Power Spectra

Mitch Withers, Res. Assoc. Prof., Univ. of Memphis

See Aster and Borchers, Time Series Analysis, chapter 4.

THE UNIVERSITY OF MEMPHIS. Center for Earthquake Research and Information

$$E = \int_{-\infty}^{\infty} |x(t)|^2 dt$$

2



$$E = \int_{-\infty}^{\infty} |x(t)|^2 dt$$

Parseval's Theorem tells us, $\int_{-\infty}^{\infty} |x(t)|^2 dt = \int_{-\infty}^{\infty} |X(f)|^2 df$



$$E = \int_{-\infty}^{\infty} |x(t)|^2 dt$$

Parseval's Theorem tells us, $\int_{-\infty}^{\infty} |x(t)|^2 dt = \int_{-\infty}^{\infty} |X(f)|^2 df$

It is also sometimes useful to know the energy per unit frequency, $|X(f)|^2$, the Energy Spectral Density.



$$E = \int_{-\infty}^{\infty} |x(t)|^2 dt$$

Parseval's Theorem tells us, $\int_{-\infty}^{\infty} |x(t)|^2 dt = \int_{-\infty}^{\infty} |X(f)|^2 df$

It is also sometimes useful to know the energy per unit frequency, $|X(f)|^2$, the Energy Spectral Density.

This is also the FT of the autocorrelation,

 $\Phi(f)\Phi^*(f) = F[\phi(t) * \phi^*(-t)] = F[\phi(t) \star \phi^*(t)]$



$$= \left(\int_{-\infty}^{\infty} \phi(t) e^{-i2\pi ft} dt\right)^* = \int_{-\infty}^{\infty} \phi^*(t) e^{-i2\pi f(-t)} dt$$



$$= \left(\int_{-\infty}^{\infty} \phi(t) e^{-i2\pi ft} dt\right)^* = \int_{-\infty}^{\infty} \phi^*(t) e^{-i2\pi f(-t)} dt$$

Let $t = -\tau$, $dt = -d\tau$, and when $t = -\infty, \tau = +\infty$

$$\Phi^*(f) = \int_{\infty}^{-\infty} \phi^*(-\tau) e^{-i2\pi f\tau}(-d\tau)$$



$$= \left(\int_{-\infty}^{\infty} \phi(t) e^{-i2\pi ft} dt\right)^* = \int_{-\infty}^{\infty} \phi^*(t) e^{-i2\pi f(-t)} dt$$

Let $t = -\tau$, $dt = -d\tau$, and when $t = -\infty$, $\tau = +\infty$

$$\Phi^*(f) = \int_{\infty}^{-\infty} \phi^*(-\tau) e^{-i2\pi f\tau}(-d\tau)$$

$$=\int_{-\infty}^{\infty}\phi^*(-\tau)e^{-i2\pi f\tau}d\tau=F[\phi^*(-\tau)]$$

THE UNIVERSITY OF MEMPHIS. Center for Earthquake Research and Information

$$= \left(\int_{-\infty}^{\infty} \phi(t)e^{-i2\pi ft}dt\right)^* = \int_{-\infty}^{\infty} \phi^*(t)e^{-i2\pi f(-t)}dt$$

Let $t = -\tau$, $dt = -d\tau$, and when $t = -\infty$, $\tau = +\infty$

$$\Phi^*(f) = \int_{\infty}^{-\infty} \phi^*(-\tau) e^{-i2\pi f\tau}(-d\tau)$$

$$= \int_{-\infty}^{\infty} \phi^*(-\tau) e^{-i2\pi f\tau} d\tau = F[\phi^*(-\tau)]$$

Now τ is just a dummy variable of integration so it doesn't matter if we call it τ or t anymore.

$$\Phi^*(f) = F[\phi^*(-t)]$$



And $\phi(t) * \phi^*(-t) = \phi(t) * \phi(t)$ so that,



And $\phi(t) * \phi^*(-t) = \phi(t) * \phi(t)$ so that,

 $\Phi(f)\Phi^*(f) = F[\phi(t) * \phi^*(-t)] = F[\phi(t) \star \phi^*(t)]$



And $\phi(t) * \phi^*(-t) = \phi(t) \star \phi(t)$ so that,

$$\Phi(f)\Phi^*(f) = F[\phi(t) * \phi^*(-t)] = F[\phi(t) \star \phi^*(t)]$$

This says that a function with a sharp narrow autocorrelation function will also have a broad energy spectral density (ESD).



And $\phi(t) * \phi^*(-t) = \phi(t) * \phi(t)$ so that,

$$\Phi(f)\Phi^*(f) = F[\phi(t) * \phi^*(-t)] = F[\phi(t) \star \phi^*(t)]$$

This says that a function with a sharp narrow autocorrelation function will also have a broad energy spectral density (ESD).

For example, purely random noise will have a delta function for it's autocorrelation. So it's ESD will be a broad scalar constant at all frequencies.



And $\phi(t) * \phi^*(-t) = \phi(t) * \phi(t)$ so that,

$$\Phi(f)\Phi^*(f) = F[\phi(t) * \phi^*(-t)] = F[\phi(t) \star \phi^*(t)]$$

This says that a function with a sharp narrow autocorrelation function will also have a broad energy spectral density (ESD).

For example, purely random noise will have a delta function for it's autocorrelation. So it's ESD will be a broad scalar constant at all frequencies.

We often refer to noise that has equal power at all frequencies as "white noise". This stems from the fact that white light contains equal contributions from all colors (or frequencies of electromagnetic radiation). More on noise later.

> THE UNIVERSITY OF MEMPHIS. Center for Earthquake Research and Information

$$E = \lim_{T \to \infty} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \infty$$



$$E = \lim_{T \to \infty} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \infty$$

And in this case, the Fourier Transform doesn't converge so we can't know the ESD.



$$E = \lim_{T \to \infty} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \infty$$

And in this case, the Fourier Transform doesn't converge so we can't know the ESD.

If a signal looks statistically similar (e.g. the mean, variance, etc) over successive windows, or finite length time periods, it is said to be stationary.



$$E = \lim_{T \to \infty} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \infty$$

And in this case, the Fourier Transform doesn't converge so we can't know the ESD.

If a signal looks statistically similar (e.g. the mean, variance, etc) over successive windows, or finite length time periods, it is said to be stationary.

Power is energy per unit time E/T and for stationary signals we can estimate the power of the signal using,

$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$

THE UNIVERSITY OF MEMPHIS. Center for Earthquake Research and Information

$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$



$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$

This converges to the true power as T becomes large, $P = \lim_{T \to \infty} P_T$.



$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$

This converges to the true power as T becomes large, $P = \lim_{T \to \infty} P_T$.

The Power Spectral Density (PSD) is found using,

$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f)$$



$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$

This converges to the true power as T becomes large, $P = \lim_{T \to \infty} P_T$.

The Power Spectral Density (PSD) is found using,

$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f)$$

Where $\Phi_T(f) = F[\phi(t)\Pi(t/T)]$



$$P_T = \frac{1}{T} \int_{-T/2}^{T/2} |\phi(t)|^2 dt = \frac{1}{T} \int_{-\infty}^{\infty} |\phi(t)\Pi(t/T)|^2 dt$$

This converges to the true power as T becomes large, $P = \lim_{T \to \infty} P_T$.

The Power Spectral Density (PSD) is found using,

$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f)$$

Where $\Phi_T(f) = F[\phi(t)\Pi(t/T)]$

The PSD will have units of $u^2/_{Hz}$ so it really is a spectral density that shows how power is distributed over each frequency.



$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f), \text{ where } \Phi_T(f) = F[\phi(t)\Pi(t/T)]$$



$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f), \text{ where } \Phi_T(f) = F[\phi(t)\Pi(t/T)]$$

$$PSD = \frac{1}{T} \left(F[\phi(t)] * F[\Pi(t/T)] \right)^2$$



$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f), \text{ where } \Phi_T(f) = F[\phi(t)\Pi(t/T)]$$

$$PSD = \frac{1}{T} \left(F[\phi(t)] * F[\Pi(t/T)] \right)^2$$
$$= \frac{1}{T} [\Phi(f) * sinc(Tf)]^2$$



$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f), \text{ where } \Phi_T(f) = F[\phi(t)\Pi(t/T)]$$

$$PSD = \frac{1}{T} \left(F[\phi(t)] * F[\Pi(t/T)] \right)^2$$
$$= \frac{1}{T} [\Phi(f) * sinc(Tf)]^2$$

This means that the simple act of selecting a finite length time period to work with is mathematically equivalent to multiplying by a boxcar. And that translates into convolution with the sinc function in the frequency domain.



$$PSD[\phi(t)] = \lim_{T \to \infty} \frac{1}{T} \Phi_T(f) \Phi_T^*(f), \text{ where } \Phi_T(f) = F[\phi(t)\Pi(t/T)]$$

$$PSD = \frac{1}{T} \left(F[\phi(t)] * F[\Pi(t/T)] \right)^2$$
$$= \frac{1}{T} [\Phi(f) * sinc(Tf)]^2$$

This means that the simple act of selecting a finite length time period to work with is mathematically equivalent to multiplying by a boxcar. And that translates into convolution with the sinc function in the frequency domain.

That convolution causes our estimate of the power at each frequency to include contributions of the power from neighboring frequencies weighted by the sinc function.





What to do to reduce this leakage or smearing?



What to do to reduce this leakage or smearing?

Remember the FT of functions with sharp edges are very wiggly and edges don't get much sharper than those of a boxcar.



What to do to reduce this leakage or smearing?

Remember the FT of functions with sharp edges are very wiggly and edges don't get much sharper than those of a boxcar.

Choose a smoother window.

Run Aster matlab program, leaktest.





Recall a decibel = db = 20log(amplitude)



Recall a decibel = db = 20log(amplitude)

Amplitude is proportional to the square root of power.



Recall a decibel = db = 20log(amplitude)

Amplitude is proportional to the square root of power.

 $db = 20 \log \left(PSD^{1/2} \right) = 10 \log (PSD)$

36



Seismic signals often span several orders of magnitude in amplitude. It is customary to plot the PSD as a function of db and log(f).

Recall a decibel = db = 20log(amplitude)

Amplitude is proportional to the square root of power.

$$db = 20 \log \left(PSD^{1/2} \right) = 10 \log (PSD)$$

Different sensors and data may have different units (e.g. velocity or acceleration) so always be sure to properly label your units as "db wrt ${}^{u^2}/{}_{Hz}$ " where u are your units. For example "db wrt ${}^{(micron/_s)^2}/{}_{Hz}$ " or "db wrt ${}^{(m/_{s^2})^2}/{}_{Hz}$ " or "db wrt ${}^{m^2/_{s^4}}/{}_{Hz}$ "



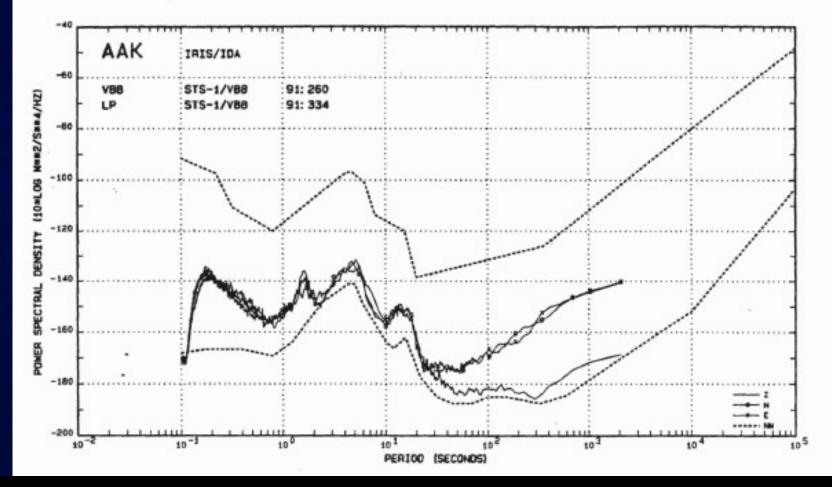


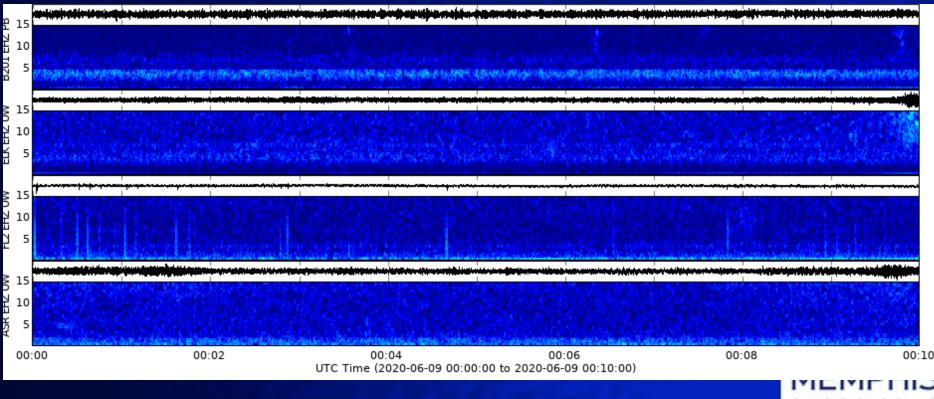
Figure 14: Earth Acceleration Power Spectral Density for background noise at the Ala Archa IRIS/IDA station as a function of period. Z, N, E refer to vertical, north, and east seismometer components. Curves labeled NM are the empirical noise model bounds of Peterson (1994) denoting to extremal PSD values from stations installed around the world. The reference (0 db) level is (1 m/s2)2/Hz. PSD estimates were obtained using Welch's method.

It is sometimes useful to track changes in the PSD over time. We just calculate the PSD estimate multiple times advancing the window by some T for each new PSD. We then plot the results in 3-d using color. This is called a spectragram.



It is sometimes useful to track changes in the PSD over time. We just calculate the PSD estimate multiple times advancing the window by some T for each new PSD. We then plot the results in 3-d using color. This is called a spectragram.

The PNSN has nice continuously updating spectragrams for their volcano monitoring. This is the one for <u>Mt St. Helens</u>.



Center for Earthquake Research and Information If a signal is stationary, it is statistically similar over any given window of fixed length.

We can then improve our PSD estimates by averaging over several windows.



Welch's Method.

If a signal is stationary, it is statistically similar over any given window of fixed length.

We can then improve our PSD estimates by averaging over several windows.



But we lose a lot of data by down weighting at the edge of the windows.

So we allow the windows to overlap.



This is Welch's method. It is common to overlap windows by 50%.



So we allow the windows to overlap.

This is Welch's method. It is common to overlap windows by 50%.

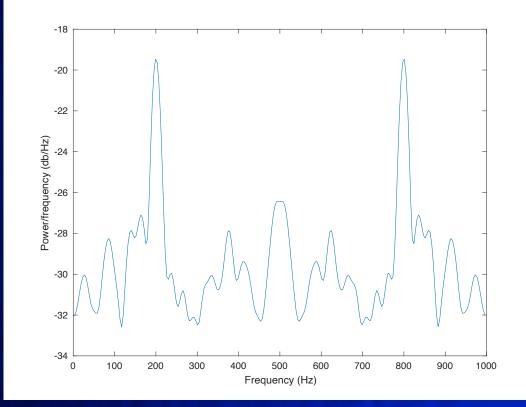
And there's a handy matlab program called pwelch to do this.

[Pxx,F] = pwelch(X,window,Noverlap,nfft,Fs)
Pxx = psd (one-sided is default)
F = frequency vector (Hz)
X = data vector, real
Window = type of window (Hamming is default)
Noverlap = number of points of overlap
nfft = number of points in each fft
Fs = sample rate

[Pxx,F] = pwelch(X,window,Noverlap,nfft,Fs)

Example:

Fs = 1000; t=0:1/Fs:.296; x = cos(2*pi*t*200)+randn(size(t)); %A cosine of 200Hz plus noise. [Pxx,F]=pwelch(x,[],[],[],Fs,'twosided'); %Uses default window, overlap, and NFFT

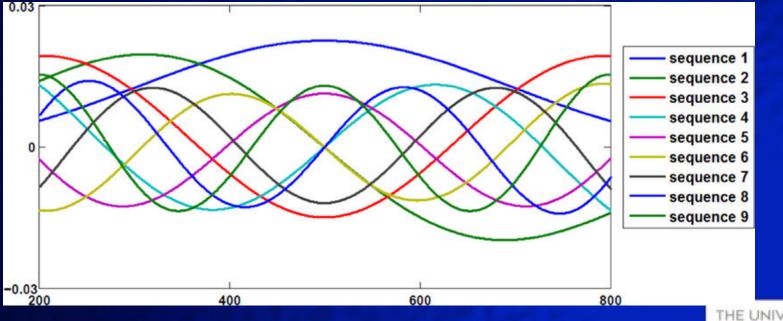


Multitaper Spectral Estimation similarly to Welch's Method computes an average PSD over several windows. This method, while computationally much more expensive, reduces data and bandwidth loss from windowing by operating on the same data segment T.



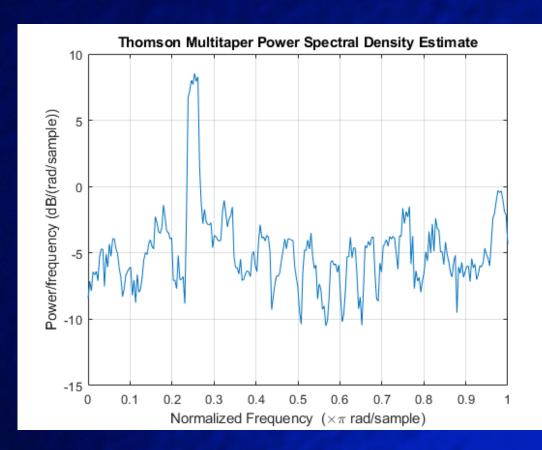
Multitaper Spectral Estimation similarly to Welch's Method computes an average PSD over several windows. This method, while computationally much more expensive, reduces data and bandwidth loss from windowing by operating on the same data segment T.

A suite of prolate spheroidal tapers is used to window the data.



The matlab pmtm command computes the PSD using the multitaper prolate spheroid method. <u>https://www.mathworks.com/help/signal/ref/pmtm.html</u>

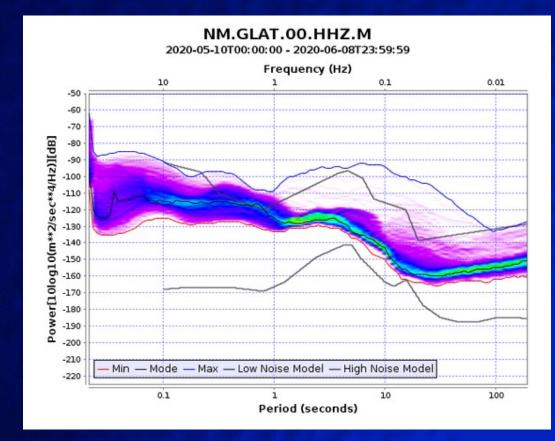
n = 0:319; x = cos(pi/4*n)+randn(size(n)); pxx = pmtm(x);





PSD PDF's are composed of combining many PSD's into an empirical Probability Density Function. They are useful for observing long term trends and signal characteristics of a seismic station.

IRIS has a tool to produce PDF plots for any channel in there archive. Start here and choose PDF Plot. <u>https://ds.iris.edu/mustang/databrowser/</u>





At each frequency, the amplitude values in the PSD's are counted.



At each frequency, the amplitude values in the PSD's are counted.

Each amplitude and frequency in the PDF is colored depending on how many PSD's share that same data point. Hot colors mean many PSD's had that amplitude at that frequency. Cooler colors mean that amplitude at that frequency is less common.



At each frequency, the amplitude values in the PSD's are counted.

Each amplitude and frequency in the PDF is colored depending on how many PSD's share that same data point. Hot colors mean many PSD's had that amplitude at that frequency. Cooler colors mean that amplitude at that frequency is less common.

One can think of extracting a vector of amplitudes at a single frequency and observing it's characteristics. It will have peaks at amplitudes that occur often and troughs at amplitudes that are rare.



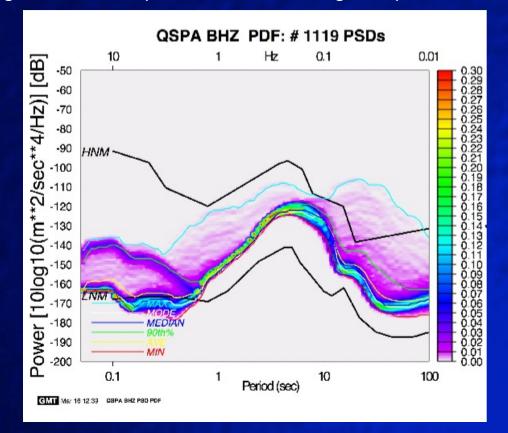
At each frequency, the amplitude values in the PSD's are counted.

Each amplitude and frequency in the PDF is colored depending on how many PSD's share that same data point. Hot colors mean many PSD's had that amplitude at that frequency. Cooler colors mean that amplitude at that frequency is less common.

One can think of extracting a vector of amplitudes at a single frequency and observing it's characteristics. It will have peaks at amplitudes that occur often and troughs at amplitudes that are rare.

It is useful for finding local noise sources, instrument problems (especially ephemeral ones), and other long term trends in the data.

From Aster, "As a final indication of the great utility of the PSD, the Figure (15) shows processed PSDs from a broadband seismometer (Guralp CMG-3Tb) located in a 255-m deep borehole in the polar icecap near the South Pole. A great many of 1-hour data length, 50% overlap, PSDs using a hamming taper, were calculated from the month of May, 2003, and the resulting individual PSDs were used to assemble an empirical probability density function for the signal characteristics at he station The bifurcation of the high frequency noise is caused by intermittent periods where tractors are moving snow near the station. Pink misty areas concentrated around 1 and 20 s are PSDs that include teleseismic earthquake signals. At short periods this is among the quietest stations on Earth. "





The High and Low Noise Models (HNM and LNM) are constructed using many PSD's from the Global Seismic Network (or some superset that includes the GSN) and extracting the highest and lowest amplitudes at each frequency across all PSD's.

The HNM and LNM may contain data points from many different stations. Each point represents the highest or lowest recorded for that frequency in the entire global dataset. (No station will be as good across all frequencies as the LNM though it is possible to construct a station that is noisier than the HNM at all frequencies).

