An Arrhenius equation (eq A.4) is fitted to the measured data to produce a time for the constant capillary rise.

$$\ln\left(\frac{h_0}{h} - b\right) = -\frac{h_0}{h} \frac{E}{RT} = -\frac{h_0}{h} \frac{E}{R} \frac{1}{T}, \quad (\text{A.4})$$

t is the time, τ is a characteristic time provided by the fit, h is the capillary rise, h_0 is the capillary rise at infinite time, and b is a constant to insure that the capillary rise has a particular value at zero time. This technique is repeated for multiple temperatures to provide adequate resolution in the $\ln(t)$ vs. $1/T$ space for the purpose of an estimation of the activation energy.

Results

The collected data are fitted using the values 6 cm for h_0 , and .75 for b . The data and associated fits are seen in figure A.2. In figure A.3, the plot of t versus $1/T$ has a slope of 12520 K where the times are found using a value of 4 cm for the constant capillary rise.

Conclusions

Restoration of Propanol contaminated slides is probably a thermally activated process as shown by the linear relationship between $\ln(t)$ and $1/T$ in figure A.3.

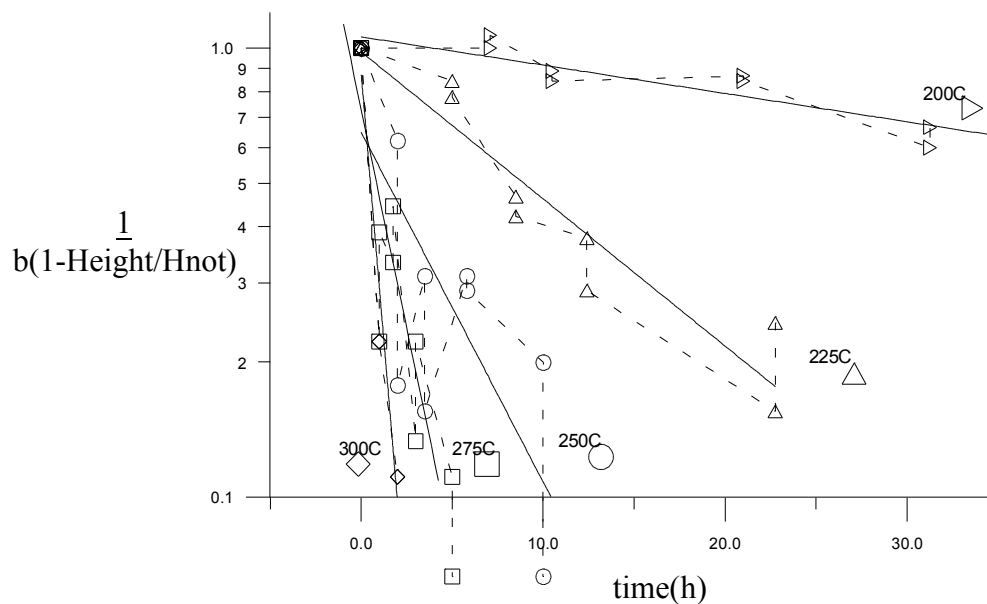


Figure A.2 The logarithmic argument in eq A.9 versus time at temperature. The slope decreases with decreasing temperature as expected.

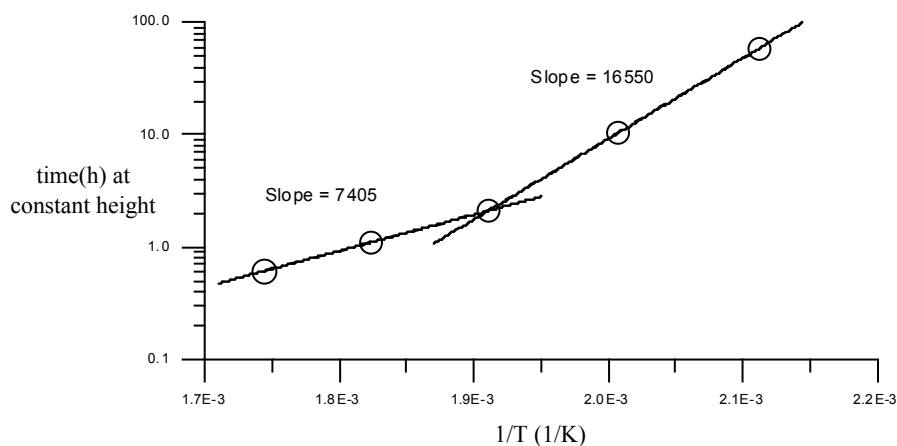


Figure A.3 t versus $1/\text{Temperature}$. Times are determined from a 4 cm capillary rise. Boltzman's constant is 1.380662×10^{-23} J/K and Avagadro's number is 6.022045×10^{23} 1/mol. A slope of 16550 K corresponds to an activation energy of 138 kJ/mol. A slope of 7405 K corresponds to an activation energy of 62 kJ/mol.

Using Boltzman's constant, $k_B = 1.380662 \times 10^{-23}$ J/K, and Avagadro's number, $N_A = 6.022045 \times 10^{23}$ 1/mol, we find that the activation energy is 138 kJ/mol for a slope of 16550 K and 62 kJ/mol for a slope of 7405 K. These experiments also suggest that the effects of contamination are reversible. It would be highly unlikely that two irreversible processes would occur under less than extreme conditions and with these relatively tame molecules; Propanol, water, and silica.

A reversible, thermally activated process implies that bonds between molecules are involved. The activation energy of the high temperature data is on the order of a hydrogen bond, ~ 80 kJ/mol. A common assumption is that water hydrogen bonds to a silica surface. Under certain conditions¹, Propanol is likely to hydrogen bond to surface adsorbed water and other surface sites.

In terms of the restricted flow behavior, a possible explanation may be the following. When a solid is heated and then allowed to cool in an environment with some humidity, water will adsorb to the surface of the solid. The Propanol exposed to the surface bonds to the water molecules by means of a hydrogen bond. The energy in a hydrogen bond between water and Propanol molecules is approximately 84 kJ/mol (Barney Ellison, personal communication), roughly 30 % greater than our smallest measured value for the activation energy. The explanation for this discrepancy is unknown. It is possible that the distance over which the bond is acting is smaller or greater and that there is the presence of an additional attraction between the solid surface and the Propanol molecules.

¹ A clean sample must be contaminated and allowed to fully dry.

Appendix B

DIRECT MEASUREMENTS of CYLINDRICAL SAMPLE STIFFNESS

The samples stiffness is derived by measuring the displacement of the upper holding ring relative to the lower, and the force that caused the displacement (figure B.1).

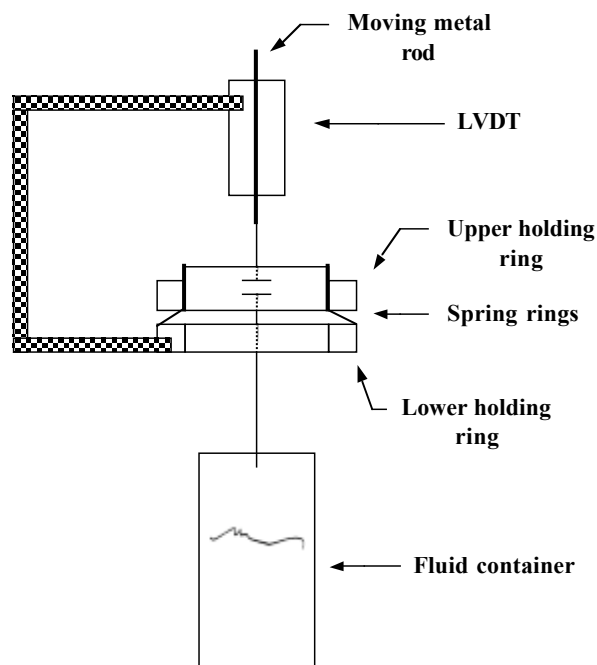


Figure B.1 Schematic of measurement assembly. The lower holding ring is held fixed with respect to the linear variable differential transformer, LVDT. The upper holding ring and metal rod move in response to a change in mass of the fluid container. The response is dependent on the stiffness of the spring rings. A voltage is produced by the LVDT to conclude the measurement.

The samples displacement is measured with a linear variable differential transformer, LVDT. The LVDT has two leads wound around a central hole. In the central hole is placed a metal rod that couples the two leads and produces a voltage. Within a modest range, the voltage output of the transformer is linearly dependent on the position of the rod (figure B.2).

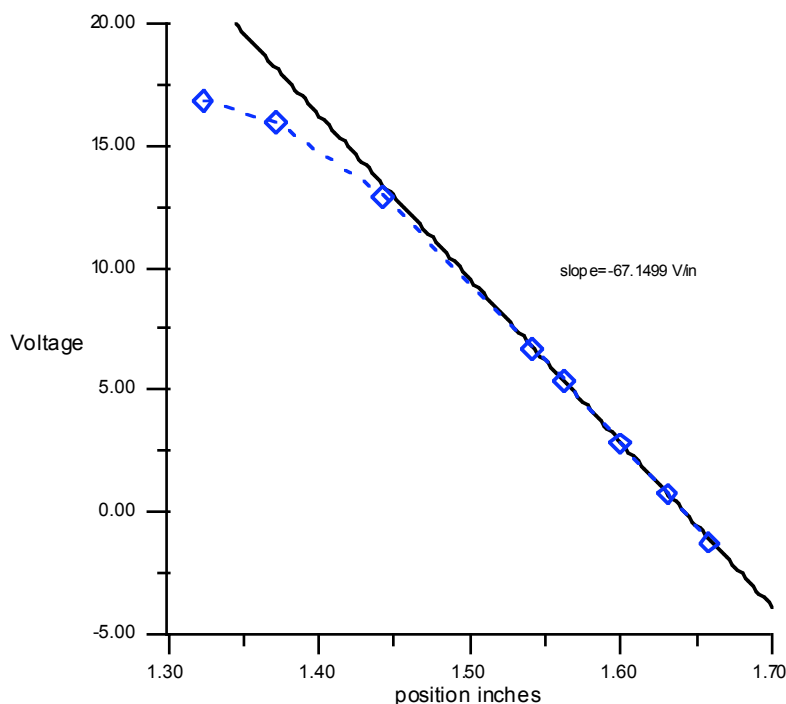


Figure B.2 LVDT output voltage versus position of the metal rod within the LVDT, calibration measurement. The position measurements are all made with a micrometer relative to a point fixed with respect to the LVDT. At zero voltage, the metal rod is centered vertically within the LVDT. The relationship between output voltage and position of the metal rod is linear between $\pm 5V$ and $\pm .1$ in.

The rod is securely fastened to the upper holding ring of the sample. The displacement of the upper holding ring is therefore directly proportional to the voltage produced by the transformer. Force is applied to the sample in the form of a weight. A container in which a measured amount of water is poured, is hung from the upper

holding ring. The force acting on the upper holding ring is the mass of water times the acceleration due to gravity, ΔVg . Having measured both the force and displacement simultaneously, the stiffness can be calculated.

At an input voltage of 15 V, the transformer is initially calibrated and found to have a change in voltage with respect to a change in displacement of ~ 2560 V/m. Mass versus displacement is measured multiple times for the same approximate conditions for both the cylindrical sample and reference (figures B.3 & B.4).

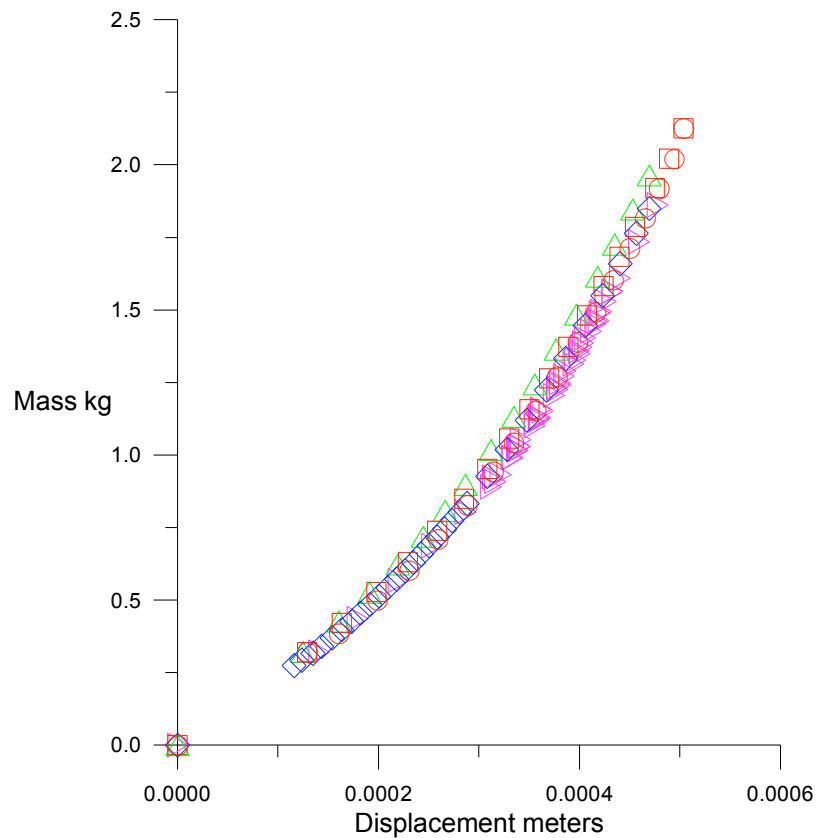


Figure B.3 Mass versus displacement for the cylindrical sample (all measurements). The stiffness is the ratio of mass to displacement times the acceleration due to gravity, the slope of this curve. The slope is changing indicating a stiffness that is not constant.

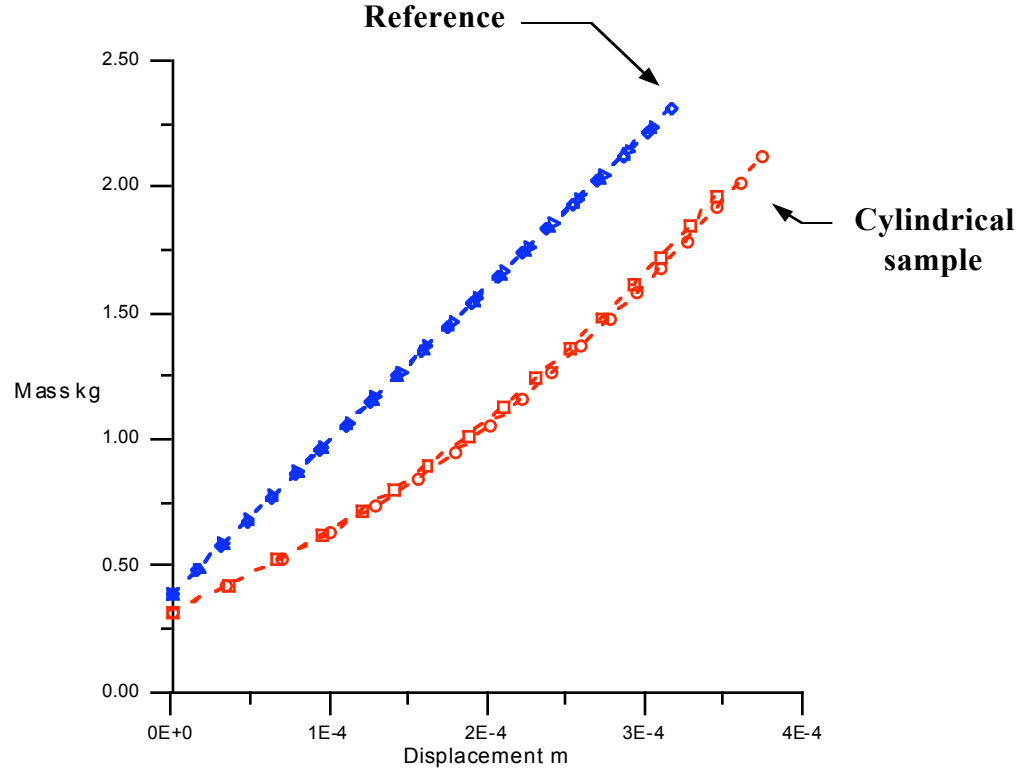


Figure B.4 Mass versus displacement for both the cylindrical sample and reference. The reference has a constant slope, therefore it has a constant stiffness.

The stiffness of the sample and reference is calculated by dividing the change in mass by the change in position and multiplying by the acceleration due to gravity (figures B.5 & B.6). The stiffness of the reference is constant at ~ 59.5 kN/m. The stiffness of the cylindrical sample has a linear relationship with respect to the displacement,

$$k = 1.38 \times 10^8 x + 2.23 \times 10^4 \text{ N/m}, \quad (\text{B.1})$$

and the mass,

$$k = 2.29 \times 10^4 m + 2.52 \times 10^4 \text{ N/m}. \quad (\text{B.2})$$

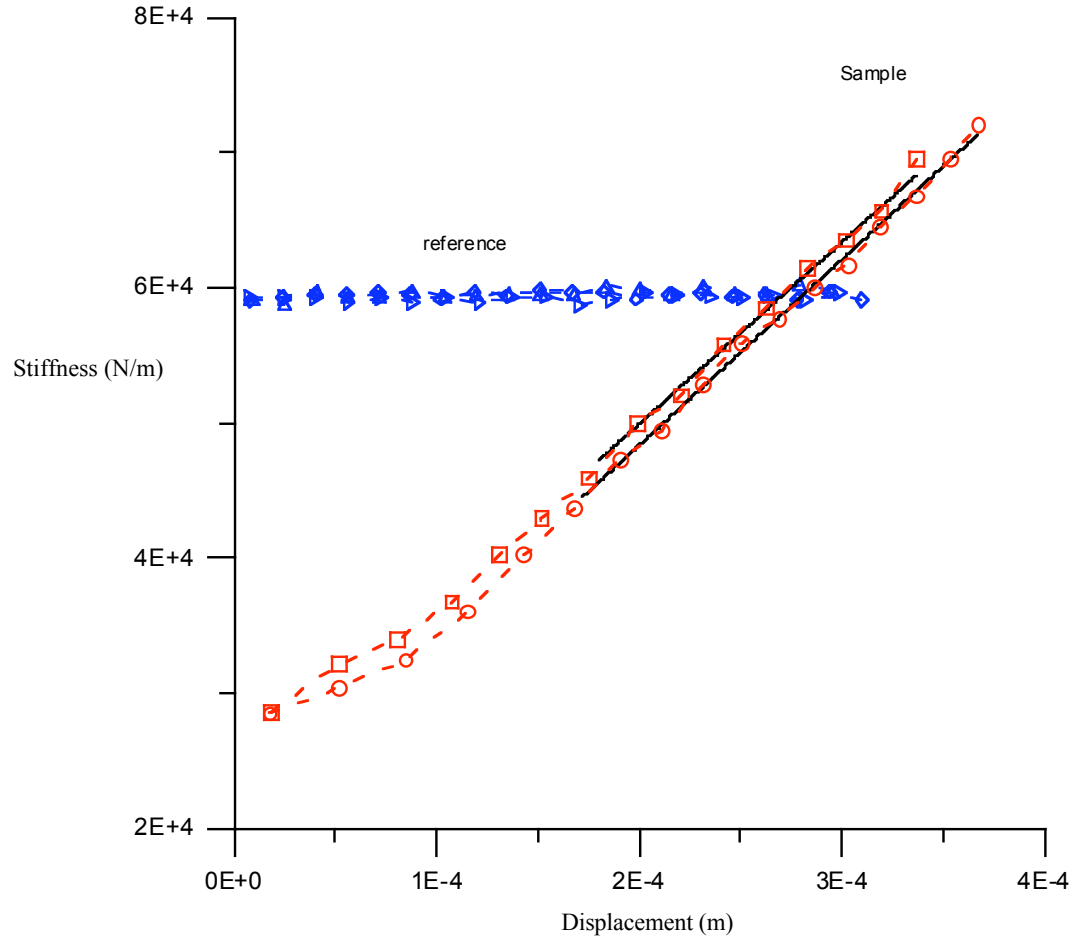


Figure B.5 Stiffness versus displacement for cylindrical sample and reference (2 measurements). The stiffness of the reference is nearly constant. The stiffness of the cylindrical sample is linear with respect to displacement over a range of ~ 0.2 mm.

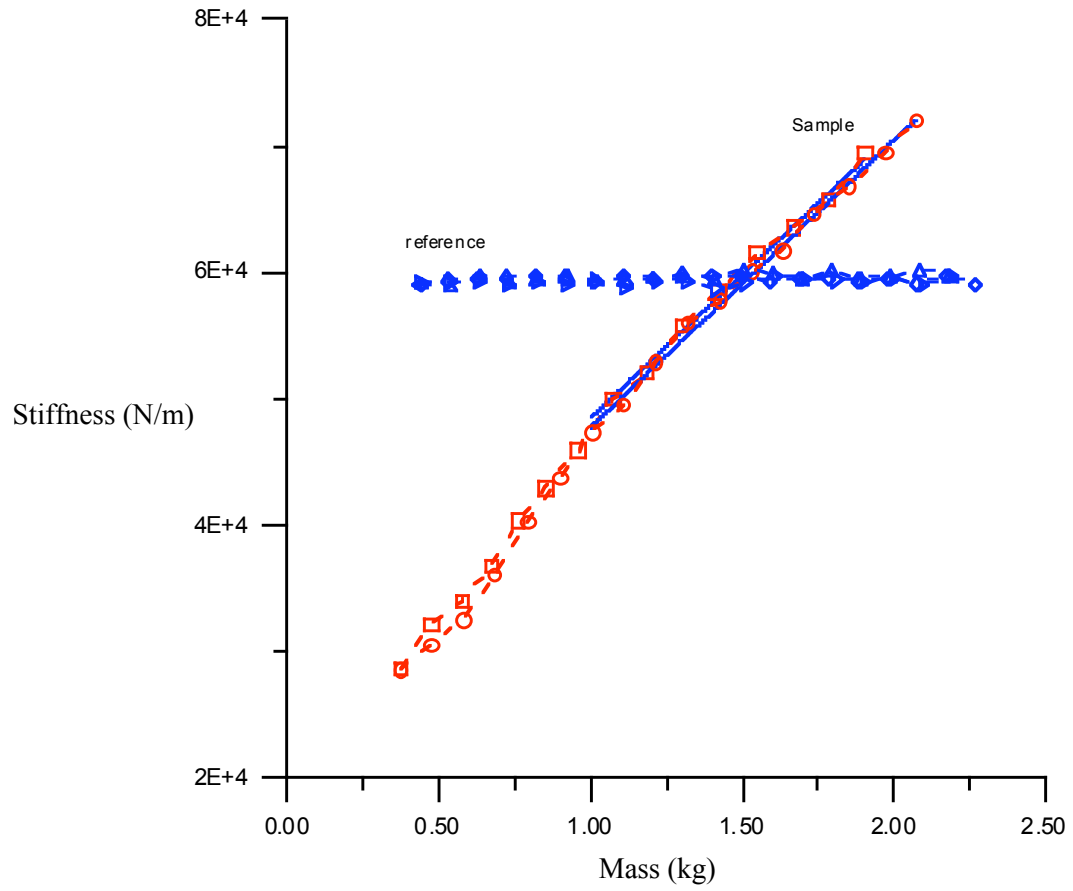


Figure B.6 Stiffness versus mass for cylindrical sample and reference. The reference stiffness is independent of applied force between .5 and 2.0 kg. The cylindrical sample stiffness increases linearly between 1.0 and 2.0 kg.

Appendix C

COMPUTER MODEL

Introduction

The following program is written in C++. The program could be compiled to run on any computer with any operating system supporting a C++ compiler. The program calculates the energy lost due to both viscous effects, E_{lossv} , and resistive contact line motion, E_{loss} , in dependence on frequency of sinusoidal deformation, $freq$. The program consists of a main routine, `main()`, and three subroutines, `guess()`, `setsvv()`, and `vol()`. Its 9 input arguments are equilibrium contact angle in degrees, $dcont$, contact angle hysteresis, $cont_hyst$, initial radius of fluid volume, Ri , initial gap separation, $orig_ht$, amplitude of deformation, Amp , velocity of the contact line at the contact angle hysteresis, vm , constant b controlling slope of sticking regime, dry stiffness scaling factor, kl_{dry} , and slope of slipping regime in deg s/m, a .

Within the `main()` routine, a moving contact line and changing contact angle are described for multiple frequencies between .001 and 100 Hz. Each frequency occurs over 1/4 of a cycle and is broken up into 1000 time steps. During each of these increments, the change in gap separation is calculated. Due to conservation of volume, the deformation results in a change in contact angle. As the contact angle moves from its equilibrium state, the probability of the radius changing increases.

The energy lost is calculated and summed over every time step. At the end of the time loop the stiffness and energy stored due to contact line movement, and the energy lost due to viscous effects are calculated.

The following is the basic structure of the program

Subroutines()

main()

I. Loop over frequencies, .001 Hz to 100 Hz.

A. Loop over quarter cycle with increment dt

1. Change in contact angle, dc, radius, R

2. Calculate Eloss

B. Calculate k1, Estor, Elossv, Atten

II. Output data

Code

```

#include <stdio.h>
#include <math.h>
#include <stdlib.h>
#include <time.h>
#define small .000000001
#define pi 3.14159265
#define rho .001
#define mslwmirpl .6

//GUESS ROUTINE USED TO DETERMINE NEW CONTACT ANGLE
double guess(double pre, double pre_pre, double dif, double pre_dif)
{
    double slope, predicted, differ;
    if (!(differ = dif - pre_dif)) differ = small;
    slope = dif * (pre - pre_pre)/differ;
    while (1)
    {
        predicted = pre - slope;
        if (predicted < 0)
        {
            printf("Error: predicted %e, pre %f pre_pre %f dif %f pre_dif\n", predicted, pre, pre_pre, dif, pre_dif);
            exit(-2);
        }
        if (predicted < pi/2.0) break;
        slope /= 2.0;
    }
    return predicted;
}

double svvRsqr, svvhsqr; // pi*R*R*h and pi*h*h*R

// SUBROUTINE TO SET VOLUME COEFFICIENTS

void setsvv(double R, double h)
{
    svvhsqr = pi * h * h * R;
    svvRsqr = pi * h * R * R;
}

long num_vol_calls;
//VOLUME SUBROUTINE USING EQ. II.20

double vol(double contact_angle)
{
    double beta, sinbeta;
    ++num_vol_calls;
    if (!(beta = pi/2.0-contact_angle)) return svvRsqr;
    sinbeta = sin(beta);
    return svvRsqr - svvhsqr * .5 * (beta - cos(beta)*sinbeta)/(sinbeta*sinbeta);
}

//MAIN ROUTINE

int main(int narg, char *argv[])
{
    // All units in millimeters, kilograms, seconds unless otherwise noted.

    double xm, X, xcap, xini, x1;
    int it, i, flag;
    double Amp, orig_ht, dhdt, Uh;
    double new_ht;
    double k1, k2 = 59459.3; //k1, stiffness of sample. k2, stiffness of reference in kN/m
    double volume_cur, volume_pre, Area, Areai, R, Ri, dR;
    double cont_an_pre, dif, dif_pre, can_pre_pre;
    double surface_tension = .07275, n = .0000010019; //Surface tension and viscosity of water

```

[illegible]

```

        dhdt = Amp * 2 * pi * freq * cos(2 * pi * freq * t);
        sum += dhdt*dt;
    }
    scale = Amp/sum;

//
    THE TIME LOOP

    for(t = 0; t <= qcycle; t += dt)
    {
        if ((dc = contact_angle2 - contact_angle1) > 0)
        {
            if (dc < (cont_hyst * pi/180))
                dR = dt * vm * pow(180 * dc/(pi * cont_hyst),b);
            else
                dR = ((dc - cont_hyst * pi/180)/a + vm) * dt;
            R = R + dR;
        }
        Elost = Elost + 4 * pi * surface_tension * (cos(contact_angle1) -
            cos(contact_angle2))*R*dR;
        fprintf(outmat2, "%f\t%f\t%f\t%f\t%f\n", t, log10(freq), R, dc*180/pi);
        dhdt = Amp * 2 * pi * freq * cos(2 * pi * freq * t);
        new_ht := dhdt*dt*scale;
        setsvv(R, new_ht); //Set up volume constants
        dif_pre = volume_pre - vol(contact_angle2);
        double sv_ca_in = cont_an_pre = contact_angle2;
        contact_angle2 += prev_delta_ca;
        it = 0;
        do
        {
            volume_cur = vol(contact_angle2);
            if (!(dif = volume_pre - volume_cur)) break;
            can_pre_pre = cont_an_pre;
            cont_an_pre = contact_angle2;
            contact_angle2 = guess(cont_an_pre, can_pre_pre, dif, dif_pre);
            dif_pre = dif;
            it = it + 1;
        }
        while(dif_pre > small || dif_pre < -small);
        prev_delta_ca = contact_angle2 - sv_ca_in;
    } //End of time loop

    Areai = pi * Ri * Ri;
    Area = pi * R * R;
    cap_pres_i = 4 * surface_tension * cos(contact_angle1)/orig_ht;
    cap_pres = 4 * surface_tension * cos(contact_angle2)/new_ht;
    dFcap = Areai * cap_pres_i - Area * cap_pres;
    k1 = dFcap/Amp + k1dry * k2;

//
    Calculate energy stored and energy lost

    Estor = dFcap * dFcap / (2 * k2);
    Uh = Amp * freq;
    Elossv = 3*pi*n*Ri*Ri*Ri*Ri*Uh/(2*orig_ht*orig_ht);
    Atten = (Elost + Elossv)/(2 * pi * (k2 * Amp * Amp/2 + Estor));
    Atten_no_v = Elost/(2 * pi * (k2 * Amp * Amp/2 + Estor));
    Atten_no_c = (Elossv)/(2 * pi * (k2 * Amp * Amp/2));
    argatten = wr*wr * tan(Atten)/(wr*wr-freq*freq);
    attenwrs = atan(argatten);

    fprintf(out, "%f,%f,%f,%f,%f,%f\n", freq, Atten, Kave/k2, Eave, Rave, cont_ang_ave * 180/pi);
    fprintf(outmat, "%f\t%f\t%f\t%f\t%f\n", log10(freq), a, Atten_no_v, k1);
    fprintf(outdat, "%f,%f,%f,%f,%f,%f,%f,%f,%f,%f\n", freq, Atten, attenwrs,
        Atten_no_v, Atten_no_c, k1/k2, Elost, Estor, R, contact_angle2 * 180/pi);
    } //End of frequency loop
fclose(outmat2);
fclose(outmat);
fclose(outdat);
return 0;
}

```

Appendix D

DATA AQUISITION PROGRAM

This program is written in basic and requires the Hi-Tech Basic and DOS operating systems.

```

1! RE-STORE "Q2da"
2!
3! THE FOLLOWING IS USED TO LOAD THE NATIONAL INSTRUMENTS AT-GPIB
4! DRIVER AND FOLLOWING COMMAND SETS COM1 TO THE RS232
6! LOAD BIN "/HTB386/GPIBN;BO AT-GPIB BA 2C0 IN 11"
7! LOAD BIN "SERIAL;DRIVER 1 DISABLE"
8! THE MAJORITY OF THIS PROGRAM WAS WRITTEN BY JOSEF PAFFENHOLZ
9! MODIFICATIONS MADE BY ROB ANDZIK (6/18/90) AND OLIVER BOYD (6/10/94)
11 COM /Nicdata/ REAL Vrange2,Period,Freq,Ckldiv
14 COM /Filename/User$[5],Path$[20],Date$[10],INTEGER Nr
15 COM /Flag/ INTEGER Flag,Storeflag$[15],Stackingflag$[15]
16 COM /Set_up/ INTEGER Setflag
17 COM /Ct/ REAL Ct_output,@Rs232
18 COM /Mirmove/ REAL Ab,Bc,Cd,De,Ef,Fa
20 COM /Infor/Ident$[3],Samp$[20],Satur$[20],Comm$[80],Prmir$[1]
21 COM /Extra/Fate$[4]
22 COM /Zeros/Nr$[3]
24 User$="Q2"
25 ASSIGN @Wavetek TO 709
26 ASSIGN @Rs232 TO 11
27 CALL Init_rs232
28 CALL Init_motor
29 CALL Look_for_date
30 Nr=0
31 Ghj=1
33 Ab=-154000
34 Bc=-259000
35 Cd=-149000
36 De=-162000
37 Ef=-233000
38 Fa=-195000
39 Prmir$="a"
40 Path$="/"&User$&"/DAT/"&Date$
41 OUTPUT @Rs232;"SC1"
42 OUTPUT @Rs232;"ST1"
43 ON ERROR CALL Read_dir
44 CREATE DIR Path$
45 OFF ERROR
46 Storeflag$="STORAGE ON"
47 CS=CHR$(255)&"K"
48 CLEAR SCREEN
49 PRINT "YOUR OPTIONS: (1)PLAY WITH WAVETEK"
50 PRINT "          (2)PLAY WITH MIRRORS"
51 PRINT "          (3)CONFIGURE ANALOGIC"
52 PRINT "          (4)TAKE MEASUREMENTS"
53 PRINT ""
54 PRINT ""
55 PRINT ""
56 IF Ghj=1 THEN
58 PRINT " ***** NOTE *****"
59 PRINT ""
60 PRINT " If measurement has been interrupted, you must"
61 PRINT "   make sure that the variable prmir$(line 39),"
62 PRINT "   initially mir a, has the value of the"
63 PRINT "   present mirror"
64 PRINT ""
65 PRINT "present mirror should be: ",Prmir$
67 Ghj=2
68 END IF
69 INPUT I
70 IF I=1 THEN GOSUB Wavetek
71 IF I=2 THEN GOSUB Mirrors
72 IF I=3 THEN CALL Config_ana
73 IF I=4 THEN CALL Set_up
74 GOTO 48
75 Wavetek:
76 OUTPUT @Wavetek;"R"

```

```

77 CLEAR SCREEN
78 INPUT "AMPLITUDE =?",Amp
79 INPUT "FREQUENCY =?",Freqw
80 O_freq=Freqw*8192
81 OUTPUT @Wavetek;"C0F"&VAL$(O_freq)&"A"&VAL$(Amp)&"I"
82 RETURN
83 Mirrors:!
84 CLEAR SCREEN
85 REAL J
86 Ch=2
87 INPUT "By step(1) or by mirror(enter)?",Ch
88 IF Ch=1 THEN
89   INPUT "HOW MANY STEPS?",J
90 ELSE
91   INPUT "Which Mirror to go to (lower case)?",Mir$
92   IF Prmir$="a" THEN
93     IF Mir$="a" THEN J=0
94     IF Mir$="b" THEN J=Ab
95     IF Mir$="c" THEN J=Ab+Bc
96     IF Mir$="d" THEN J=-Fa-Ef-De
97     IF Mir$="e" THEN J=-Fa-Ef
98     IF Mir$="f" THEN J=-Fa
99   END IF
100  IF Prmir$="b" THEN
101    IF Mir$="b" THEN J=0
102    IF Mir$="c" THEN J=Bc
103    IF Mir$="d" THEN J=Bc+Cd
104    IF Mir$="e" THEN J=-Ab-Fa-Ef
105    IF Mir$="f" THEN J=-Ab-Fa
106    IF Mir$="a" THEN J=-Ab
107  END IF
108  IF Prmir$="c" THEN
109    IF Mir$="c" THEN J=0
110    IF Mir$="d" THEN J=Cd
111    IF Mir$="e" THEN J=Cd+De
112    IF Mir$="f" THEN J=-Bc-Ab-Fa
113    IF Mir$="a" THEN J=-Bc-Ab
114    IF Mir$="b" THEN J=-Bc
115  END IF
116  IF Prmir$="d" THEN
117    IF Mir$="d" THEN J=0
118    IF Mir$="e" THEN J=De
119    IF Mir$="f" THEN J=De+Ef
120    IF Mir$="a" THEN J=-Cd-Bc-Ab
121    IF Mir$="b" THEN J=-Cd-Bc
122    IF Mir$="c" THEN J=-Cd
123  END IF
124  IF Prmir$="e" THEN
125    IF Mir$="e" THEN J=0
126    IF Mir$="f" THEN J=Ef
127    IF Mir$="a" THEN J=Ef+Fa
128    IF Mir$="b" THEN J=-De-Cd-Bc
129    IF Mir$="c" THEN J=-De-Cd
130    IF Mir$="d" THEN J=-De
131  END IF
132  IF Prmir$="f" THEN
133    IF Mir$="f" THEN J=0
134    IF Mir$="a" THEN J=Fa
135    IF Mir$="b" THEN J=Fa+Ab
136    IF Mir$="c" THEN J=-Ef-De-Cd
137    IF Mir$="d" THEN J=-Ef-De
138    IF Mir$="e" THEN J=-Ef
139  END IF
140  Prmir$=Mir$
141  IF Gl=1 THEN Gl=Gl+1
142 END IF
143 OUTPUT @Rs232;"ST0"
144 WAIT .1
145 OUTPUT @Rs232;"A1.5 V3 D"&VAL$(J)&" G"
146 OUTPUT @Rs232;"ST1"

```

```

147 RETURN
148 END
149!*****
150 SUB Init_rs232
151 CONTROL 11,0;1
152 CONTROL 11,13;9600
153 CONTROL 11,14;3
154 CONTROL 11,5;3
155 CONTROL 11,12;16+32+128
156 SUBEND
157!*****
158 SUB Init_motor
159 COM /Ct/ REAL Ct_output,@Rs232
160 OUTPUT @Rs232;"8LD3"
161 OUTPUT @Rs232;"A3"
162 OUTPUT @Rs232;"K"
163 SUBEND
164!*****
165 SUB Look_for_date
166 COM /Filename/User$,Path$,Date$,INTEGER Nr
167 COM /Extra/Fate$
168 D$=DATE$(TIMEDATE)
169 M$=D$[4,6] !MONTH
170 D$=D$[1,2] !DAY
171 D$=TRIM$(D$)
172 Add_zero(D$)
173 SELECT M$
174 CASE "Jan"
175 M$="01"
176 CASE "Feb"
177 M$="02"
178 CASE "Mar"
179 M$="03"
180 CASE "Apr"
181 M$="04"
182 CASE "May"
183 M$="05"
184 CASE "Jun"
185 M$="06"
186 CASE "Jul"
187 M$="07"
188 CASE "Aug"
189 M$="08"
190 CASE "Sep"
191 M$="09"
192 CASE "Oct"
193 M$="10"
194 CASE "Nov"
195 M$="11"
196 CASE "Dec"
197 M$="12"
198 END SELECT
199 CALL Add_zero(M$)
200 Fate$=M$&D$
201 Date$=M$&D$&"_"
202 SUBEND
203 !*****
204 SUB Set_up
205 COM /Niddata/ REAL Vrange2,Period,Freq,Ckldiv
206 COM /Set_up/ INTEGER Setflag
207 COM /Filename/User$,Path$,Date$,INTEGER Nr
208 COM /Mirmove/ REAL Ab,Bc,Cd,De,Ef,Fa
209 COM /Infom/Ident$[3],Samp$[20],Satur$[20],Comm$[80],Prmir$[1]
210 COM /Extra/Fate$
211 COM /Ct/ REAL Ct_output,@Rs232
212 DIM Response$[1],Accept$[20]
213 DIM Freqa(30)
214 DIM Frtyp$[4]
215 INTEGER Num
216 ASSIGN @Rs232 TO 11

```



```

217 ASSIGN @Wavetek TO 709
218 Anafig$="N"
219 OUTPUT @Wavetek;"R"
220 CLEAR SCREEN
221 IF Pmir$="a" THEN
222   GOSUB Info
223 ELSE
224   REAL Mo
226   PRINT ""
227   PRINT ""
228   PRINT "          LASER LIGHT MUST BE AT MIRROR A."
229   PRINT ""
230   PRINT "    It is currently on mirror ",Pmir$
231   PRINT ""
232   PRINT "    Now moving to a, are you ready to play?"
234   IF Pmir$="b" THEN Mo=-Ab
235   IF Pmir$="c" THEN Mo=-Bc-Ab
236   IF Pmir$="d" THEN Mo=-Cd-Bc-Ab
237   IF Pmir$="e" THEN Mo=Ef+Fa
238   IF Pmir$="f" THEN Mo=Fa
239   OUTPUT @Rs232;"ST0"
240   WAIT .1
241   OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Mo)&" G"
242   OUTPUT @Rs232;"ST1"
243   INPUT "    ** Press return to continue **",Dummy
245   GOSUB Info
246 END IF
247!The variable H corresponds to the frequency
249 FOR H=1 TO Num
250   Freq=Freqa(H)
251   Out_freq=Freq*8192
252   Amp_s=.04*LOG(Freq)+Amp
254   OUTPUT @Wavetek;"C0F"&VAL$(Out_freq)&"A"&VAL$(Amp_s)&"I"
255   IF H=1 THEN GOSUB Start_up
256   GOSUB Start
257 NEXT H
258 GOSUB End
259 Info:
260 CLEAR SCREEN
261 ASSIGN @Ff TO "c:\q2\prg\q2dr.inf";FORMAT ON
262 INPUT "Y-intercept ( = 3.6050 for 2.8 fringes at current gains)? ",Amp
263 INPUT "Type of frequency range (low,test,manu)? ",Frtyp$
264 IF Frtyp$="test" THEN
265   Num=1
266   Freqa(1)=10
267   OUTPUT @Ff;Fate$,VAL$(Nr+1),VAL$(Num*6),"xxx"
268   GOTO 293
269 END IF
270 IF Frtyp$="manu" THEN
271   INPUT "How many frequencies (max=21)? ",Num
272   FOR I=1 TO Num
273     PRINT "Frequency #"&VAL$(I)&"?"
274     INPUT K
275     Freqa(I)=K
276   NEXT I
277 ELSE
278   ASSIGN @Ffile TO "q2da.fre";FORMAT ON
279   ENTER @Ffile;Num
280   PRINT Num
281   FOR I=1 TO Num
282     ENTER @Ffile;Freqa(I)
283     PRINT Freqa(I)
284   NEXT I
285 END IF
286 INPUT "Three letter identifier? ",Ident$
287 PRINT TABXY(1,30),"Capital X is a variable which you supply."
288 INPUT "Sample type (X glass slides, rock type, glass beads)?",Samp$
289 INPUT "Saturant (none, water, methanol, etc.)?",Satur$
290 PRINT TABXY(1,30),"sampXX, X.Xmm sl., Xx prestress, and other comments"
291 INPUT "Comments (within double quotes)?",Comm$

```

```

292 OUTPUT @Ff,Fate$,VAL$(Nr+1),VAL$(Num*6),Ident$,Samp$,Satur$,Comm$
293 RETURN
294 Start:!
295 !The variable M corresponds to the mirror at a given frequency
296 FOR M=1 TO 6
297 GOSUB Ttl
298 GOSUB Output_ttl
299 IF M=1 THEN GOTO 308
300 OUTPUT @Rs232;"ST0"
301 WAIT .1
302 IF M=2 THEN OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Ab)&" G"
303 IF M=3 THEN OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Bc)&" G"
304 IF M=4 THEN OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Cd)&" G"
305 IF M=5 THEN OUTPUT @Rs232;"A1.5 V3 D"&VAL$(De)&" G"
306 IF M=6 THEN OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Ef)&" G"
307 OUTPUT @Rs232;"ST1"
308 PRINT "ATTEMPT #"&VAL$(M)&" AT FREQ:"&VAL$(Freq)&"("&VAL$(H)&") using WaveTek voltage:
"&VAL$(Amp_s)
309 WAIT 20
310 CALL Read_data
311 NEXT M
312 OUTPUT @Rs232;"ST0"
313 WAIT .1
314 OUTPUT @Rs232;"A1.5 V3 D"&VAL$(Fa)&" G"
315 OUTPUT @Rs232;"ST1"
316 RETURN
317 Ttl:!
318 IF Freq<.0625 THEN
319   Clkdiv=2
320   GOTO 339
321 END IF
322 IF Freq<.125 THEN
323   Clkdiv=4
324   GOTO 339
325 END IF
326 IF Freq<.25 THEN
327   Clkdiv=8
328   GOTO 339
329 END IF
330 IF Freq<.5 THEN
331   Clkdiv=16
332   GOTO 339
333 END IF
334 IF Freq<1 THEN
335   Clkdiv=32
336   GOTO 339
337 END IF
338 Clkdiv=64
339 RETURN
340 Output_ttl:!
341 Ttlport=&H3BC
342 Divisor=256-Clkdiv
343 OUT Ttlport,Divisor
344 RETURN
345 Start_up:!
346 CLEAR SCREEN
347 INPUT "CONFIGURE ANALOGIC DIGITIZER ?, Y/N [N]",Anafig$
348 IF Anafig$="Y" OR Anafig$="y" THEN CALL Config_ana
349 !
350 IMAGE 16/,30A,18A
351 CLEAR SCREEN
352 BEEP 240,.1
353 BEEP 440,.1
354 Setflag=0
355 RETURN
356 End:!
357 QUIT
358 SUBEND
359 !*****
360 SUB Read_dir

```

```

361 COM /Filename/User$,Path$,Date$,INTEGER Nr
362 DIM Cat$(1:400)[80],File$(1:400)[14]
363 CAT "/"&User$&"/DAT/"&Date$ TO Cat$(*);NO HEADER,COUNT Nr
364 ERROR SUBEXIT
365 SUBEND
366 !*****
367 SUB Read_data
368 COM /Nicdata/ REAL Vrange2,Period,Freq,Ckdiv
369 COM /Filename/User$,Path$,Date$,INTEGER Nr
370 COM /Flag/ INTEGER Flag,Storeflag[15],Stackingflag[15]
371 COM /Tmdt/Datetime$[20],Time$[20]
372 COM /Set_up/ INTEGER Setflag
373 INTEGER N
374 ON KEY 0 CALL Not_used
375 ON KEY 1 CALL Not_used
376 ON KEY 2 CALL Not_used
377 GOTO 379
378 PAUSE
379 ON KEY 3 LABEL Storeflag$,2 CALL Storeflag
380 ON KEY 4 CALL Not_used
381 ON KEY 5 CALL Not_used
382 ON KEY 6 CALL Not_used
383 ON KEY 7 CALL Not_used
384 ON KEY 8 LABEL "STORED NR:"&VAL$(Nr) CALL Not_used
385 ON KEY 9 CALL Not_used
386 Nad=715
387 ASSIGN @Nic TO 715
388 OUTPUT Nad;"FLDDL(3)=5" !SET THE FEILD DELIMITER TO COMMA
389 OUTPUT Nad;"FLDLEN(3)=10" !SET THE FIELD LENGTH FOR NUMERIC PRECISION=10
390 OUTPUT Nad;"OMODE=1" !SET FOR DATA TRANSFER ONLY
391 OUTPUT Nad;"FORMAT=1" !SET FOR AN ASCII FORMAT
392 OUTPUT Nad;"LINLEN=80" !SET LINE LENGTH TO 80
393 OUTPUT Nad;"NPTS" !FIND OUT HOW MANY POINTS IN RECORD
394 ENTER @Nic;N
395 ALLOCATE INTEGER Data(1:N),REAL Stackdata(1:N)
396 !
397 CALL Dataacq(Data(*),N)
398 !
399 ON KEY 3 LABEL Storeflag$,5 CALL Storeflag
400 Cont: IF Storeflag$="STORAGE ON" THEN CALL Store_data(Data(*),N)
401 Setflag=1
402 BEEP 523.25,.2
403 BEEP 1046.5,.2
404 BEEP 783.99,.25
405 SUBEND
406 !*****
407 SUB Config_ana
408 !THIS PROGRAM WILL CONFIGURE THE D6100 FOR Q DATA AQUISITION
409 DIM Command$[100],Darm$[4],Filter$(1:4)[12],Ch$[1],Trig_source$[10],Junk$(50)[20]
410 ASSIGN @Ana TO 715
411 !
412 CLEAR SCREEN
413 OUTPUT @Ana;"RESET"
414 INPUT "TO QUIT NOW PRESS 1, Else enter",K
415 IF K=1 THEN GOTO 479
416 DISP "INITIALIZING ANALOGIC: PRESS F2 WHEN ANALOGIC IS READY"
417 BEEP
418 PAUSE
419 DISP
420 OUTPUT @Ana;"DISARM" !SEND DARM COMMAND
421 !
422 FOR I=1 TO 4
423 Ch$=VAL$(I) !PUT THE VALUE OF I INTO A CHARACTER STRING
424 Filter$(I)="FILTER("&Ch$&")=1" !COMBINE FILTER INTO ONE STRING
425 OUTPUT @Ana;Filter$(I) !SEND FILTER COMMANDS TO ANA
426 NEXT I !FOR ALL FOUR CHANNELS
427 !
428 ! THE FOLLOWING COMMANDS CAME FROM "MODEL 6100 UNIVERSAL
429 ! WAVEFORM ANALIZER VOLUME II: COMMAND INDEX" pF7C-1 TO F7C-145
430 ! THE COMMANDS ARE LISTED ALPHABETICALLY

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```

431 !
432 OUTPUT @Ana;"TRIG"      !SET TRIGGER SOURCE TO
433 OUTPUT @Ana;"TRGSRC=7"   !TO EXT. TRIGGER
434 OUTPUT @Ana;"TRGLEV=2"   !SET TRIG. LEVEL TO TWO VOLTS
435 OUTPUT @Ana;"TRGM=1"     !SET MODE TO NORMAL
436 !
437 !
438 !   THESE LINES TAKEN OUT TO SPEED UP THE ARMING PROCESS
439 !
440 !
441 !
442 !
443 !
444 OUTPUT @Ana;"BUFR"      !SELECT BUFFER MENU
445 OUTPUT @Ana;"INPSEL=2"   !SELECT BUF.A2
446 OUTPUT @Ana;"REC=1"     !MODE OFF
447 OUTPUT @Ana;"INPSEL=3"   !SELECT BUF.A3
448 OUTPUT @Ana;"REC=1"     !MODE OFF
449 OUTPUT @Ana;"INPSEL=4"   !SELECT BUF.A4
450 OUTPUT @Ana;"REC=1"     !MODE OFF
451 OUTPUT @Ana;"TMBSEL=2"   !TIMEBASE B
452 OUTPUT @Ana;"INPSEL=2"   !SELECT INPUT CHANNEL 2
453 OUTPUT @Ana;"REC=2"     !MODE ON
454 OUTPUT @Ana;"TMB"       !SELECT TMB MENU
455 OUTPUT @Ana;"TMBSEL=1"   !SELECT TMB A
456 OUTPUT @Ana;"NPTS=16384" !SET #OF POINTS TO 16384
457 OUTPUT @Ana;"PERSRC=3"   !SELECT EXTCLK TTL
458 OUTPUT @Ana;"TMBSEL=2"   !SELECT TMB B
459 OUTPUT @Ana;"NPTS(2)=16384" !SET #OF POINTS TO 16384 FOR 2
460 OUTPUT @Ana;"PERSRC=3"   !SELECT EXTCLK TTL
461 OUTPUT @Ana;"TRCSRC(1)=BUF.A1"!SELECT BUF.A1 AS UPPER SOURCE
462 OUTPUT @Ana;"X"         !SELECT X MENU
463 OUTPUT @Ana;"XSCL(1)=1X1/4" !SET X SCALE TO 1/4
464 OUTPUT @Ana;"XOFF(1)=.7678" !ADJUST TRACE TO START AT RIGHT
465 OUTPUT @Ana;"TRCSRC(2)=BUF.B2"!SELECT BUF.B2 AS LOWER SOURCE
466 OUTPUT @Ana;"XSCL(2)=1X1/4" !SET X SCALE TO 1/4
467 OUTPUT @Ana;"XOFF(2)=.7678" !ADJUST TRACE TO START AT RIGHT
468 OUTPUT @Ana;"TRIG"      !SELECT TRIG MENU
469 OUTPUT @Ana;"HLD OFF=1"  !SET HOLD OFF TO NONE
470 OUTPUT @Ana;"MARK"      !SELECT MARK MENU
471 OUTPUT @Ana;"TRACE=2"    !SELECT TRACE 2
472 OUTPUT @Ana;"MARKER=2"   !SELECT BASELINE
473 OUTPUT @Ana;"BLINE=2"    !TURN BASE LINE ON
474 OUTPUT @Ana;"MARKER=3"   !SELECT CROSSHAIR
475 OUTPUT @Ana;"CROSS=2"    !SET CROSSHAIR MODE TO CENTER
476 OUTPUT @Ana;"TMB"       !SELECT TIMEBASE DISPLAY
477 OUTPUT @Ana;"INTEN=1"
478 OUTPUT @Ana;"LOCAL"     !GOTO LOCAL MODE
479 !
480 SUBEND
481 !*****
482 SUB Store_data(INTEGER Data(*),N)
483 COM /Nicdata/ REAL Vrange2,Period,Freq,Ckldiv
484 COM /Filename/User$,Path$,Date$,INTEGER Nr
485 COM /Flag/ INTEGER Flag,Storeflag[15],Stackingflag[15]
486 COM /Tmdt/Datetime$,Time$
487 COM /Zeros/Nr$[3]
488 INTEGER Resol,Types,Mrows,Ncols,Imagf
489 REAL Number,Resolution
490 DIM Name$[25]
491 Resol=16
492 Nr=Nr+1
493 PRINT Nr
494 Nr$=VAL$(Nr)
495 CALL Add_zero(Nr$)
496 Filename$="/"&Date$&Nr$
497 Ptspercycle=8192/Ckldiv !ASSUMES 8192 POINTS PER CYCLE
498 PRINT "THE FILE IS BEING STORED AT PATH"&Path$&Filename$&".DTA"
499 PRINT ""
500 Datetime$=DATES(TIMEDATE)

```

```

501 Time$=TIME$(TIMEDATE)
502 !
503 CREATE Path$&Filename$&".DTA",0 !CREATE A DOS FILE
504 ASSIGN @File TO Path$&Filename$&".DTA";FORMAT OFF
505 Types=0 !0 FOR PC'S
506 Mrows=N !NUMBER OF ROWS IN MATRIX
507 Ncols=1 !NUMBER OF COLUMNS IN MATRIX
508 Imagf=0 !IMAGINARY FLAG
509 Number=N !REAL VERSION OF N
510 Resolution=Resol !REAL VERSION OF RESOL
511 !
512 ! OUTPUT DATA MATRIX
513 !
514 Name$="qdata"
515 CALL Head(@File,Types,Mrows,Ncols,Imagf,Name$)
516 OUTPUT @File;Data(*)
517 !
518 ! OUTPUT VRANGE2
519 !
520 Name$="vrange2"
521 CALL Head(@File,Types,Mrows,Ncols,Imagf,Name$)
522 OUTPUT @File;Vrange2
523 !
524 ! OUTPUT FREQUENCY
525 !
526 Name$="freq"
527 CALL Head(@File,Types,Mrows,Ncols,Imagf,Name$)
528 OUTPUT @File;Freq
529 !
530 ! OUTPUT PTSPERCYCLE
531 !
532 Name$="ptspercycle"
533 CALL Head(@File,Types,Mrows,Ncols,Imagf,Name$)
534 OUTPUT @File;Ptspercycle
537 !
538 Flag=1
539 Endsub: SUBEND
540 !*****
541 SUB Add_zero(I$)
542 COM /Zeros/Nr$
543 IF (VAL(I$)<10) AND (POS(I$,"0")=0) THEN I$="0"&I$
544 IF I$=Nr$ AND (VAL(I$)<100) THEN I$="0"&I$
545 SUBEND
546 !*****
547 SUB Not_used
548 BEEP
549 SUBEND
550 !*****
551 SUB Storeflag
552 COM /Flag/ INTEGER Flag,Storeflag$[15],Stackingflag$[15]
553 SELECT Storeflag$
554 CASE "STORAGE OFF"
560 Storeflag$="STORAGE ON"
570 CASE "STORAGE ON"
580 Storeflag$="STORAGE OFF"
590 END SELECT
600 Flag=1
610 SUBEND
620 !*****
630 SUB Stackingflag
631 COM /Flag/ INTEGER Flag,Storeflag$[15],Stackingflag$[15]
632 SELECT Stackingflag$
633 CASE "STACKING OFF"
634 Stackingflag$="STACKING ON"
635 CASE "STACKINGFLAG ON"
636 Stackingflag$="STACKING OFF"
637 END SELECT
638 Flag=1
640 SUBEND
650 !*****

```

```

660 SUB Dataacq(INTEGER Data(*),N)
670 !DATA OUTPUT FROM ANALOGIC 6100
680 COM /Nicdata/ REAL Vrange2,Period,Freq,Ckdiv
690 DIM A$(99),Pause_flag$(1)
700 Pause_flag$="N"
710 ON KEY 5 LABEL "READING" CALL Not_used
720 ON KEY 3 LABEL "PAUSE ASAP",2 GOSUB Asap
730 Nad=715
740 ASSIGN @Nic TO 715
750 !
760 OUTPUT Nad;"FLDDLM(3)=5" !SET THE FIELD DELEMETER TO A COMMA
770 OUTPUT Nad;"FLDLEN(3)=10" !SET THE FIELD LENGTH FOR NUMERIC PRECISION=10
780 OUTPUT Nad;"OMODE=1" !SET FOR DATA TRANSFER ONLY
790 OUTPUT Nad;"FORMAT=1" !SET FOR AN ASCII FORMAT
800 OUTPUT Nad;"LINLEN=80" !SET LINE LENGTH TO 80
810 OUTPUT Nad;"TRGSEL=2" !SET THE TRIGGER TO THE ARM TRIGGER
820 OUTPUT Nad;"TRGSRC=7" !SET THE SOURCE TO BE EXTERNAL
830 OUTPUT Nad;"HLDOFF=9" !TURN HOLD ON UNTIL ARMED
840 OUTPUT Nad;"CLRERR" !CLEAR THE ERROR STATUS BIT
850 OUTPUT Nad;"CLRAQU" !CLEAR THE AQUISITION COMPLETE BIT
860 OUTPUT Nad;"CLRKEY" !CLEAR THE KEY CODE BIT
870 OUTPUT Nad;"ERRSRQ=1" !TURN OFF ERROR STATUS BYTE
880 OUTPUT Nad;"KEYSRQ=1" !TURN OFF KEY SERVICE REQUEST
890 OUTPUT Nad;"OUTSRQ=1" !CLEAR BIT 7 OFF STATUS BYTE
900 Start: OUTPUT Nad;"CLRSRQ" !CLEAR THE SERVICE REQUEST BYTE
910 OUTPUT Nad;"AQUSSQ=2" !TURN ON THE AQUISITION COMPLETE BYTE
920 OUTPUT Nad;"ARM" !ARM THE DIGITIZER
930 REPEAT
940 OUTPUT Nad;"SRQ" !REQUEST THE SERVICE REQUEST STATUS BYTE
950 ENTER @Nic;Service !GET STATUS BYTE
960 UNTIL Service>=19 !WAIT TILL DATA ACQD OR ERROR
970 ON KEY 5 LABEL "TRANSFER DATA" CALL Not_used
1380 OUTPUT Nad;"BUF.B2?" !TELL 6100 TO DUMP HEADER IN ASCII
1390 ENTER @Nic;A$ !GET THE HEADER OF CH 2
1400 Count=0
1410 I=0
1420 REPEAT !DECOMPOSE THE HEADER
1430 I=I+1
1440 IF A$(I,1)<>" THEN !NOT START OF A FIELD
1450 GOTO Cnt
1460 ELSE
1470 Count=Count+1
1480 END IF
1490 SELECT Count
1500 CASE 8
1510 Vrange2=VAL(A$(I+1,1)&A$(I+2,1)&A$(I+3,1)&A$(I+4,1)&A$(I+5,1)&A$(I+6,1))
1520 END SELECT
1530 Cnt: UNTIL Count=8
1540 OUTPUT Nad;"FLDDLM(3)=7" !SET THE FIELD DELIMITER TO A LINE FEED
1550 OUTPUT Nad;"FLDLEN(3)=10" !SET THE FIELD LENGTH FOR NUMERIC PREC.=10
1560 OUTPUT Nad;"FORMAT=4"
1570 ALLOCATE INTEGER A2(1:N)
1610 OUTPUT Nad;"SRC(2)"
1620 ENTER @Nic USING "W";A2(*)
1630 PRINT "2"
1640 OUTPUT Nad;"LOCAL" !RESTORES LOCAL CAPABILITY TO ANALOGIC
1660 MAT Data=A2
1670 IF Pause_flag$="Y" THEN
1680 BEEP
1690 PAUSE
1700 Pause_flag$="N"
1710 END IF
1720 GOTO End
1730 Asap: Pause_flag$="Y"
1740 RETURN
1750 End:
1760 SUBEND
1770 !*****
1780 SUB Head(@File,INTEGER Types,Mrows,Ncols,Imagf,Name$)
1810 INTEGER Length

```

```
1820 Length=LEN(Name$)+1
1830 OUTPUT @File;Types,Mrows,Ncols,Imagf,Length,Name$
1840 SUBEND
```

Appendix E

DATA REDUCTION PROGRAM

The following material and program was written by Rainer Moerig based on a prior Matlab version written by Ivan Getting. The data reduction program is written in Fortran and requires any computer and operating system supporting a fortran compiler.

Remarks

Q2DR.EXE

Q2DR.EXE is an executable Fortran program. The name Q2DR means **Q2 Data Reduction**. To some extent Q2DR.EXE is “optimized” for the Q2 experiment.

Differences between Q2DR and qdr2.m (matlab)

Q2DR.EXE is based on the matlab program qdr2.m. In principle all features of the matlab program qdr2.m (including the channel1 fit) are contained in Q2DR.EXE. In addition Q2DR.EXE offers running modes (see: Q2DR.EXE running modes) in which the guesses for the fit-function of channel2 are done automatically.

Old data acquisition program QDANLMAT

The Basic program QDANLMAT was used to control the measurement. Two channels are measured and the data (2-byte integers) are transferred to the computer. QDANLMAT converts the 2-byte integers to 8-byte reals. A similar conversion (2-byte integers to 4-byte integers) is done for other stored parameters. This blows up the size of each stored data file to approximately 262,000 bytes (simply by storing the channel data as 4-byte integers (readable by matlab), would reduce the file size by a factor of 2).

Q2 requirements

Q2 only needs channel2 as 2-byte integers (= 16 bits = resolution of analogic) and a few of the many stored parameters. Therefore, QDANLMAT was modified.

New data acquisition program Q2DA

Q2DA (**Q2 Data Acquisition**) is basically the same program as QDANLMAT, but only channel2 is transferred to the computer. The data are **not** converted. They are stored as they come in from the analogic as 2-byte integers. The parameters VRANGE2, FREQUENCY, and PTSPERCYCLE are stored (VRANGE2 is unnecessarily stored). Each data file has a size of approximately 33,000 bytes.

Q2DA reads the frequencies to be measured from the file Q2DA.FRE. Q2DA writes information such as date, # of first file, # of files to be reduced, 3 character identifier, sample type (20 char.), saturant (20 char.), and comments (50 char.) to the file Q2DR.INF. These information are read by Q2DR.EXE.

When the data acquisition is finished, a QUIT command in Q2DA transfers control back to DOS (necessary to run Q2DA in Q2RUN.BAT).

Q2RUN.BAT

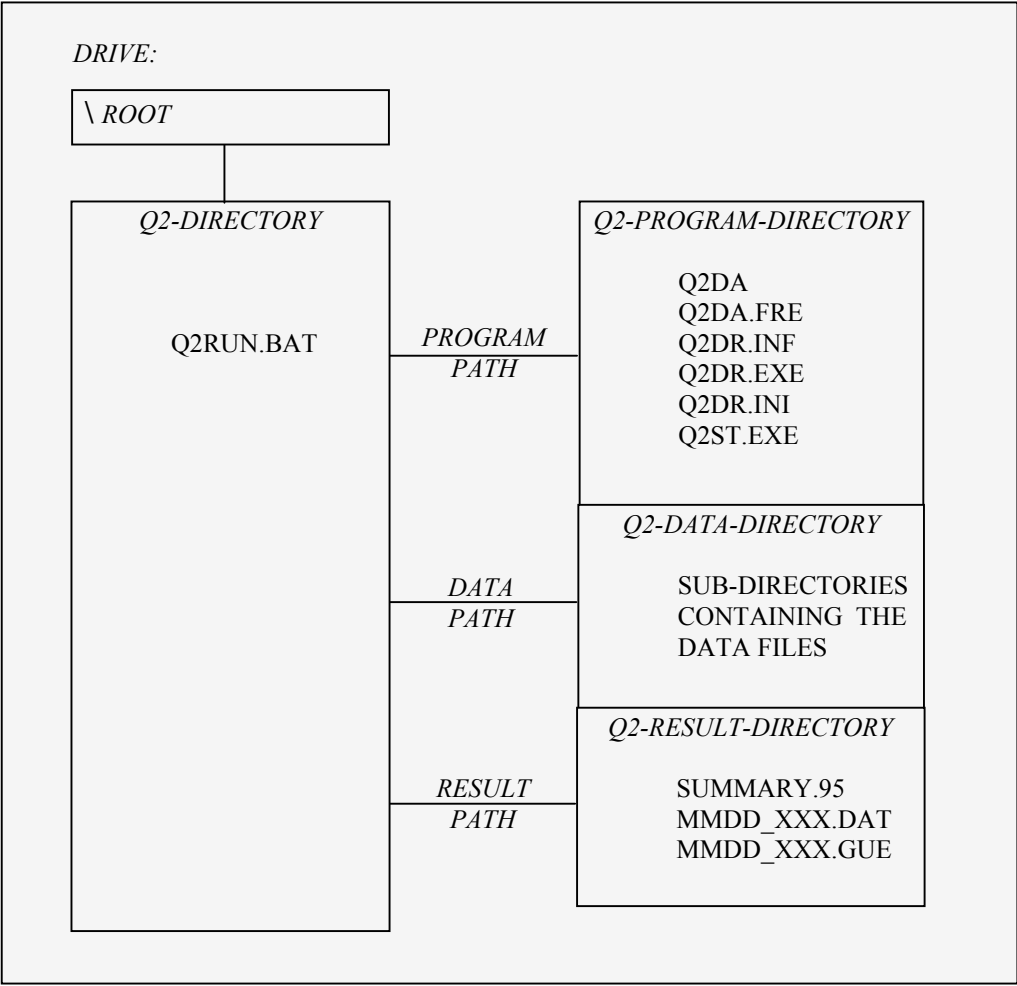
Q2RUN.BAT is the file that has to be executed to take a measurement with a subsequent data reduction.

Q2RUN.BAT calls HTB.BAT. Q2DA has to be loaded and started manually. When the measurement is started (option 4 in Q2DA), control is given to HTB executing Q2DA. At the end of the data acquisition the QUIT command returns control to DOS, that means to Q2RUN.BAT. Q2RUN.BAT now calls Q2DR.EXE. If Q2DR.EXE runs in an automatic mode (see: Q2DR.EXE running modes), the guesses and the data reduction are done automatically, otherwise the guesses have to be made manually (similar to qdr2.m).

Q2RUN.BAT can be called from any directory because the Q2-directory containing Q2RUN.BAT is given in the path of the Q2-autoexec files.

Directory structure

The Figure below shows a possible (current, May 1995) Q2-directory structure. The files in this directory structure are described below:



<i>Q2-DIRECTORY:</i>	
Q2RUN.BAT	is the batch file to run a Q2 experiment
with subsequent data reduction	
<i>Q2-PROGRAM-DIRECTORY:</i>	
Q2DA	is the Q2 Data Acquisition
program (Basic)	
Q2DA.FRE	contains the number of frequencies and
the frequencies to be	
measured.	This file is read by
Q2DA.	
Q2DA.INF	contains information such as date, file
number, sample, saturant etc.	
This file is created by Q2DA. Q2DR.EXE reads the information to	

the summary.95 file.
 Q2DR.EXE is the Q2 Data Reduction program
 (Fortran)
 Q2DR.INI contains the Q2DR.EXE
 running mode, the drive, the program path,
 the data path, and the result path.
 Q2ST.EXE is a program to change the Q2DR.EXE
 running mode, the drive, or the
 paths. Because Q2ST.EXE checks your input for validity, it is
 recommended to make changes in
 Q2DR.INI by running Q2ST.EXE.

Q2-DATA-DIRECTORY:

contains the data files
 in sub-directories.

Q2-RESULT-DIRECTORY:

SUMMARY.95 is a summary file containing information about
 the measurements
 done. This file is updated by Q2DR.EXE.
 MMDD_XXX.DAT is the result file. MMDD is the date, XXX is a 3
 character identifier. This
 file contains the data file number, the frequency, the parameters
 A1 through A7, the standard deviation, and the number of
 iterations done in the
 Levenberg-Marquardt algorithm. A -1 indicates that
 convergence was not met.
 MMDD_XXX.GUE is the guess file. It contains the data file number,
 the frequency, and the
 guesses for the parameters A1 through A7.

Q2DR.EXE Running Modes

Q2DR.EXE can be run in three different running modes. They are manu, auto, and auma.

manu

Manu means manual. The guesses for the parameters A1 through A7 have to be made manually by fitting the fit-function to the measured signal.

auto

Auto means automatic. The guesses are done automatically. If no convergence is reached in the Levenberg-Marquardt algorithm, the guessing and the reduction procedures are repeated once again with a slightly different guessing method. Irrespectively whether or not convergence is reached, after this second run, the results are stored and the next file is processed.

auma

Auma means automatic/manual. The guesses are done first automatically. If no convergence is reached for a file, the same second run is done as in the auto mode.

If all files are correctly reduced, the program terminates. If some files are incorrectly reduced, the program waits. By pressing ENTER to continue, the program switches into the manu mode. The files incorrectly reduced in the automatic mode can be reduced now once again by making the guesses manually.

In the result file the data resulting from the unsuccessful automatic run are replaced by the data resulting from the manual run.

Q2DR source code files

Q2DR.EXE is built up by the 7 Fortran source code files:

Q2MAIN.FOR
 Q2DATA.FOR
 Q2GUESSA.FOR
 Q2GUESSM.FOR
 Q2CH1FIT.FOR
 Q2CH2FIT.FOR
 Q2RESULT.FOR

These Fortran source code files can be changed by using any editor. The modified files must be saved as ASCII files. The extension FOR is necessary because the Fortran compiler uses this extension to recognize a file as a source code file. To create the executable Q2DR.EXE file, the source code files must be compiled and linked. This must be done with the Microsoft Fortran Powerstation Compiler (FL32) and Linker (LINK32) because graphics- and run-time-routines called in the source code files are special tools of Microsoft Fortran.

The following table lists the content of the different source code files (compare flow diagram):

source code file	contents
Q2MAIN.FOR	main program Q2DR and subroutine INFO
Q2DATA.FOR	subroutines FILES, Q2DATA, and SCALE
Q2GUESSA.FOR	subroutines GUESSA and POLY2
Q2GUESSM.FOR	subroutines GUESSM, START, GRAPHICSMODE,
WINDOS,	MINMAX, BACKGRD, SIGFIT, DRAW, COPY, CURSOR, ITEST,
	WRONG, ACTUG5, and ENDPROGRAM
Q2CH1FIT.FOR	subroutines CHN1FIT and QR
Q2CH2FIT.FOR	subroutines CHN2FIT, MARQ, HESSMAT, and
GAUSSJ	
Q2RESULT.FOR	subroutine RESULTS

Q2MAIN.FOR

```

program q2dr
  implicit double precision (a-h,o-z)
  character fname*80,fgues*80,fsfit*80,fssfit*80,fsumm*80
  character dat*4,txt*3,sampl*20,satur*20,comm*80
  character drive*2,path1*80,path2*80,path3*80
  common/fna/fgues,fsfit,fssfit,fsumm
  common/inf/dat,txt,sampl,satur,comm
  common/dri/drive,path1,path2,path3
  common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
  *      ppc,resol
  common/con/pi,pi2,w
  common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
  common/mod/garray(126,9),parray(126,11)
  common/gra/igraph,jgraph

  call info(nr1,nr2,imode)

  modus=1
c.....Automatic mode
  if(imode.eq.1.or.imode.eq.2) then
    do 1 i=nr1,nr2
      call files(i,fname,modus)
77   call guessa(fname)
      call chn1fit(fname,stddev1)
      call chn2fit(fname,stddev1,stddev2,indic)
      if(method.eq.1) goto 77
      call results(i,nr1,nr2,stddev1,stddev2,indic,imode)
1    continue
  end if

  if(imode.eq.2) then
    if(nrn.eq.0) then
      stop
    else
      print*,'press ENTER to continue'
      read(*,*)
      ifile=1
      igrph=nrnot(ifile)
      jgraph=nrnot(nrn)
    end if
  end if

c.....Manual mode
  if(imode.eq.0.or.imode.eq.2) then
    if(imode.eq.0) then
      igrph=nr1
      jgraph=nr2
    end if
    modus=0
88   if(imode.eq.2) then
      nr1=nrnot(ifile)
      nr2=nrnot(ifile)
    end if

```

```

do 2 i=nr1,nr2
call files(i,fname,modus)
call guessm(fname,i)
garray(i,1)=dble(i)
garray(i,2)=freq
do 2 j=1,7
garray(i,j+2)=guess(j)
2 continue

if(imode.eq.2) then
if(ifile.ne.nrn) then
ifile=ifile+1
goto 88
else
imode=-2
ifile=1
end if
end if

modus=1
99 if(imode.eq.(-2)) then
nr1=nrnot(ifile)
nr2=nrnot(ifile)
end if

do 3 i=nr1,nr2
call files(i,fname,modus)
do 4 j=1,7
guess(j)=garray(i,j+2)
4 continue
call chn2fit(fname,stddev1,stddev2,indic)
call results(i,nr1,nr2,stddev1,stddev2,indic,imode)
3 continue

if(imode.eq.(-2)) then
if(ifile.ne.nrn) then
ifile=ifile+1
goto 99
else
imode=-3
call results(i,nr1,nr2,stddev1,stddev2,indic,imode)
end if
end if
end if

stop
end
*****
subroutine info(nr1,nr2,imode)
implicit double precision (a-h,o-z)
character fguess*80,fsfit*80,fssfit*80,fsumm*80
character dat*4,txt*3,sampl*20,satur*20,comm*80,cmode*4
character drive*2,path1*80,path2*80,path3*80,path*80,sfile*80
common/fna/fgues,fsfit,fssfit,fsumm
common/inf/dat,txt,sampl,satur,comm
common/dri/drive,path1,path2,path3
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)

c.....The data acquisition program 'Q2DA' writes information
c to the file 'Q2DR.INF'. These informations are read in.
open(1,file='q2dr.inf')
rewind(1)
read(1,'(a4)')dat
read(1,'(i3)')nr1
read(1,'(i3)')nr
nr2=nr1+nr-1
read(1,'(a3)')txt
if(imode.eq.1) write(*,'(1x,a4,1x,i3,1x,i3,1x,a3)')dat,nr1,nr2,txt

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        if(txt.ne.'xxx'.and.txt.ne.'XXX') then
            read(1,'(a20)')sampl
            read(1,'(a20)')satur
            read(1,'(a80)')comm
            if(imode.eq.1) then
                write(*,'(1x,2a20)')sampl,satur
                write(*,'(1x,a80)')comm
            end if
        end if
    end if
close(1)

c.....read MODE, DRIVE, and PATHS from file 'Q2DR.INI'
c
c.....Q2DR.EXE runs automatically (auto), manually (manu), or
c    first automatically and then manually (auma).
    open(1,file='q2dr.ini')
    rewind(1)
    read(1,'(a4)')cmode
    if(cmode.eq.'auto') imode=1
    if(cmode.eq.'manu') imode=0
    if(cmode.eq.'auma') imode=2
    if(imode.lt.0.or.imode.gt.2)
        *stop 'Error: MODE (auto, manu, or auma) is not defined (Q2DR.INI)'
c.....drive
    read(1,'(a2)')drive
c.....paths
    read(1,'(a80)')path
    path1=drive//path(1:len_trim(path))
    read(1,'(a80)')path
    path2=drive//path(1:len_trim(path))
    read(1,'(a80)')path
    path3=drive//path(1:len_trim(path))
    read(1,'(a80)')sfile
    close(1)

c.....The guesses are written to 'fgues'.
    fgues=path3(1:len_trim(path3))//dat/'_'/txt/'.'.gue'

c.....The parameters resulting from chn1fit are written to 'fsfit'.
c    In this Q2-version neither the channel1-fit is done nor the
c    results are stored.
    fsfit=path3(1:len_trim(path3))//sfit.dat'

c.....The parameters resulting from chn2fit are written to 'fssfit'.
    fssfit=path3(1:len_trim(path3))//dat/'_'/txt/'.'.dat'

c.....The info data are written to 'fsumm'.
    fsumm=path3(1:len_trim(path3))//sfile(1:len_trim(sfile))

    pi=dble(4.)*datan(dble(1.))
    pi2=dble(2.)*pi

    method=0

    return
end

```

Q2DATA.FOR

```

*****
include 'flib.f'
*****
subroutine files(i,fname,modus)
include 'flib.f'
integer*4 length,handle
character fname*80,pfad*80,nr*3,ext*4
character dat*4,txt*3,sampl*20,satur*20,comm*80
character drive*2,path1*80,path2*80,path3*80
common/inf/dat,txt,sampl,satur,comm

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```

common/dri/drive,path1,path2,path3
record / file$info / info

pfad=path2(1:len_trim(path2))//dat/'_'
ext='.dta'

if(i.lt.10) then
  nr='00'//char(i+48)
elseif(i.ge.10.and.i.lt.100) then
  i1=i/10
  i2=mod(i,10)
  nr='0'//char(i1+48)//char(i2+48)
elseif(i.ge.100) then
  i1=i/100
  i2=mod(i,100)
  i3=mod(i2,10)
  i2=i2/10
  nr=char(i1+48)//char(i2+48)//char(i3+48)
end if
fname=pfad(1:len_trim(pfad))//dat/'_'//nr//ext

c.....Check whether data file exists
handle=file$first
length=getfileinfoqq(fname,info,handle)
if (handle.eq.file$last.or.handle.eq.file$error) then
  print 10,fname
  print*, 'FILE does not exist ! '
  stop
end if
if (info.length.eq.0) then
  print 10,fname
  print*, 'File is empty ! '
  stop
end if
10 format(1x,a80)

call q2data(fname,modus)

call scale

return
end
*****
subroutine q2data(fname,modus)
implicit double precision (a-h,o-z)
character fname*80,vnamei(20)*1,vname*20
integer*2 inum(10),idat
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/con/pi,pi2,w

c.....CHANNEL1 and VRANGE1 (and several other unnecessary parameters)
c  are not stored by QDANLMAT
rn=16384.d0
vrange1=1.d0
clkdiv=1.d0
resol=16.d0

if(modus.eq.1) print 90,fname(1:len_trim(fname)), ' loading '
90 format(1x,a28,a11)

open(99,file=fname,form='binary')
rewind(99)

do 1 k=1,4

do 2 i=1,5
  read(99)inum(i)
2 continue

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do 3 i=1,inum(5)
  read(99)vnamei(i)
3 continue

do 4 i=1,inum(5)-1
  vname(i:i)=vnamei(i)
4 continue

if(k.eq.1) then
  do 5 j=1,inum(3)
    do 5 i=1,inum(2)
      read(99)idat
      qdata(i,j+1)=dble(idat)
5 continue
else
  read(99)realnr
end if

if(k.eq.2) then
  vrange2=realnr
elseif(k.eq.3) then
  freq=realnr
elseif(k.eq.4) then
  ppc=realnr
end if

1 continue

close(99)

w=pi2/ppc

return
end
*****
subroutine scale
implicit double precision (a-h,o-z)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*   ppc,resol
dimension scal(2)
reso=dble(2.)*resol
scal(1)=vrange1/reso
scal(2)=vrange2/reso
do 1 i=1,int(rn)
  do 1 j=1,2
    qdata(i,j)=qdata(i,j)*scal(j)
1 continue
return
end

```

Q2GUESSA.FOR

```

*****
subroutine guessa(fname)
implicit double precision (a-h,o-z)
character fname*80
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*   ppc,resol
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(128)
dimension turn(256,2),pturn(3,2),delta(2),it1(2),fit1(2)
dimension izero(16384)

print 90,fname,' guessing'
90 format('+',a28,a11)

c.....Maximum/Minimum
qmax=-1.d20
qmin=1.d20
do 1 i=1,int(rn)

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    qmax=dmax1(qmax,qdata(i,2))
    qmin=dmin1(qmin,qdata(i,2))
1 continue

c.....Guesses g3 and g4
    guess(3)=(qmax-qmin)/2.d0
    guess(4)=(qmax+qmin)/2.d0

c.....Arcsin of scaled signal
    do 2 i=1,int(rn)
        arg=(qdata(i,2)-guess(4))/guess(3)
        if(dabs(arg).gt.1.d0) arg=dsign(1.d0,arg)
        qdata(i,3)=dasin(arg)
    2 continue

c.....Quadratic function fit to arcsin of signal
    call poly2(kiend,m)

c.....Zeros in derivative
    prodold=1.d0
    nzero=0
    do 3 i=1,kiend-1
        prodnew=qdata(i,6)*qdata(i+1,6)
        if(i.gt.1) prodold=qdata(i-1,6)*qdata(i,6)
        if(prodnew.le.0.d0.and.prodold.gt.0.d0) then
            nzero=nzero+1
            if(dabs(qdata(i,6)).le.dabs(qdata(i+1,6))) then
                izero(nzero)=i
                if(izero(nzero-1).eq.i) izero(nzero)=i+1
            else
                izero(nzero)=i+1
                if(izero(nzero-1).eq.(i+1)) izero(nzero)=i+2
            end if
        end if
    3 continue

c.....Turn around points and guess g2
    sg2=0.d0
    nturn=2*int(rn/ppc)
    do 4 i=1,nturn
        iturn=(2*i-1)*int(ppc)/4
        difmin=1.d20
        do 5 j=1,nzero
            qdata(j,8)=dabs(dble(izero(j)+m-iturn))
            difmin=dmin1(difmin,qdata(j,8))
        5 continue
        do 6 j=1,nzero
            if(dabs(difmin-qdata(j,8)).le.1.d-20) then
                if(j.gt.2) then
                    if(qdata(izero(j-1),5)*qdata(izero(j+1),5).gt.0.d0) then
                        turn(i,1)=qdata(izero(j),4)
                        turn(i,2)=qdata(izero(j),5)
                    elseif(qdata(izero(j+1),5)*qdata(izero(j+2),5).gt.0.d0) then
                        turn(i,1)=qdata(izero(j+1),4)
                        turn(i,2)=qdata(izero(j+1),5)
                    elseif(qdata(izero(j-1),5)*qdata(izero(j-2),5).gt.0.d0) then
                        turn(i,1)=qdata(izero(j-1),4)
                        turn(i,2)=qdata(izero(j-1),5)
                    else
                        turn(i,1)=qdata(izero(j),4)
                        turn(i,2)=qdata(izero(j),5)
                    end if
                else
                    turn(i,1)=qdata(izero(j),4)
                    turn(i,2)=qdata(izero(j),5)
                end if
            end if
        6 continue
        sg2=sg2+dble(iturn)-turn(i,1)
    4 continue

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guess(2)=pi2*sg2/dbl(nturn)/ppc

c.....first part of signal
  istart=2*int(turn(1,1))-1-m
  if(-qdata(istart,6).ge.dble(0.)) then
    pf=1.d0
  else
    pf=-1.d0
  end if
  slope=pf

c.....unfold signal
  vorz=1.d0
  ja=1
  st=0.d0
  do 7 i=1,nturn
    do 8 j=ja,nzero
      st=st+pf
      if(dabs(turn(i,1)-qdata(izero(j),4)).lt.1.d-20) then
        if(i.eq.1) then
          if(pf*qdata(izero(1),5).gt.0.d0) st=st-pf
          turn(i,2)=st*pi+(-1)**st*turn(i,2)
          if(turn(i,2).lt.0.d0) then
            vorz=-1.d0
            turn(i,2)=vorz*turn(i,2)
          end if
          pf=-pf
          ja=j
          goto 7
        else
          st=st-dble(2.)*pf
          turn(i,2)=st*pi+(-1)**st*turn(i,2)
          turn(i,2)=vorz*turn(i,2)
          pf=-pf
          ja=j
          goto 7
        end if
      end if
    8 continue
  7 continue

c.....check for 2pi-jumps
  frac=dble(.75)
  if(method.eq.0) then
    do 30 i=3,nturn-2
      hope=dble(2.)*turn(i,2)-turn(i-2,2)
      diff=hope-turn(i+2,2)
      if(dabs(diff).gt.frac*pi2) then
        idiff=int(diff/pi2)
        if(idiff.eq.0) idiff=int(dsign(dble(1.),diff))
        do 31 j=i+2,nturn
          turn(j,2)=turn(j,2)+dble(idiff)*pi2
        31 continue
      end if
    30 continue
  else
    diff=turn(3,2)-turn(1,2)
    if(dabs(diff).gt.frac*pi2) then
      difsum=0.d0
      do 23 i=3,5,2
        difsum=difsum+dabs(turn(i+2,2)-turn(i,2))
      23 continue
      difsum=difsum/dbl(2.)
      if(difsum.le.dabs(diff)) then
        turn(1,2)=turn(1,2)+dsign(1.d0,diff)*pi2
      else
        do 9 i=3,nturn,2
          turn(i,2)=turn(i,2)-dsign(1.d0,diff)*pi2
        9 continue
      end if
    end if
  end if

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```

end if
diff=turn(4,2)-turn(2,2)
if(dabs(diff).gt.frac*pi2) then
  difsum=0.d0
  do 24 i=4,6,2
    difsum=difsum+dabs(turn(i+2,2)-turn(i,2))
24  continue
    difsum=difsum/dbl(2.)
    if(difsum.le.dabs(diff)) then
      turn(2,2)=turn(2,2)+dsign(1.d0,diff)*pi2
    else
      do 10 i=4,nturn,2
        turn(i,2)=turn(i,2)-dsign(1.d0,diff)*pi2
10  continue
    end if
  end if

  do 11 k=0,1
    do 11 i=3+k,nturn-2,2
      np=0
      sx=0.d0
      sx2=0.d0
      sy=0.d0
      sxy=0.d0
      do 12 j=1+k,i,2
        np=np+1
        sx=sx+turn(j,1)
        sx2=sx2+turn(j,1)*turn(j,1)
        sy=sy+turn(j,2)
        sxy=sxy+turn(j,1)*turn(j,2)
12  continue
      den=np*sx2-sx*sx

      if(dabs(den).lt.1.d-300) goto 11

      sl=(np*sxy-sx*sy)/den
      bl=(sx2*sy-sx*sxy)/den
      hope=sl*turn(i+2,1)+bl
      diff=hope-turn(i+2,2)
      if(dabs(diff).ge.frac*pi2) then
        idiff=int(diff/pi2)
        if(idiff.eq.0) idiff=int(dsign(dbl(1.),diff))
        do 13 j=i+2,nturn,2
          turn(j,2)=turn(j,2)+dbl(idiff)*pi2
13  continue
        end if
      end if
    end if
  end if

c.....may be there is a 2pi-jump between each two turn around points
  sum=0.d0
  do 17 i=3,4
    do 17 j=i,nturn,2
      sum=sum+dabs(turn(j,2)-turn(j-2,2))
17  continue
    sum=sum/(nturn-2)
    if(sum.gt.frac*pi2) then
      dif1=turn(3,2)-turn(1,2)
      vorz1=dsign(1.d0,dif1)
      dif2=turn(4,2)-turn(2,2)
      vorz2=dsign(1.d0,dif2)
      vorz=vorz1*vorz2
      if(dabs(dif1).gt.frac*pi2.and.dabs(dif2).gt.frac*pi2.and.
*        vorz.gt.0.d0) then
        do 18 i=3,4
          do 18 j=i,nturn,2
            turn(j,2)=turn(j,2)-vorz1*dbl(int((j-1)/2))*pi2
18  continue
          end if
        end if
      end if
    end if
  end if

```

```

c.....quadratic function fit to unfolded turn around points
  dn=dble(nturn/2)
  do 19 i=1,2
    sx=0.d0
    sx2=0.d0
    sx3=0.d0
    sx4=0.d0
    sy=0.d0
    syx=0.d0
    syx2=0.d0
    do 20 j=i,nturn,2
      x=turn(j,1)
      y=turn(j,2)
      x2=x*x
      x3=x2*x
      x4=x2*x2
      sx=sx+x
      sx2=sx2+x2
      sx3=sx3+x3
      sx4=sx4+x4
      sy=sy+y
      syx=syx+y*x
      syx2=syx2+y*x2
    20 continue
    det=dn*(sx4*sx2-sx3*sx3)-sx4*sx*sx-sx2*sx2*sx2+
      * dble(2.)*sx3*sx2*sx
    cdet=syx2*(sx3*sx-sx2*sx2)+syx*(sx3*sx2-sx4*sx)+
      * sy*(sx4*sx2-sx3*sx3)
    pturn(1,i)=cdet/det
    cdet=syx2*(sx2*sx-dn*sx3)+syx*(dn*sx4-sx2*sx2)-
      * sy*(sx4*sx-sx3*sx2)
    pturn(2,i)=cdet/det
    cdet=syx2*(dn*sx2-sx*sx)+syx*(sx2*sx-dn*sx3)+
      * sy*(sx3*sx-sx2*sx2)
    pturn(3,i)=cdet/det
  19 continue

c.....guesses g1, g6 and g7
  guess(1)=(pturn(1,1)-pturn(1,2))/pi2
  guess(6)=(pturn(2,1)+pturn(2,2))/dble(2.)
  guess(7)=(pturn(3,1)+pturn(3,2))/dble(2.)

c.....guess g5
  pia1sin=pi*guess(1)*dsin(w+guess(2))+guess(6)+guess(7)
  if(slope.ge.0.d0) then
    arg=qdata(1,3)-pia1sin
  else
    arg=dsign(dble(1.),qdata(1,3))*pi-qdata(1,3)-pia1sin
  end if
  guess(5)=arg

c.....correction of g5 (fitting first or second turn around point)
  do 21 j=1,2
    delta(j)=0.d0
    do 22 i=j,nturn,2
      iturn=(2*i-1)*int(ppc)/4-int(guess(2)*ppc/pi2)
      di=dble(iturn)
      fit=guess(3)*dsin(pi*guess(1)*dsin(w*di+guess(2))+guess(5)+
        * guess(6)*di+guess(7)*di*di)+guess(4)
      if(i.eq.1.or.i.eq.2) then
        it1(i)=iturn
        fit1(i)=fit
      end if
      delta(j)=delta(j)+dabs(qdata(iturn,2)-fit)
    22 continue
  21 continue
  delg5=dmax1(delta(1),delta(2))

  if(dabs(delg5-delta(1)).lt.1.d-20) then

```

```

        iwhere=1
    else
        iwhere=2
    end if
    arg0=(qdata(it1(iwhere),2)-guess(4))/guess(3)
    if(dabs(arg0).gt.1.d0) arg0=dsign(1.d0,arg0)
    phi0=dasin(arg0)
    arg1=(fit1(iwhere)-guess(4))/guess(3)
    if(arg1.gt.1.d0) arg1=dsign(1.d0,arg1)
    phi1=dasin(arg1)
    delg5=dabs(phi0-phi1)

    dit1=dbl(it1(iwhere))
    argu1=pi*guess(1)*dsin(w*dit1+guess(2))+guess(5)+
    *      guess(6)*dit1+guess(7)*dit1*dit1
    argu2=w*pi*guess(1)*dcos(w*dit1+guess(2))+guess(6)+
    *      dbl(2.)*guess(7)*dit1
    argu3=w*w*pi*guess(1)*dsin(w*dit1+guess(2))-dbl(2.)*guess(7)
    fitcur=-guess(3)*(dsin(argu1)*argu2+dcos(argu1)*argu3)

    vorz=(-1.d0)**(iwhere+1)
    if(qdata(it1(iwhere)-m,7).ge.0.d0.and.fitcur.ge.0.d0) then
        if(qdata(it1(iwhere),2).gt.fit1(iwhere)) then
            signum=-vorz
        else
            signum=vorz
        end if
    elseif(qdata(it1(iwhere)-m,7).le.0.d0.and.fitcur.le.0.d0) then
        if(qdata(it1(iwhere),2).gt.fit1(iwhere)) then
            signum=vorz
        else
            signum=-vorz
        end if
    elseif(qdata(it1(iwhere)-m,7).ge.0.d0.and.fitcur.le.0.d0) then
        delg5=pi-phi0-phi1
        signum=vorz
    elseif(qdata(it1(iwhere)-m,7).le.0.d0.and.fitcur.ge.0.d0) then
        delg5=pi-phi0-phi1
        signum=-vorz
    end if
    guess(5)=guess(5)+signum*delg5

    return
end
*****
subroutine poly2(kiend,m)
implicit double precision (a-h,o-z)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol

n=int(ppc/16.d0)+1
m=(n-1)/2
kiend=int(rn)-2*m

sx2=0.d0
sx4=0.d0
do 1 k=1,m
    x2=dbl(k)*dbl(k)
    sx2=sx2+x2
    sx4=sx4+x2*x2
1 continue
sx2=dbl(2.)*sx2
sx4=dbl(2.)*sx4
sn422=dbl(n)*sx4-sx2*sx2

do 2 k=m+1,int(rn)-m
sy=0.d0
syx=0.d0
syx2=0.d0

```

```

do 3 l=-m,m
li=l+k
sy=sy+qdata(li,3)
syx=syx+dbble(l)*qdata(li,3)
syx2=syx2+dbble(l)*dbble(l)*qdata(li,3)
3 continue

ki=k-m
qdata(ki,4)=dbble(k)
qdata(ki,5)=(sx4*sy-sx2*syx2)/sn422
qdata(ki,6)=syx/sx2
qdata(ki,7)=(dbble(n)*syx2-sx2*sy)/sn422
2 continue

return
end

```

Q2GUESSM.FOR

```

*****
include 'fgraph.fi'
*****
subroutine guessm(fname,i)
implicit double precision (a-h,o-z)
include 'fgraph.fd'
integer*1 input
integer*2 allo,run,color
character fname*80
common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
common/para/ixa,ixb,icycle,mcycle
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/gra/igraph,jgraph

allo=0
run=0
call start(run)

if(i.eq.igraph) then
  call graphicsmode()
  call backgrd
end if
call windos

input=0
do 1 while(input.ne.9)

if(input.eq.0) then
  run=run+1
  call clearscreen($gclearscreen)
  call sigfit
  if(allo.eq.0) then
    color=12
    call draw(2,color,fname)
    call copy(allo,run)
  else
    call copy(allo,run)
  end if
  color=2
  call draw(4,color,fname)
end if

call cursor(input)

if(input.eq.8) then
  call start(run)
  call windos

```

```

        allo=0
    end if

1 continue

    if(input.eq.9) call copy(allo,0)

    if(i.eq.jgraph) call endprogram()

c   0:black 1:blue 2:green 3:cyan 4:red 5:magenta 6:brown 7:white
c   8:gray 9:light 10:light 11:light 12:light 13:light 14:yellow 15:bright
c       blue  green  cyan  red  magenta  white
    return
end
*****
    subroutine copy(allo,run)
    implicit double precision (a-h,o-z)
    include 'fgraph.fd'
    integer*1 buffer[allocatable](:)
    integer*2 error,status,allo,run
    integer*4 imsize
    common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
    if(run.eq.0) then
        deallocate(buffer,stat=error)
        return
    end if
    if(allo.eq.0) then
        if(run.gt.1) deallocate(buffer,stat=error)
        imsize=imagesize_w(xmin,ymax,xmax,ymin)
        allocate(buffer(imsize),stat=error)
        if(error.ne.0) then
            status=setvideomode($defaultmode)
            stop 'Error: insufficient memory'
        end if
        call getimage_w(xmin,ymax,xmax,ymin,buffer)
        allo=1
    else
        call putimage_w(xmin,ymax,buffer,$gpset)
    end if
    return
end
*****
    subroutine cursor(input)
    implicit double precision (a-h,o-z)
    include 'fgraph.fd'
    character t0*51,t1*51,t2*15,text(8)*3,ti*3,str*14,empty*14
    character in*10,intin*3
    integer*1 input
    integer*2 dummy2,ipos(8,2),row,col,rowstep,colstep
    record/rcCOORD/ cp
    common/con/pi,pi2,w
    common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
    common/para/ixa,ixb,icycle,mcycle
    data ((ipos(i,j),i=1,8),j=1,2) / 44,45,44,45,44,45,44,45,
    * 4,4,30,30,56,56,82,82 /
    data text / 'A1=','A2=','A3=','A4=','A5=','A6=','A7=','#C=' /
    empty=' '
    colstep=3

    call settextrposition(40,4,cp)
    dummy2=settextrcolor(4)
    t0='0: display, 1-7: parameters, '
    t1='8: # of cycles (#C), 9: next file '
    call outtext(t0)
    call settextrposition(40,55,cp)
    call outtext(t1)

    do 1 i=1,8
    row=ipos(i,1)
    col=ipos(i,2)

```



```

call settextposition(row,col,cp)
ti=text(i)
call outtext(ti)

if(i.lt.8) then
  write(str,'(f14.10)')guess(i)
else
  write(str,'(i5)')icycle
end if
col=col+colstep
call settextposition(row,col,cp)
call outtext(str)
1 continue

call settextposition(42,4,cp)
dummy2=settextcolor(9)
t2='INPUT (0-9) : '
call outtext(t2)
in=' '
do 2 while(len_trim(in).gt.1.or.ichar(in).lt.48.or.
*          ichar(in).gt.57)
  call settextposition(42,19,cp)
  call outtext(empty)
  call settextposition(42,19,cp)
  read(*,'(a10)')in
  if(len_trim(in).gt.1.or.ichar(in).lt.48.or.
*    ichar(in).gt.57) then
    row=42
    col=19
    call wrong(row,col)
  end if
2 continue
input=ichar(in)-48

if(input.gt.0.and.input.lt.9) then
  call settextposition(42,4,cp)
  dummy2=settextcolor(4)
  call outtext(t2)

  row=ipos(input,1)
  col=ipos(input,2)
  dummy2=settextcolor(9)
  call settextposition(row,col,cp)
  ti=text(input)
  call outtext(ti)

  col=col+colstep
  call settextposition(row,col,cp)
  call outtext(empty)

  rowstep=(-1)**input
  row=row+rowstep
  call settextposition(row,col,cp)
  if(input.lt.8) then
    write(str,'(f14.10)')guess(input)
  else
    write(str,'(i5)')icycle
  end if
  call outtext(str)
  row=row-rowstep

  col=col+2
  call settextposition(row,col,cp)
  if(input.lt.8) then
    read(*,*)guess(input)
    if(input.ne.5) call actug5
  else
    icycle=-1
    do 3 while(icycle.le.0.or.icycle.gt.mcycle)
      read(*,'(a3)')intin

```

```

        call itest(intin,icycle)
        call settextrposition(row,col,cp)
        if(icycle.le.0.or.icycle.gt.mcycle) call wrong(row,col)
3      continue
    end if

    row=row+rowstep
    col=col-2
    call settextrposition(row,col,cp)
    call outtext(empty)

  end if

  return
end
*****
include 'flib.f'
*****

subroutine wrong(row,col)
include 'fgraph.fd'
include 'flib.fd'
integer*2 row,col
integer*4 duration,frequency
record/rccoord/ cp
duration=500
frequency=440
do 1 i=1,6
  frequency=frequency+(i-1)*100
  call beepqq(duration,frequency)
1 continue
  call settextrposition(row,col,cp)
  call outtext('invalid input, repeat!')
  duration=1000
  call sleepqq(duration)
  call settextrposition(row,col,cp)
  call outtext(' ')
  call settextrposition(row,col,cp)
  return
end
*****

subroutine itest(in,ires)
character in*3
dimension ii(3)
ind=0
iend=len_trim(in)
do 1 i=1,iend
  if(ichar(in(i:i)).lt.48.or.
*   ichar(in(i:i)).gt.57) ind=1
1 continue
  if(ind.eq.0) then
    ires=0
    do 2 i=1,iend
      ii(i)=ichar(in(i:i))-48
      ires=ires+ii(i)*10*(iend-i)
2   continue
  end if
  return
end
*****

subroutine actug5
implicit double precision (a-h,o-z)
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*   ppc,resol
pia1 sin=pi*guess(1)*dsin(w+guess(2))+guess(6)+guess(7)
arg=(qdata(1,2)-guess(4))/guess(3)
if(dabs(arg).gt.1.d0) arg=dsign(1.d0,arg)
arg=dasin(arg)
slope=qdata(2,2)-qdata(1,2)

```

```

    if(slope.ge.0.d0) then
        arg=arg-pia1sin
    else
        arg=dsign(1.d0,arg)*pi-arg-pia1sin
    end if
    guess(5)=arg
    return
end
*****

subroutine draw(nr,color,fname)
implicit double precision (a-h,o-z)
include 'fgraph.fd'
character*9 strx,stry
character fname*80
integer*2 status,color,colo
record/rcoord/ cp
record/wxycoord/wxy
common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
common/para/ixa,ixb,icycle,mcycle
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol

if(color.eq.12) then
    colo=7
    status=setcolor(colo)
    status=rectangle_w($gborder,xmin,ymax,xmax,ymin)
    status=setcolor(colo)
end if

if(color.eq.2) then
    colo=7
    status=setTextcolor(colo)
    call settextposition(3,3,cp)
    write(strx,'(e9.3)')xmin
    write(stry,'(e9.3)')ymax
    call outtext('('//strx//','//stry//')')
    call settextposition(37,98,cp)
    write(strx,'(e9.3)')xmax
    write(stry,'(e9.3)')ymin
    call outtext('('//strx//','//stry//')')
    call settextposition(3,91,cp)
    call outtext(fname(1:len_trim(fname)))
end if

c.....nr=2 is measured signal, nr=4 is fit-function
c.....if you want to look at the measured signal only, activate the if statement below.
c.....The initial number of cycles displayed is icycle2 (subroutine start).
c    if(nr.eq.2) then
        status=setcolor(color)
        call moveto_w(qdata(ixa,3),qdata(ixa,nr),wxy)
c.....nrdisp is the displayed number of points per cycle for the fit-function
        nrdisp=128
        istep=1
        if(color.eq.2) istep=ixb/icycle/nrdisp
        do 1 i=ixa,ixb,istep
            status=lineto_w(qdata(i,3),qdata(i,nr))
1        continue
c    end if

    return
end
*****

subroutine sigfit
implicit double precision (a-h,o-z)
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
common/para/ixa,ixb,icycle,mcycle
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol

```

```

do 1 i=ixa,ixb
di=dble(i)
qdata(i,4)=guess(3)*dsin(pi*guess(1)*dsin(w*di+guess(2))+guess(5)+
*          guess(6)*di+guess(7)*di*di)+guess(4)
1 continue
return
end
*****

subroutine start(run)
implicit double precision (a-h,o-z)
integer*2 run
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
common/para/ixa,ixb,icycle,mcycle
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol

if(run.eq.0) then
icycle1=1
icycle2=2
icycle=icycle2-(icycle1-1)
ixa=(icycle1-1)*int(ppc)+1
ixb=icycle2*int(ppc)
mcycle=int(rn/ppc)
ch2min=1.d300
ch2max=-1.d300
do 1 i=1,int(rn)
qdata(i,3)=dble(i)
ch2min=dmin1(ch2min,qdata(i,2))
ch2max=dmax1(ch2max,qdata(i,2))
1 continue
guess(1)=3.d-6
guess(2)=-.09d0
guess(3)=(ch2max-ch2min)/2.d0
guess(4)=(ch2max+ch2min)/2.d0
guess(6)=0.d0
guess(7)=0.d0
pia1sin=pi*guess(1)*dsin(w+guess(2))+guess(6)+guess(7)
arg=(qdata(1,2)-guess(4))/guess(3)
if(dabs(arg).gt.1.d0) arg=dsign(1.d0,arg)
arg=dasin(arg)
slope=qdata(2,2)-qdata(1,2)
if(slope.ge.0.d0) then
arg=arg-pia1sin
else
arg=dsign(1.d0,arg)*pi-arg-pia1sin
end if
guess(5)=arg
else
icycle2=icycle-(icycle1-1)
ixa=(icycle1-1)*int(ppc)+1
ixb=icycle2*int(ppc)
end if
return
end
*****

subroutine backgrd
include 'fgraph.fd'
integer*2 r,g,b
integer*4 color,oldbgd
rgb(r,g,b)=(#3f3f3f.and.(r.or.ishft(g,8).or.ishft(b,16)))
color=rgb(63,63,63)
oldbgd=setbkcolor(color)
return
end
*****

subroutine windos
implicit double precision (a-h,o-z)
include 'fgraph.fd'

```

```

logical*2 switch/.TRUE./
integer*2 status,maxx,maxy
common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
common maxx,maxy

call minmax
call setviewport(50,48,maxx-50,3*(maxy+1)/4-2)
status=setwindow(switch,xmin,ymin,xmax,ymax)

call settetwindow(1,5,60,130)

return
end
*****
subroutine minmax
implicit double precision (a-h,o-z)
common/mima/xmin,xmax,ymin,ymax,ch2min,ch2max
common/para/ixa,ixb,icycle,mcycle
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol

xmin=dbl(xmin)
xmax=dbl(xmax)
ymin=1.d300
ymax=-1.d300
do 1 i=ixa,ixb
  ymin=dmin1(ymin,qdata(i,2))
  ymax=dmax1(ymax,qdata(i,2))
1 continue
vorz=dsign(1.d0,ymin)
ymin=(1.d0-vorz*5.d-2)*ymin
vorz=dsign(1.d0,ymax)
ymax=(1.d0+vorz*5.d-2)*ymax
return
end
*****
subroutine graphicsmode()
include 'fgraph.fd'
integer*2      modestatus,maxx,maxy
record/videoconfig/ myscreen
common      maxx,maxy
c
c  set highest resolution graphics mode
c
modestatus = setvideomode( $xres256color )
if(modestatus.eq.0) stop 'Error: cannot set graphics mode'
c
c  determine the minimum and maximum dimensions
c
call getvideoconfig( myscreen )
maxx = myscreen.numxpixels - 1
maxy = myscreen.numypixels - 1
return
end
*****
subroutine endprogram()
c
c  ENDPROGRAM resets the screen before returning
c
include 'fgraph.fd'
integer*2 dummy
dummy = setvideomode( $defaultmode )
return
end

```

Q2CH1FIT.FOR

```

*****
subroutine chn1fit(fname,stddev1)

```

```

implicit double precision (a-h,o-z)
character fname*80
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(128)
dimension x(5)

c.....Chn1 fit is NOT done !
  print 10,fname,'  fitting '
  10 format('+',a28,a11)
  goto 20
c.....Chn1 fit is NOT done !

  n=int(rn)
  m=5
  do 1 i=1,n
    xi=dble(i)
    wx=w*xi
    qdata(i,3)=dsin(wx)
    qdata(i,4)=dcos(wx)
    qdata(i,5)=1.d0
    qdata(i,6)=xi
    qdata(i,7)=xi*xi
    qdata(i,8)=qdata(i,1)
  1 continue

  call qr(m,n,x)

  sfit(1)=dsqrt(x(1)*x(1)+x(2)*x(2))
  if(x(1).lt.0.d0) sfit(1)=-sfit(1)
  sfit(2)=datan(x(2)/x(1))
  sfit(3)=x(3)
  sfit(4)=x(4)
  sfit(5)=x(5)

  sum=0.d0
  do 2 i=1,n
    xi=dble(i)
    fit=sfit(1)*dsin(w*xi+sfit(2))+sfit(3)+sfit(4)*xi+sfit(5)*xi*xi
    resid1=qdata(i,1)-fit
    sum=sum+resid1*resid1
  2 continue
  stddev1=dsqrt(sum/(n-5))

  print 11,fname,'  fitting  stddev1:',stddev1
  11 format('+',a28,a20,f7.4)

20 return
end
*****
subroutine qr(n,m,x)

c  n=number of unknowns x(i)
c  m=number of equations, m>=n
c  qdata(m,n+2)=model matrix
c  qdata(m,1)=vector of right side, original y(i)
c  x(n)=solution of qdata(m,n+2)*x(n)=qdata(m,1)

implicit double precision (a-h,o-z)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
dimension x(n),p(5)

  do 1 k=1,n
    p(k)=dble(k)
  1 continue
  do 2 k=1,n
    r0=0.d0
    r1=dble(k)

```

```

      if(k+1.le.n) then
        do 3 i=1,m
          r0=r0+qdata(i,int(p(k))+2)*qdata(i,int(p(k))+2)
3        continue
          r2=r0
          do 4 j=k+1,n
            r0=0.d0
            do 5 i=1,m
              r0=r0+qdata(i,int(p(j))+2)*qdata(i,int(p(j))+2)
5            continue
              if(r0.le.r2) then
                r1=dbl(j)
                r2=r0
              end if
            end if
          4        continue
        end if
        r0=p(k)
        p(k)=p(int(r1))
        p(int(r1))=r0
        do 6 i=1,m
          qdata(i,9)=qdata(i,int(p(k))+2)
6        continue
          r1=0.d0
          do 7 i=1,m
            r1=r1+qdata(i,9)*qdata(i,9)
7          continue
            if(dabs(r1).lt.1.d-50) then
              print*,'execution stopped'
              return
            end if
            qdata(int(p(k)),int(p(k))+2)=1.d0
            r0=0.d0
            do 8 i=1,m
              r0=r0+qdata(i,9)*qdata(i,8)
8            continue
              x(int(p(k)))=r0/r1
              do 9 i=1,m
                qdata(i,8)=qdata(i,8)-x(int(p(k)))*qdata(i,9)
9              continue
                if(k+1.le.n) then
                  do 10 j=k+1,n
                    r0=0.d0
                    do 11 i=1,m
                      r0=r0+qdata(i,9)*qdata(i,int(p(j))+2)
11                    continue
                      qdata(int(p(j)),int(p(k))+2)=r0/r1
                      do 12 i=1,m
                        qdata(i,int(p(j))+2)=qdata(i,int(p(j))+2)-
*                          qdata(int(p(j)),int(p(k))+2)*qdata(i,9)
12                    continue
                  10                continue
                end if
              2            continue
            do 13 i=n,1,-1
              r0=0.d0
              if(i+1.le.n) then
                do 14 k=i+1,n
                  r0=r0+qdata(int(p(k)),int(p(i))+2)*x(int(p(k)))
14                continue
                  x(int(p(i)))=x(int(p(i)))-r0
                end if
              13            continue
            return
          end

```

Q2CH2FIT.FOR

```

*****
subroutine chn2fit(fname,stddev1,stddev2,indic)

```

```

implicit integer*4 (i-n)
implicit double precision (a-h,o-z)
character fname*80
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/con/pi,pi2,w
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(128)
dimension a(7),deltaa(7)

print 10,fname,' fitting '
10 format('+',a28,a11)

icount=0
dlamb=1.d-3

m=7
n=int(rn)

do 1 i=1,m
a(i)=guess(i)
deltaa(i)=dble(1.)
1 continue

c.....Starting with a large ISTEP and reducing it while approaching convergence,
c      could speed up chn2fit. ISTEP is introduced but NOT changed during the run.
c      istep=int(ppc)/16
istep=1
istop=0
do 2 while (istop.eq.0)
call marq(a,deltaa,m,dlamb,istep)
icount=icount+1
indi=0
if(dabs(deltaa(2)).lt.1.d-6.and.
*      dabs(deltaa(1)/a(1)).lt.1.d-4) then
istep=1
istop=1
dlamb=dble(0.)
call marq(a,deltaa,m,dlamb,istep)
elseif(icount.ge.8) then
istep=1
istop=1
indi=-1
dlamb=dble(0.)
call marq(a,deltaa,m,dlamb,istep)
end if
2 continue

sum=0.d0
do 3 i=1,n
x=dble(i)
sinarg=pi*a(1)*dsin(w*x+a(2))+a(5)+a(6)*x+a(7)*x*x
qdata(i,1)=a(3)*dsin(sinarg)+a(4)
resid2=qdata(i,2)-qdata(i,1)
sum=sum+resid2*resid2
3 continue
stddev2=dsqrt(sum/(n-7))

do 4 i=1,m
ssfit(i)=a(i)
4 continue

if(icount.lt.8) then
print 20,fname,' fitted   ampl.:',ssfit(1),' phase:',
*      ssfit(2),icount
indic=icount
method=0
else
print 20,fname,' NOT fitted   ampl.:',ssfit(1),' phase:',
*      ssfit(2),indi
indic=indi

```



```

        method=method+1
    end if
20 format('+',a28,a20,f7.3,a9,f7.4,i3)

    return
end
*****
subroutine marq(a,deltaa,m,dlamb,istep)
implicit integer*4 (i-n)
implicit double precision (a-h,o-z)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/con/pi,pi2,w
dimension a(m),deltaa(m)
dimension beta(7),hmat(7,7)

n=int(rn)
istop=0
call hessmat(a,beta,hmat,m,dlamb,n,istep)
do 1 while (istop.eq.0)
    chisq=0.d0
    do 2 i=1,n,istep
        chisq=chisq+qdata(i,1)*qdata(i,1)
2    continue

    if(dabs(dlamb).lt.1.d-20) then
        istop=1
    else
        call gaussj(hmat,m,m,beta,1,1)
        do 3 i=1,m
            deltaa(i)=beta(i)
3        continue

        do 4 i=1,m
            a(i)=a(i)+deltaa(i)
4        continue
        chisqnw=0.d0
        do 5 i=1,n,istep
            x=dble(i)
            sinarg=pi*a(1)*dsin(w*x+a(2))+a(5)+a(6)*x+a(7)*x*x
            qdata(i,1)=qdata(i,2)-a(3)*dsin(sinarg)-a(4)
            chisqnw=chisqnw+qdata(i,1)*qdata(i,1)
5        continue

        if(chisqnw.le.chisq) then
            dlamb=1.d-1*dlamb
            istop=1
        else
            dlamb=1.d1*dlamb
            call hessmat(a,beta,hmat,m,dlamb,n,istep)
        end if
    end if
1 continue
return
end
*****
subroutine hessmat(a,beta,hmat,m,dlamb,n,istep)
implicit integer*4 (i-n)
implicit double precision (a-h,o-z)
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/con/pi,pi2,w
dimension a(m),beta(m),hmat(m,m)

do 1 i=1,n,istep
    x=dble(i)
    sinarg=pi*a(1)*dsin(w*x+a(2))+a(5)+a(6)*x+a(7)*x*x
    ssa=dsin(sinarg)
    csa=dcos(sinarg)
    qdata(i,1)=qdata(i,2)-a(3)*ssa-a(4)

```

```

arg=w*x+a(2)
qdata(i,3)=a(3)*pi*csa*dsin(arg)
qdata(i,4)=a(1)*a(3)*pi*csa*dcos(arg)
qdata(i,5)=ssa
qdata(i,6)=dble(1.)
qdata(i,7)=a(3)*csa
qdata(i,8)=qdata(i,7)*x
qdata(i,9)=qdata(i,8)*x
1 continue

do 2 j=1,m
beta(j)=dble(0.)
do 3 k=1,n,istep
beta(j)=beta(j)+qdata(k,1)*qdata(k,j+2)
3 continue
do 2 i=j,m
hmat(i,j)=dble(0.)
do 4 k=1,n,istep
hmat(i,j)=hmat(i,j)+qdata(k,i+2)*qdata(k,j+2)
4 continue
hmat(j,i)=hmat(i,j)
if(i.eq.j) hmat(i,j)=(dble(1.)+dlamb)*hmat(i,j)
2 continue

return
end
*****
*   W.H. Press, B.P. Flannery, S.A. Teukolsky, W.T. Vetterling   *
*   Numerical Recipes, Cambridge University Press, 1986, p.28-29 *
*****
subroutine gaussj(a,n,np,b,m,mp)
parameter (nmax=7)
implicit integer*4 (i-n)
implicit double precision (a-h,o-z)
dimension a(np,np),b(np,mp),ipiv(nmax),indxr(nmax),indxc(nmax)
do 11 j=1,n
ipiv(j)=0
11 continue
do 22 i=1,n
big=dble(0.)
do 13 j=1,n
if(ipiv(j).ne.1) then
do 12 k=1,n
if(ipiv(k).eq.0) then
if(dabs(a(j,k)).ge.big) then
big=dabs(a(j,k))
irow=j
icol=k
end if
elseif(ipiv(k).gt.1) then
print*,'singular matrix'
return
end if
12 continue
end if
13 continue
ipiv(icol)=ipiv(icol)+1

if(irow.ne.icol) then
do 14 l=1,n
dum=a(irow,l)
a(irow,l)=a(icol,l)
a(icol,l)=dum
14 continue
do 15 l=1,m
dum=b(irow,l)
b(irow,l)=b(icol,l)
b(icol,l)=dum
15 continue
end if

```

```

indxr(i)=irow
indxc(i)=icol
if(a(icol,icol).eq.dble(0.)) then
  print*,'Singular matrix'
  return
end if
pivinv=dble(1.)/a(icol,icol)
a(icol,icol)=dble(1.)
do 16 l=1,n
  a(icol,l)=a(icol,l)*pivinv
16 continue
do 17 l=1,m
  b(icol,l)=b(icol,l)*pivinv
17 continue
do 21 ll=1,n
  if(ll.ne.icol) then
    dum=a(ll,icol)
    a(ll,icol)=dble(0.)
    do 18 l=1,n
      a(ll,l)=a(ll,l)-a(icol,l)*dum
18 continue
    do 19 l=1,m
      b(ll,l)=b(ll,l)-b(icol,l)*dum
19 continue
    end if
21 continue
22 continue
do 24 l=n,1,-1
  if(indxr(l).ne.indxc(l)) then
    do 23 k=1,n
      dum=a(k,indxr(l))
      a(k,indxr(l))=a(k,indxc(l))
      a(k,indxc(l))=dum
23 continue
    end if
24 continue
  return
end

```

Q2RESULT.FOR

```

*****
subroutine results(i,nr1,nr2,stddev1,stddev2,indic,imode)
implicit double precision (a-h,o-z)
character fguess*80,fsfit*80,fssfit*80,fsum*80
character dat*4,txt*3,sampl*20,satur*20,comm*80,resfil*12
character str*37
common/fna/fguess,fsfit,fssfit,fsum
common/inf/dat,txt,sampl,satur,comm
common/dat/qdata(16384,9),rn,vrange1,vrange2,freq,clkdiv,
*      ppc,resol
common/par/method,guess(7),sfit(5),ssfit(7),nrn,nrnot(126)
common/mod/garray(126,9),parray(126,11)
dimension g(9),p(11)

method=0

if(imode.gt.(-2)) then
  if(i.eq.nr1) then
    ibeg=i
    open(1,file=fguess)
c    open(2,file=fsfit)
    open(3,file=fssfit)
c    rewind(1)
    rewind(2)
    rewind(3)
    indisum=0
    if(txt.ne.'xxx'.and.txt.ne.'XXX') then
      open(9,file=fsum,access='append')
      resfil=dat/'_'//txt/'.'dat'

```

```

        write(9,90)resfil,sampl,satur,comm
90      format(1x,a12,6x,2a20,a80)
        close(9)
      end if
    end if

    write(1,10)i,freq,(guess(j),j=1,7)
c      write(2,20)i,freq,(sfit(j),j=1,5),stddev1
    write(3,30)i,freq,(ssfit(j),j=1,7),stddev2,indic

    if(indic.eq.(-1)) then
      indisum=indisum+1
      nrn=indisum
      nrnot(nrn)=i
    end if

    if(i.eq.nr2) then
      iend=i
      close(1)
c      close(2)
      close(3)

      str=' file(s) may be reduced incorrectly !'
      write(*,40)indisum,str
40      format(1x,'E',72('I'),»',/,
*        1x,'^',72(' '),'^',/,
*        1x,'^',16x,i3,a37,16x,'^',/,
*        1x,'^',72(' '),'^',/,
*        1x,'E',72('I'),'_' )
      end if
    end if

    if(imode.eq.(-2)) then
      parray(i,1)=dble(i)
      parray(i,2)=freq
      do 1 j=1,7
        parray(i,j+2)=ssfit(j)
1      continue
      parray(i,10)=stddev2
      parray(i,11)=dble(indic)
    end if

    if(imode.eq.(-3)) then
      open(1,file=fgues)
      open(3,file=fssfit)
      rewind(1)
      rewind(3)
      l=1
      k=nrnot(1)
      do 2 i=ibeg,iend
        read(1,10)integ1,(g(j),j=2,9)
        g(1)=dble(integ1)
        read(3,30)integ1,(p(j),j=2,10),integ2
        p(1)=dble(integ1)
        p(11)=dble(integ2)
        if(integ1.ne.k) then
          do 3 j=1,9
            garray(i,j)=g(j)
            parray(i,j)=p(j)
3          continue
          do 4 j=10,11
            parray(i,j)=p(j)
4          continue
        else
          l=l+1
          k=nrnot(l)
        end if
2      continue

      rewind(1)

```

```

rewind(3)

do 5 i=ibeg,iend
  write(1,10)int(garray(i,1)),(garray(i,j),j=2,9)
  write(3,30)int(parray(i,1)),(parray(i,j),j=2,10),
*                                int(parray(i,11))
5  continue
  close(1)
  close(3)
end if

10 format(1x,i3,2f9.4,f10.5,2f8.3,f9.3,2e11.4)
c 20 format(1x,i3,7e12.5)
30 format(1x,i3,2f9.4,f10.5,2f8.3,f9.3,2e11.4,f7.4,i3)

return
end

```